# Maximizing Direct Detection with HYPER Dark Matter

Gilly Elor MITP, JGU

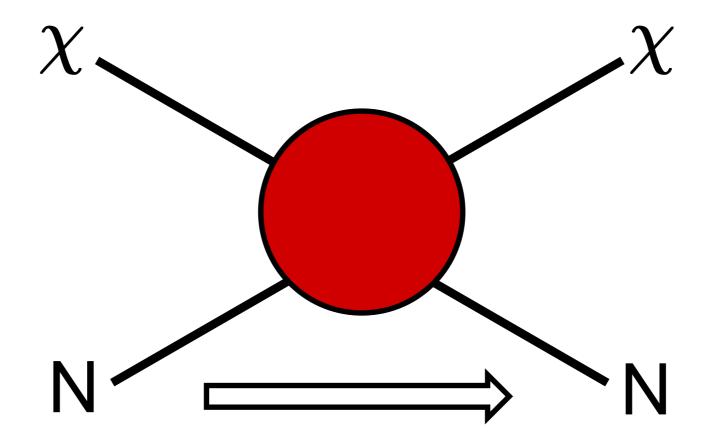
Based on:

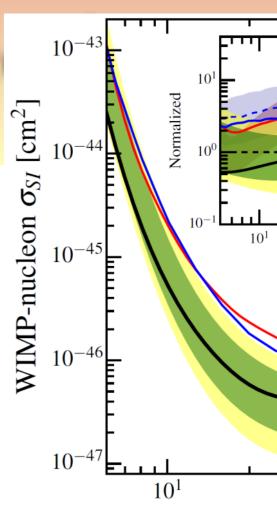
soon to be posted work with Robert McGehee and Aaron Pierce

