

The QCD Axion and Unification

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arXiv: 1908.01772

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8 October 2019, University of Zurich

$$\mathcal{L}_{CP}(\theta_Y) = -\overline{q_L} \mathcal{M} e^{i\theta_Y} q_R + \text{h.c.}$$

$$q_L \to e^{i\alpha/2} q_L$$

$$q_R \to e^{-i\alpha/2} q_R$$

$$\mathcal{L}_{\mathcal{CP}}(\theta_{Y}, \theta_{\text{QCD}}) = -\underbrace{\overline{q_{L}} \mathcal{M} e^{i\theta_{Y}} q_{R} + \text{h.c.}}_{\mathcal{L}_{Y}} - \theta_{\text{QCD}} \frac{g_{s}^{2}}{32\pi^{2}} G_{\mu\nu}^{a} \tilde{G}_{a}^{\mu\nu}$$

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$$q_{L} \rightarrow e^{i\alpha/2} q_{L}$$

$$q_{R} \rightarrow e^{-i\alpha/2} q_{R} \quad J \quad U(1)_{A} \quad \left\{ \begin{array}{c} \text{Quark masses} \\ \text{Color anomaly} \end{array} \right.$$

$$\mathcal{L}_{CP}(\theta_{Y}, \theta_{\text{QCD}}) = -\overline{q_{L}} \underbrace{\mathcal{M} e^{i\theta_{Y}} q_{R} + \text{h.c.}}_{\mathcal{L}_{Y}} - \theta_{\text{QCD}} \frac{g_{s}^{2}}{32\pi^{2}} G_{\mu\nu}^{a} \tilde{G}_{a}^{\mu\nu}$$

$$\begin{array}{c} q_{L} \rightarrow e^{i\alpha/2} q_{L} \\ q_{R} \rightarrow e^{-i\alpha/2} q_{R} \end{array} \right\} \quad U(\mathcal{I})_{A} \quad \left\{ \begin{array}{c} \text{Quark masses} \\ \text{Color anomaly} \end{array} \right.$$

$$\begin{array}{c} \text{Quark mass} \\ \text{matrix} \end{array} \quad \text{Topological} \quad \begin{array}{c} \text{Adler, Bell, Jackiw, 1969} \\ \text{Fujikawa, 1979-1980} \end{array}$$

$$\begin{array}{c} \mathcal{L}_{CP}(\theta_{Y}, \theta_{\text{QCD}}) \quad \leftrightarrow \quad \mathcal{L}_{CP}(\theta_{Y} - \alpha, \theta_{\text{QCD}} + N_{i}\alpha) \end{array}$$

$$\begin{array}{c} N = \sum_{\substack{\text{colored} \\ \text{chiral} \\ \text{fields}}} C_{\Psi_{L}^{i}} T_{D}[R_{\Psi_{L}^{i}}] \times \text{mult}[\Psi_{L}^{i}] \end{array}$$

$$\mathcal{L}_{OP}(\bar{\theta}) = -\frac{g_s^2}{32\pi^2} \underbrace{(\arg(\operatorname{Det} M) + \theta_{\operatorname{QCD}})}_{\bar{\theta}} G^a_{\mu\nu} \tilde{G}^{\mu\nu}_a$$

• Strong CP problem:



Crewther, Vecchia, Veneziano, Witten, 1979



• Strong CP problem: Why does QCD seem to conserve CP?





• Strong CP problem: Why does QCD seem to conserve CP?

$$U(1)_A \equiv U(1)_{PQ} \qquad \mathcal{L}_Y = \bar{q}_L \mathcal{M} e^{i\frac{a}{f_a}} q_R + \text{h.c.}$$
$$\mathcal{L}_{\mathcal{P}} \left(\bar{\theta} + \frac{a}{f_a}\right) = -\frac{g_s^2}{32\pi^2} \left(\bar{\theta} + \frac{a}{f_a}\right) G^a_{\mu\nu} \tilde{G}^{\mu\nu}_a + \partial_\mu a \text{ couplings}$$
$$\frac{\text{Peccei-Quinn, 1977}}{2\pi^2} = -\frac{g_s^2}{32\pi^2} \left(\bar{\theta} + \frac{a}{f_a}\right) G^a_{\mu\nu} \tilde{G}^{\mu\nu}_a + \partial_\mu a \text{ couplings}$$

• Strong CP problem: Why does QCD seem to conserve CP?

$$U(1)_{A} \equiv U(1)_{PQ}$$

$$\mathcal{L}_{QP}\left(\bar{\theta} + \frac{a}{f_{a}}\right) = -\frac{g_{s}^{2}}{32\pi^{2}}\left(\bar{\theta} + \frac{a}{f_{a}}\right)G_{\mu\nu}^{a}\tilde{G}_{a}^{\mu\nu} + \partial_{\mu}a \text{ couplings}$$
Peccei-Quinn, 1977
$$2 c^{2} \sqrt{1 - 4m_{u}m_{d}} + 2\left[1\left(\bar{a} + a\right)\right]$$

Wilczek, 1978 Weinberg, 1978

Shift: $a \to a - \bar{\theta} f_a$

• Strong CP problem: Why does QCD seem to conserve CP?

 $U(1)_A \equiv U(1)_{PQ}$

$$\mathcal{L}_{CP}\left(\bar{\theta} + \frac{a}{f_a}\right) = -\frac{g_s^2}{32\pi^2} \underbrace{\left(\bar{\theta} + \frac{a}{f_a}\right)}_{a} G^a_{\mu\nu} \tilde{G}^{\mu\nu}_a + \partial_\mu a \text{ couplings}$$
Peccei-Quinn, 1977

$$V(a) = -m_{\pi}^2 f_{\pi}^2 \sqrt{1 - \frac{4 m_u m_d}{(m_u + m_d)^2} \sin^2 \left[\frac{1}{2} \left(\frac{\bar{\theta} + \frac{a}{f_a}}{f_a}\right)\right]}$$



$$\langle a \rangle = -\bar{\theta} f_a \quad \longrightarrow \quad d_n \propto \frac{a}{f_a} + \bar{\theta} = 0$$

Strong CP problem solved dynamically!



Pseudo-Nambu-Goldstone boson: AXION

Wilczek, 1978 Weinberg, 1978

Axions

• Hypothetical pseudoscalar particle:

$$\mathcal{L}_{a} = \frac{1}{2} (\partial_{\mu} a)^{2} + \frac{a}{f_{a}} \frac{g_{s}^{2}}{32\pi^{2}} G_{\mu\nu} \tilde{G}^{\mu\nu} + \frac{1}{4} a \, g_{a\gamma\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\partial_{\mu} a}{2f_{a}} (c_{q} \bar{q} \gamma^{\mu} \gamma_{5} q)$$



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• Axion mass:

$$V(a) \Rightarrow m_a = 5.70(6)(4) \,\mu \text{eV} \left(\frac{10^{12} \text{ GeV}}{f_a}\right)$$

Grilli di Cortona et al., 2016



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• Axion origin:

$$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y \otimes \underbrace{U(1)_A}_{\text{global}}$$

Peccei-Quinn (PQ) symmetry

 $Axion: Pseudo-Nambu-Goldstone\ boson$



Peccei, Quinn, Weinberg, Wilczek

• Relevant Lagrangian:

$$\mathcal{L}_Y \supset Y_u \overline{Q_L} \,\tilde{H}_u \, u_R + Y_d \,\overline{Q_L} \,H_d \, d_R + \text{h.c.}$$

• PQ symmetry:

$$\begin{aligned} H_u &\to e^{iX_u} H_u & u_R \to e^{iX_u} u_R \\ H_d &\to e^{-iX_d} H_d & d_R \to e^{iX_d} d_R \end{aligned}$$

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$$H_d \to e^{-iX_d} H_d \qquad d_R \to e^{iX_d} d_R$$

• Scalar potential:

$$V \stackrel{\text{PQ}}{=} V(H_u^{\dagger} H_u, H_d^{\dagger} H_d)$$
2 Goldstones
$$\begin{cases} 1 \text{ eaten by } Z_\mu \\ 1 \text{ axion} \end{cases}$$

• Axion coupling to matter:

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$$V \stackrel{\text{SSB}}{\to} \quad H_u \propto e^{i \frac{v_d}{v_u} \frac{a}{v}}, \quad H_d \propto e^{-i \frac{v_u}{v_d} \frac{a}{v}}$$

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• PQ rotation on fermions:

$$\mathcal{L} \supset \frac{a}{v} \frac{g_s^2}{32\pi^2} N G \tilde{G}, \quad N = \frac{1}{2} \left(\frac{v_d}{v_u} + \frac{v_u}{v_d} \right)$$

Original PQ Axion Ruled out

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• Experimental bounds:

$$\underbrace{\operatorname{Br}(K^+ \to \pi^+ + a)}_{3 \times 10^{-5} \left(\frac{v_1}{v_2} + \frac{v_2}{v_1}\right)^2} \xrightarrow{\operatorname{KEK bound}}_{\operatorname{Br}(K^+ \to \pi^+ + \operatorname{nothing})} 3.8 \times 10^{-8}$$

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$$f_a \gg v$$
 Invisible axion models

Dine, Fischler, Srednicki, 1981

Zhitnitsky, 1980

 $S \sim (1, 1, 0)$

"Simple generalisation of the PQ scheme with a harmless axion"

$$\mathcal{L}_Y \supset Y_u \overline{Q_L} \, \tilde{H}_u \, u_R + Y_d \, \overline{Q_L} \, H_d \, d_R + \text{h.c.}$$

• PQ charges:

 $H_u \to e^{iX_u} H_u, \quad H_d \to e^{-iX_d} H_d, \quad S \to e^{iX_s} S$

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• Scalar potential:



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• Axion coupling to matter:

 $\mathcal{L}_Y \supset Y_u \overline{Q_L} \,\tilde{H}_u \, u_R + Y_d \,\overline{Q_L} \,H_d \, d_R + \text{h.c.}$ $H_u \propto e^{2i \frac{v_d^2}{v^2} \frac{a}{v_S}}, \quad H_d \propto e^{-2i \frac{v_u^2}{v^2} \frac{a}{v_S}}, \quad S \propto e^{-i \frac{a}{v_S}}$

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$$\mathcal{L} \supset \frac{a}{v_S} \frac{g_s^2}{32\pi^2} NG\tilde{G}, \quad N = N_f$$

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$$\bullet \quad f_a \equiv v_S \gg v \qquad ()$$

Invisible Axions: KSVZ Kim, 1979

Shifman, Vainshtein, Zakharov, 1980

• "The $U(1)_A$ is spontaneously broken and is realized by the existence of the axion a that does not couple to ordinary quarks at tree level."

$$\mathcal{L}_Y \supset \mathcal{L}_Y^{\mathrm{SM}} + Y_q \,\overline{q_L} \, q_R \, S + M_q \,\overline{q_L} \, q_R + \mathrm{h.c.}$$

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$$V = V(H^{\dagger}H, |S|^{2})$$

$$S = \frac{S_{0} + v_{S}}{\sqrt{2}} \,e^{ia/v_{S}} \qquad \underbrace{\checkmark}_{2} \,\mathrm{Goldstones} \,\left\{ \begin{array}{c} 1 \,\mathrm{eaten} \,\mathrm{by} \, Z_{\mu} \\ 1 \,\mathrm{axion} \end{array} \right.$$

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$$S = \frac{S_{0} + v_{S}}{\sqrt{2}} e^{ia/v_{S}} \longrightarrow 2 \text{ Goldstones } \left\{ \begin{array}{c} 1 \text{ eaten by } Z_{\mu} \\ 1 \text{ axion} \end{array} \right.$$

$$- \frac{2}{\zeta} \left(\sum_{k=0}^{\infty} \mathcal{G} \\ \sum_{k=0}^{\infty} \mathcal{G} \end{array} \right) \Rightarrow \frac{a}{v_{S}} \underbrace{2T_{D}[q]}_{N} \underbrace{\frac{g_{S}^{2}}{32\pi^{2}}}_{32\pi^{2}} G_{\mu\nu} \tilde{G}^{\mu\nu},$$

$$\downarrow \quad f_{a} \equiv v_{S} \gg v$$

"The only requirement is the existence of an anomalous U(1) symmetry broken at a large energy scale"





Irastorza, Redondo, 2018

"The only requirement is the existence of an anomalous U(1) symmetry broken at a large energy scale"



• Astrophysical constraints: $f_a \gtrsim 10^7 \text{ GeV}$





Experimental Status

$$\mathcal{L}_a = \frac{1}{2} (\partial_\mu a)^2 + \frac{a}{f_a} \frac{g_s^2}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu} + \frac{1}{4} a g_{a\gamma\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\partial_\mu a}{2f_a} (c_q \bar{q} \gamma^\mu \gamma_5 q)$$

• Axion halosopes: exploit coherent effects to detect axion DM Axions highly non-relativistic \rightarrow monochromatic photons $\frac{\Delta \omega}{\omega} \sim \sigma_v^2 \sim 10^{-6}$

Experimental Status

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Axion halosopes: exploit coherent effects to detect axion DM
 Axions highly non-relativistic → monochromatic photons



ties $\frac{\Delta \omega}{\omega} \sim \sigma_v^2 \sim 10^{-6}$ If $m_a \subset \omega_c \pm \frac{\omega_c}{Q}$ on $P_s \propto \frac{Q}{m_s} g_{a\gamma\gamma}^2 B_e |G_m|^2 V \rho_a$

Resonantly enhanced conversion

 $a \rightarrow \gamma$


And what if we decouple the detector V from the axion ω_a ?

ABRACADABRA Kahn, Said, Thaler, 2016

A Broadband/Resonant Approach to Cosmic Axion Detection with an Amplifying B-field Ring Apparatus

And what if we decouple the detector V from the axion ω_a ?



 $\vec{B_0}$ generates an $\vec{E_a}$ through Ampère's circuit law

ABRACADABRA



$CASPEr \quad {\rm Graham, \, Rajendran, \, 2013}$

 $Cosmic \ Axion \ Precession \ Experiment$



Detecting an oscillating nuclear-spin-dependent energy shift using **NMR techniques**

CASPEr

 $Cosmic \ Axion \ Precession \ Experiment$

$$\mathcal{L}_{a} = \frac{1}{2} (\partial_{\mu} a)^{2} + \frac{a}{f_{a}} \frac{g_{s}^{2}}{32\pi^{2}} G_{\mu\nu} \tilde{G}^{\mu\nu} + \frac{1}{4} a g_{a\gamma\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\partial_{\mu} a}{2f_{a}} (c_{q} \bar{q} \gamma^{\mu} \gamma_{5} q)$$

$$E < \Lambda_{QCD}$$

$$-\frac{i}{2} g_{d} a \overline{\Psi_{n}} \sigma_{\mu\nu} \gamma_{5} \Psi_{n} F^{\mu\nu} + \frac{g_{aNN} (\partial_{\mu} a) \overline{\Psi_{n}} \gamma^{\mu} \gamma_{5} \Psi_{n}}{\mathcal{L}_{spin}}$$

$$\mathcal{L}_{EDM}$$

$$\mathcal{L}_{EDM}$$

$$\mathcal{L}_{EDM}$$

$$\mathcal{L}_{Spin}$$

$$\mathcal{L}_{S$$

Detecting an oscillating nuclear-spin-dependent energy shift using **NMR techniques**

CASPEr

 $Cosmic \ Axion \ Precession \ Experiment$

CASPEr ELECTRIC

CASPEr WIND



Budker, Graham, Bedbetter, Rajendran, Sushkov, 2014

Strongest magnet 30 T $\Rightarrow m_a < 10^{-6} \text{ eV}$

Motivation





Wise, Georgi, Glashow, 1981





SU(5)

- Rank 4
 15 Weyl d.o.f. fit perfectly in 5 and 10

$$Q_{L} = \begin{pmatrix} u_{L} \\ d_{L} \end{pmatrix}, \ L_{L} = \begin{pmatrix} \nu_{L} \\ e_{L} \end{pmatrix}, \ \begin{pmatrix} u^{c} \\ e_{L} \end{pmatrix}, \ \begin{pmatrix} d^{c} \\ d^{$$

SU(5)

- Rank 4
- 15 Weyl d.o.f. fit perfectly in $\overline{5}$ and 10
- Charge quantisation

$$Q = T_{23} (\equiv \frac{1}{2}\sigma_z) + \sqrt{\frac{5}{3}}T_{24}$$

Matter content $\overline{5}, 10$

SU(5)

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Scalar content

 $5_H, 24_H$

Matter content

 $\bar{5}, 10$

$$5_{H} = \begin{pmatrix} T \\ H_{1} \end{pmatrix} \qquad "SU(5) \stackrel{\langle 24_{H} \rangle}{\to} SM \stackrel{\langle 5_{H} \rangle}{\to} SU(3)_{C} \otimes U(1)_{Q}"$$
$$24_{H} = \begin{pmatrix} \Sigma_{8} - \frac{2}{\sqrt{30}} \Sigma_{0} & \Sigma_{(\bar{3},2)} \\ \Sigma_{(3,2)} & \Sigma_{3} + \frac{3}{\sqrt{30}} \Sigma_{0} \end{pmatrix}$$

$$SU(5) \qquad \begin{array}{c} V_{24} \sim \underbrace{(8,1,0)}_{G_{\mu}} \oplus \underbrace{(1,3,0)}_{W_{\mu}} \oplus \underbrace{(3,2,-5/6)}_{V_{\mu}^{c}} \oplus \underbrace{(\bar{3},2,5/6)}_{V_{\mu}} \oplus \underbrace{(1,1,0)}_{\lambda_{\mu}} \\ \\ 5_{H} \sim \underbrace{(1,2,1/2)}_{H_{1}} \oplus \underbrace{(3,1,-1/3)}_{T} \\ 24_{H} \sim \underbrace{(8,1,0)}_{\Sigma_{8}} \oplus \underbrace{(1,3,0)}_{\Sigma_{3}} \oplus \underbrace{(3,2,-5/6)}_{\Sigma_{(3,2)}} \oplus \underbrace{(\bar{3},2,5/6)}_{\Sigma_{(\bar{3},2)}} \oplus \underbrace{(1,1,0)}_{\Sigma_{24}} \end{array}$$

Scalar content Gaug $5_H, 24_H$

Gauge bosons Matter content 24_V $\overline{5}, 10$

$$5_{H} = \begin{pmatrix} T \\ H_{1} \end{pmatrix} \qquad A_{\mu} = \begin{pmatrix} G_{\mu} - \frac{1}{\sqrt{15}} B_{\mu} & V_{\mu}^{c} \\ V_{\mu} & W_{\mu} + \frac{3}{2\sqrt{15}} B_{\mu} \end{pmatrix}$$
$$24_{H} = \begin{pmatrix} \Sigma_{8} - \frac{2}{\sqrt{30}} \Sigma_{0} & \Sigma_{(\bar{3},2)} \\ \Sigma_{(3,2)} & \Sigma_{3} + \frac{3}{\sqrt{30}} \Sigma_{0} \end{pmatrix}$$







 $SU(5)_{GG}$ • Gauge couplings do not unify Georgi, Glashow 60 50 40 α_1^{-1} $lpha_{
m GUT}^{-1}$ 30 α_2^{-1} α_3^{-1} 20 10 1000 10⁷ **10**¹⁵ 10¹¹ **10**¹⁹ $\alpha_i = \frac{g_i^2}{4\pi}$ $\Lambda_{GUT}(GeV)$

- $SU(5)_{GG}$ Gauge couplings do not unify Wrong fermion mass relations

 $\mathcal{L}_Y \supset Y_1 \,\overline{5} \, 10 \, 5_H^* + Y_3 \, 10 \, 10 \, 5_H \epsilon_5 + \text{h.c.}$

$$M_d = M_e^T = Y_1 \frac{v^*}{2},$$
$$M_u = 4(Y_3 + Y_3^T) \frac{v}{2}$$

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@ GUT scale
$$M_d = M_e^T = Y_1 \frac{v^*}{2},$$

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 - \bullet Neutrinos are massless (as in the SM)



Not allowed by Gauge Symmetry

$SU(5)_{GG}$

- Gauge couplings do not unify
- Wrong fermion mass relations
- Neutrinos are massless (as in the SM)



 $SU(5)_{GG} \otimes U(1)_{PQ}$ Wise, Georgi, Glashow

- Gauge couplings do not unify
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$$24_H \stackrel{U(1)}{\to} 24_H^* \implies 24_H \supset \frac{1}{\sqrt{2}} |\Sigma_0| e^{\frac{ia(x)}{v_{\Sigma}}}$$

	$Y_e \neq Y_d$	$M_{\nu} \neq 0$	Unification	$U(1)_{PQ}$
Renorm.				
SU(5)				
extension?				

 $SU(5)_{GG}$ $\otimes U(1)_{PQ}$

- Gauge couplings do not unify
- Wrong fermion mass relations
- Neutrinos are massless (as in the SM)
- Extra content for PQ!

 $\mathcal{L} \supset Y_5 \,\overline{5} \, 10 \, 5_H^* + Y_{10} \, 10 \, 10 \, 5_H' \, \epsilon_5 + \lambda \, 5_H^* 24_H^2 \, 5_H' + \text{h.c.}$



Wise, Georgi, Glashow

 $SU(5)_{GG} \otimes U(1)_{PQ}$

- Gauge couplings do not unify
- Wrong fermion mass relations
- Neutrinos are massless (as in the SM)
- Extra content for PQ!

 $\mathcal{L} \supset Y_5 \,\overline{5} \, 10 \, 5_H^* + Y_{10} \, 10 \, 10 \, 5_H' \, \epsilon_5 + \lambda \, 5_H^* 24_H^2 \, 5_H' + \text{h.c.}$

	$Y_e \neq Y_d$	$M_{\nu} \neq 0$	Unification	$U(1)_{PQ}$	
$5'_H$	×	×	X	\checkmark	

Wise, Georgi, Glashow



 $45_{H} \sim \underbrace{(8,2,1/2)}_{\Phi_{1}} \oplus \underbrace{(\bar{6},1,-1/3)}_{\Phi_{2}} \oplus \underbrace{(3,3,-1/3)}_{\Phi_{3}} \oplus \underbrace{(\bar{3},2,-7/6)}_{\Phi_{4}} \oplus \underbrace{(3,1,-1/3)}_{\Phi_{5}} \oplus \underbrace{(\bar{3},1,4/3)}_{\Phi_{6}} \oplus \underbrace{(1,2,1/2)}_{H_{2}} \oplus \underbrace{(\bar{1},2,1/2)}_{H_{2}} \oplus \underbrace{(\bar{1},2,1/2)}$



$\mathcal{L} \supset M_{\nu} \mathbf{1}_i \mathbf{1}_i + Y_{\nu}^i \,\overline{5} \,\mathbf{1}_i \,\mathbf{5}_H$





 $\mathcal{L} \supset \lambda \,\overline{5} \,\overline{5} \,10_H + \overline{5} \,10 \left(Y_1^* 5_H^* + Y_2^* 45_H^*\right)$ $+ \mu \,5_H 45_H 10_H^* + \text{h.c.}$





$$M_{\nu} \sim \lambda Y_{\nu} \frac{\langle \Phi \rangle^2}{M_{\Delta_T}^2}$$



$$24 = \begin{pmatrix} \rho_8 - \frac{1}{\sqrt{15}}\rho_0 & \rho_{(\bar{3},2)} \\ \rho_{(3,2)} & \rho_3 + \frac{3}{2\sqrt{15}}\rho_0 \end{pmatrix} \qquad \overbrace{\rho_0}^{\text{type-III}} + \overbrace{\rho_3}^{\text{type-III}} \\ \rho_3 \end{array}$$

 $\mathcal{L} \supset \lambda \mathrm{Tr}\{24^2 24_H\} + \mathrm{h.c.}$

$Y_e \neq Y_d$	$M_{\nu} \neq 0$	Unification	$U(1)_{PQ}$
	1_i	\checkmark	×
15	10_H	\checkmark	X
40H	15_H	\checkmark	X
	24	\checkmark	\checkmark
	1_i	×	\checkmark
$5' + \bar{5}'$	15_H	X	\checkmark
	24	×	\checkmark

 $\mathcal{L} \supset Y_5 \,\overline{5}' 24_H 5'$



	$Y_e \neq Y_d$	$M_{\nu} \neq 0$	Unification	$U(1)_{PQ}$	
		1_i	\checkmark	×	
45_H	15	10_H	\checkmark	×	
	49 <u>H</u>	15_H	\checkmark	×	
KSVZ		24	\checkmark	\checkmark	Adjoint $SU(5)$
					Fileviez Pérez, 2007
		1_i	X	\checkmark	
	$5' + \bar{5}'$	15_H	×	\checkmark	
		24	×	\checkmark	

 \bullet Minimal GUT theory (# of Reps.) where the PQ-symmetry can be realized

• PQ charges:

 $\bar{5} \to e^{-3i\theta}\bar{5}, \qquad 10 \to e^{+i\theta}10, \qquad 5_H \to e^{-2i\theta}5_H,$ $24_H \to e^{-10i\theta}24_H, \qquad 45_H \to e^{-2i\theta}45_H, \qquad 24 \to e^{+5i\theta}24.$

• Charged fermion masses:

 $\mathcal{L} \supset \overline{5} \, 10 \, (Y_1 \, 5_H^* + Y_2 \, 45_H^*) + 10 \, 10 \, (Y_3 \, 5_H + Y_4 \, 45_H) \epsilon_5 + \text{h.c.}$

$$M_E = \frac{1}{2} (Y_1^T v_5^* - 6Y_2^T v_{45}^*),$$

$$M_D = \frac{1}{2} (Y_1 v_5^* + 2Y_2 v_{45}^*),$$

$$M_U = \frac{1}{2\sqrt{2}} (4(Y_3 + Y_3^T) v_5 - 8(Y_4 - Y_4^T) v_{45}).$$

• PQ charges:

 $\bar{5} \to e^{-3i\theta}\bar{5}, \qquad 10 \to e^{+i\theta}10, \qquad 5_H \to e^{-2i\theta}5_H,$ $24_H \to e^{-10i\theta}24_H, \qquad 45_H \to e^{-2i\theta}45_H, \qquad 24 \to e^{+5i\theta}24.$

• Neutral fermion masses:

 $\mathcal{L} \supset h_1^i \,\bar{5}_i \, 24 \, 5_H + h_2^i \,\bar{5}_i \, 24 \, 45_H + \lambda \, \text{Tr}\{24^2 24_H\} + \text{h.c.}$



 $M_{\nu} = M_{\nu}^{I} + M_{\nu}^{III} + M_{\nu}^{cs}$

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 $\mathcal{L} \supset h_1 \overline{5} \, 24 \, 5_H + h_2 \, \overline{5} \, 24 \, 45_H + \lambda \, \text{Tr} \{ 24^2 24_H \} + \text{h.c.}$

• New fermion masses:

 M_{24} Tr{ 24^2 }

$$M_{\rho_0} = \frac{1}{3} M_{\rho_3}, \quad M_{\rho_8} = \frac{2}{3} M_{\rho_3}, \quad M_{\rho_{(3,2)}} = M_{\rho_{(\bar{3},2)}} = \frac{1}{6} M_{\rho_3}.$$

 $M_{24} \equiv M_{\rho_3}$

$$24_H \rightarrow \frac{v_{\Sigma}}{\sqrt{15}} \operatorname{diag}(-1, -1, -1, 3/2, 3/2) e^{ia(x)/v_{\Sigma}}$$

$$\mathcal{L} \supset \lambda \operatorname{Tr} \{ 24^2 \langle 24_H \rangle \} + \text{h.c.}$$

= $\frac{\lambda}{\sqrt{15}} v_{\Sigma} e^{ia(x)/v_{\Sigma}} \left(-\operatorname{Tr} \{ \rho_8 \rho_8 \} + \frac{1}{2} \operatorname{Tr} \{ \rho_{(\bar{3},2)} \rho_{(3,2)} \} + \frac{3}{2} \operatorname{Tr} \{ \rho_3 \rho_3 \} + \frac{1}{2} \rho_0^2 \right) + \text{h.c.}$
$$24_H \rightarrow \frac{v_{\Sigma}}{\sqrt{15}} \operatorname{diag}(-1, -1, -1, 3/2, 3/2) e^{ia(x)/v_{\Sigma}}$$

 $\begin{aligned} \mathcal{L} \supset \operatorname{Tr} \{ D_{\mu} \langle 24_{H} \rangle^{\dagger} D_{\mu} \langle 24_{H} \rangle \} &= \operatorname{Tr} \{ (ig_{\mathrm{GUT}} [A_{\mu}, \langle 24_{H} \rangle])^{\dagger} (ig_{\mathrm{GUT}} [A^{\mu}, \langle 24_{H} \rangle] \} \\ &= \underbrace{g_{\mathrm{GUT}}^{2} \frac{5}{6} v_{\Sigma}^{2}}_{M_{\mathrm{GUT}}^{2}} V_{\mu}^{\dagger} V^{\mu} \\ \mathcal{L} \supset \lambda \operatorname{Tr} \{ 24^{2} \langle 24_{H} \rangle \} + \text{h.c.} \\ &= \frac{\lambda}{\sqrt{15}} v_{\Sigma} e^{ia(x)/v_{\Sigma}} \left(-\operatorname{Tr} \{ \rho_{8} \rho_{8} \} + \frac{1}{2} \operatorname{Tr} \{ \rho_{(\bar{3},2)} \rho_{(3,2)} \} + \frac{3}{2} \operatorname{Tr} \{ \rho_{3} \rho_{3} \} + \frac{1}{2} \rho_{0}^{2} \right) + \text{h.c.} \end{aligned}$

• Same scalar field to break both theories $(SU(5) \otimes U(1)_{PQ})$

$$24_H \rightarrow \frac{v_{\Sigma}}{\sqrt{15}} \operatorname{diag}(-1, -1, -1, 3/2, 3/2) e^{ia(x)/v_{\Sigma}}$$





GUT scale window

Axion mass window

 $M_{\rm GUT} \simeq (10^{15.06} - 10^{15.74}) \,\,{\rm GeV}$

 $m_a \simeq (3 - 13) \times 10^{-9} \text{ eV}$



Res. I $B_{\text{max}} = 5 \ T, V_B = 1 \text{ m}^3$ Res. II $B_{\text{max}} = 20 T, V_B = 1 \text{ m}^3$ Res. III $B_{\text{max}} = 5 \ T, V_B = 100 \text{ m}^3$

SPOILER

The theory can be fully tested!



$$\frac{1}{\alpha_i(M_Z)} - \frac{1}{\alpha_i(\mu)} = \frac{1}{2\pi} b_i^{\rm SM} \, \ln\frac{\mu}{M_Z} + \frac{1}{2\pi} \sum_I b_{iI} \, \Theta(\mu - M_I) \, \ln\frac{\mu}{M_I}$$

$$\frac{1}{\alpha_i(M_Z)} - \frac{1}{\alpha_i(\mu)} = \frac{1}{2\pi} b_i^{\text{SM}} \operatorname{Ln} \frac{\mu}{M_Z} + \frac{1}{2\pi} \sum_I b_{iI} \Theta(\mu - M_I) \operatorname{Ln} \frac{\mu}{M_I}$$



$$\frac{1}{\alpha_i(M_Z)} - \frac{1}{\alpha_i(\mu)} = \frac{1}{2\pi} b_i^{\rm SM} \, \mathrm{Ln} \frac{\mu}{M_Z}$$







$$\frac{1}{\alpha_i(M_Z)} - \frac{1}{\alpha_i(\mu)} = \frac{1}{2\pi} b_i^{\text{SM}} \operatorname{Ln} \frac{\mu}{M_Z} + \frac{1}{2\pi} \sum_I b_{iI} \Theta(\mu - M_I) \operatorname{Ln} \frac{\mu}{M_I}$$

$$B_{ij}^I = (b_i^I - b_j^I) r_I$$

$$B_{ij}^{I} = (b_i^I - b_j^I) r_I$$

$$\frac{B_{23}}{B_{12}} = 0.718$$

$$\operatorname{Ln} \left(\frac{M_{GUT}}{M_Z}\right) = \frac{184.87}{B_{12}}$$

$\mathcal{L} \supset M_{24} \mathbb{T} \{24^2\} + \lambda \operatorname{Tr} \{24^2 24_H\} + \mathrm{h.c}$	•

	5_H		24		4	45_H							
B_{ij}	H_1	T	$ ho_8$	$ ho_3$	$\rho_{(3,2)} + \rho_{(\bar{3},2)}$	Φ_1	Φ_2	Φ_3	Φ_4	Φ_5	Φ_6	H_2	
B_{12}	$-\frac{1}{15}$	$\frac{1}{15}r_T$	0	$-\frac{4}{3}r_{ ho_3}$	$\frac{4}{3}r_{32}$	$-rac{8}{15}r_{\Phi_1}$	$\frac{2}{15}r_{\Phi_2}$	$-rac{9}{5}r_{\Phi_3}$	$rac{17}{15}r_{\Phi_4}$	$rac{1}{15}r_{\Phi_5}$	$rac{16}{15}r_{\Phi_6}$	$-rac{1}{15}r_{H_2}$	
B_{23}	$\frac{1}{6}$	$-\frac{1}{6}r_T$	$-2r_{\rho_8}$	$\frac{4}{3}r_{ ho_3}$	$\frac{2}{3}r_{32}$	$-\frac{2}{3}r_{\Phi_1}$	$-rac{5}{6}r_{\Phi_2}$	$rac{3}{2}r_{\Phi_3}$	$rac{1}{6}r_{\Phi_4}$	$-rac{1}{6}r_{\Phi_5}$	$-\frac{1}{6}r_{\Phi_6}$	$\frac{1}{6}r_{H_2}$	

$$\frac{1}{\alpha_i(M_Z)} - \frac{1}{\alpha_i(\mu)} = \frac{1}{2\pi} b_i^{\text{SM}} \operatorname{Ln} \frac{\mu}{M_Z} + \frac{1}{2\pi} \sum_I b_{iI} \Theta(\mu - M_I) \operatorname{Ln} \frac{\mu}{M_I}$$

$$B_{ij}^I = (b_i^I - b_j^I) r_I$$

$$B_{ij}^{IJ} = 0.718$$

$$\operatorname{Ln} \left(\frac{M_{GUT}}{M_Z}\right) = \frac{184.87}{B_{12}}$$
Help towards unification
$$\mathcal{L} \supset M_{24} \operatorname{Te} 24^2 + \lambda \operatorname{Tr} \{24^2 24_H\} + \text{h.c.}$$
Help towards unification
$$\frac{45_H}{1 - \frac{1}{15} r_T} = \frac{1}{0} - \frac{4}{3} r_{\rho_3} - \frac{4}{3} r_{\rho_3}} - \frac{8}{3} r_{\rho_3} - \frac{2}{3} r_{\rho_1} - \frac{5}{6} r_{\rho_2}} - \frac{3}{3} r_{\rho_3} - \frac{4}{5} r_{\rho_5}}{\frac{3}{2} r_{\rho_3}} - \frac{4}{3} r_{\rho_3}}$$

 B_{ij}

 B_{12}

 B_{23}



Allowed by unification constraints





Miralles, Pich $M_{\Phi_1} \ge 1 \text{ TeV} \Rightarrow M_{\text{GUT}} \le 10^{15.87} \text{ GeV}$







$$\Gamma(p \to M\bar{\ell}) = \frac{m_p}{32\pi} \left(1 - \left(\frac{m_p}{m_M}\right)^2 \right)^2 \left| \sum_i \langle M | \mathcal{O}_i^{X,Y} | p \rangle \right|^2$$
• Proton decay mediators:

$$\mathcal{L}_p \supset i\bar{5}\gamma^\mu D_\mu \bar{5} + \frac{i}{2} \operatorname{Tr}\{\overline{10}\gamma^\mu D_\mu 10\}$$

$$- M_{GUT}^2 \left(\operatorname{Tr}\{X_\mu^{cT} C X^\mu + Y_\mu^{cT} C Y^\mu\} \right) + \text{h.c.}$$

$$V_\mu \sim (3, 2, -5/6)$$
• Integrate X and Y fields out
$$V_\mu^c = \begin{pmatrix} X_\mu \\ Y_\mu \end{pmatrix} \xrightarrow{\to} \mathcal{Q}_X = 4/3$$

$$\to \mathcal{Q}_Y = 1/3$$

|2

$$\Gamma(p \to \pi^0 e^+) = \frac{\pi m_p}{2} \left(\frac{\alpha_{\rm GUT}}{M_{\rm GUT}^2}\right)^2 A_R^2 V_e |\langle \pi^0 | (ud)_R u_L | p \rangle|^2$$







Window for GUT Scale



$$M_{\rm GUT} = [10^{15.06} - 10^{15.74}] \,\,{\rm GeV}$$

Axion couplings

$$\mathcal{L} \supset \frac{g_S^2}{32\pi^2} \frac{a}{v_{\Sigma}} NG\tilde{G} \equiv \frac{g_S^2}{32\pi^2} \frac{a}{f_a} G\tilde{G}$$

$$N = \sum_{\substack{\text{colored}\\\text{chiral}\\\text{fields}}} C_{\Psi_L^i} T_D[R_{\Psi_L^i}] \times \text{mult}[\Psi_L^i] = \left(\underbrace{3 \times 1}_{\rho_8} + \underbrace{\frac{1}{2} \times 2}_{\rho_{(3,2)}} + \underbrace{\frac{1}{2} \times 2}_{\rho_{(\bar{3},2)}} \right) = 5$$



$$f_a \equiv \frac{N}{v_{\Sigma}} = \sqrt{\frac{6}{5\pi\alpha_{\rm GUT}}} \frac{M_{\rm GUT}}{10}$$



Axion-EDM coupling

$$\mathcal{L} \supset -\frac{i}{2} g_{aD} a (\overline{\Psi_N} \sigma^{\mu\nu} \gamma_5 \Psi_N) F_{\mu\nu}$$
$$d_n = g_{aD} a \approx 2.4 \times 10^{-16} \frac{a}{f_a} e \cdot \text{cm}$$



$$g_{a\gamma\gamma} = \frac{\alpha_{\rm EM}}{2\pi f_a} \left(\frac{E}{N} - 1.92(4)\right)$$

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$$E = \sum_{\substack{\text{charged}\\\text{chiral}\\\text{fields}}} C_{\Psi_L^I} \mathcal{Q}(\Psi_L^I)^2 \times \text{mult}[\Psi_L^I]$$

$$g_{a\gamma\gamma} = \frac{\alpha_{\rm EM}}{2\pi f_a} \left(\frac{E}{N} - 1.92(4)\right) = \frac{\alpha_{\rm EM}}{2\pi f_a} \left(\frac{8}{3} - 1.92(4)\right)$$

$$E = \sum_{\substack{\text{charged} \\ \text{chiral} \\ \text{fields}}} C_{\Psi_L^I} \mathcal{Q}(\Psi_L^I)^2 \times \text{mult}[\Psi_L^I] \xrightarrow{\text{Complete SU(5)}}_{\text{representation}}$$

$$\rho_{(3,2)} \qquad \rho_{(3,2)} \qquad \rho_{($$

$$g_{a\gamma\gamma} = \frac{\alpha_{\rm EM}}{2\pi f_a} \left(\frac{E}{N} - 1.92(4)\right) = \frac{\alpha_{\rm EM}}{2\pi f_a} \left(\frac{8}{3} - 1.92(4)\right)$$



• Most economical renormalizable SU(5) where the PQ mechanism can be realized: The Adjoint SU(5): 45н, 24

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- The theory could be fully tested by the **ABRACADABRA** and the **CASPEr-Electric** experiments.
- Appealing theory!
Thank you for your attention!!

Alpha GUT

 $B_{23}, B_{12} = 0$ RGE only sensitive to splitting in 24 BUT $b_i \neq 0 \Rightarrow \alpha_{\text{GUT}} = \alpha_{\text{GUT}}(M_{\rho_3})$

Range of M_{ρ_3} ?

• Lower bound

• Upper bound



$$M_{\nu} \simeq \frac{h_1^2 v_0^2}{M_{\rho_3}}$$

Perturbativity Yukawa: $\Rightarrow M_{\rho_3} \lesssim 10^{15} \text{ GeV}$

 10^{33}

 $A_c Y_A A \equiv Y_A^{\rm diag}$



15.4

15.2

15.0

15.6

15.8

Alpha GUT

 B_{23}, B_{12} only sensitive to splitting in **24** BUT $b_i = f(M_{\rho_3}) \Rightarrow \alpha_{\text{GUT}} = f(M_{\rho_3})$



 $10^{4.51} \text{ GeV} \lesssim M_{\rho_3} \lesssim 10^{15} \text{ GeV}$