#### Dark Matter Direct Detection

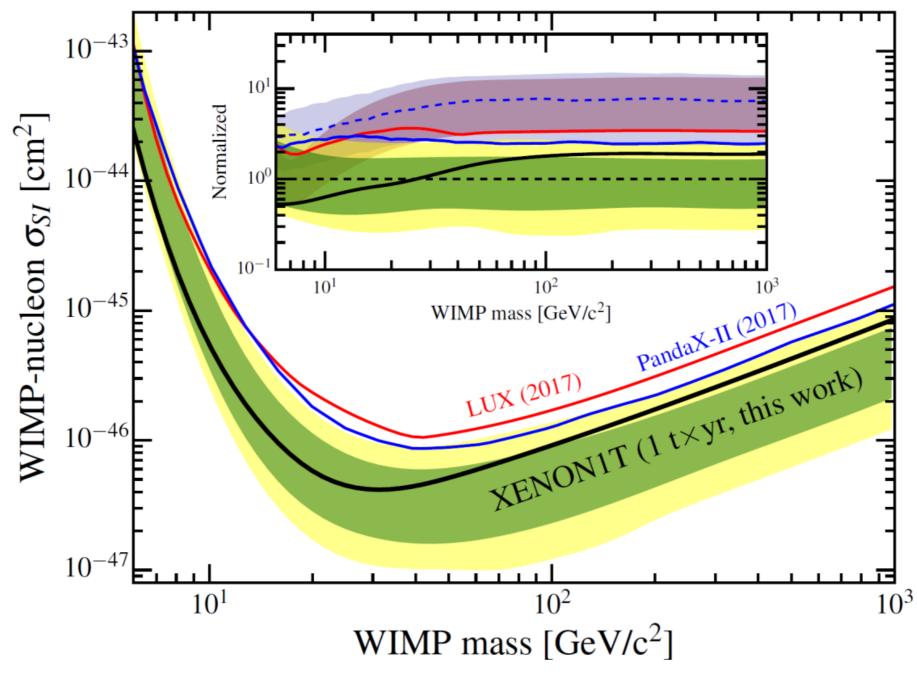


### Dark Matter Direct Dete





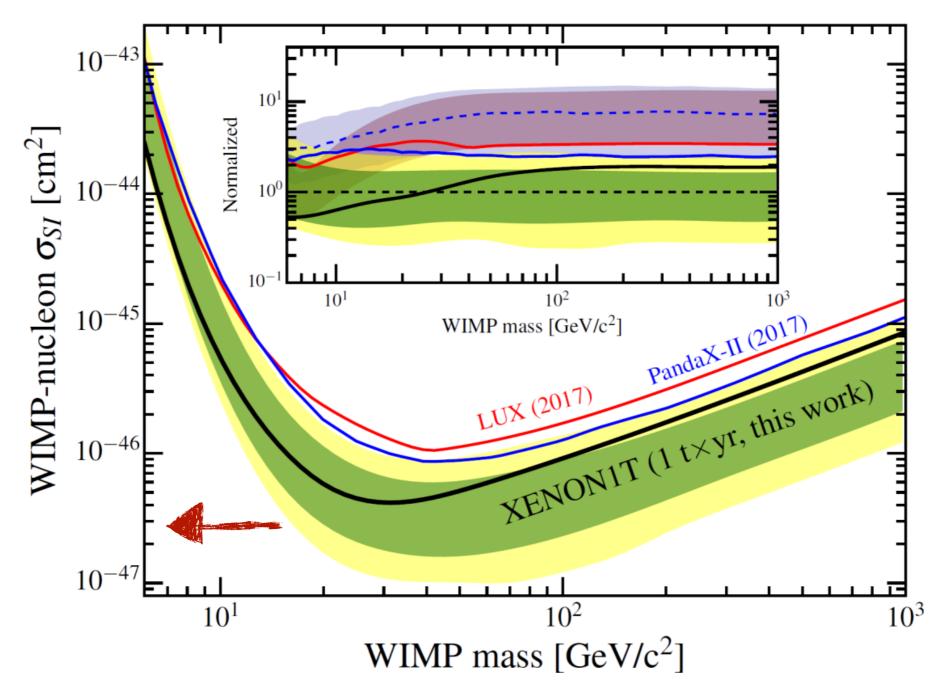
#### WIMP Dark Matter



XENON Collaboration PRL 121 (2018) no. 11, 111302

G. Elor

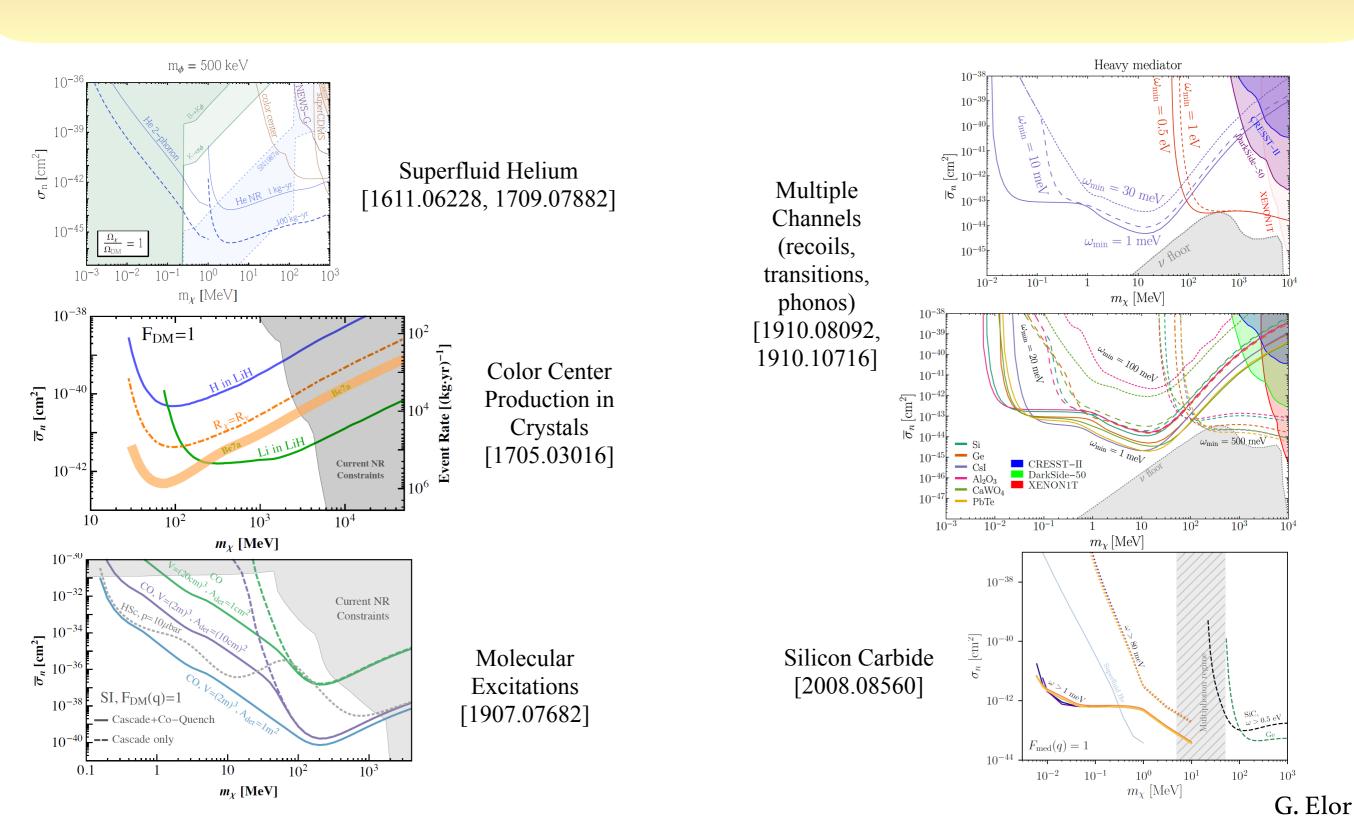
### Sub-GeV Dark Matter



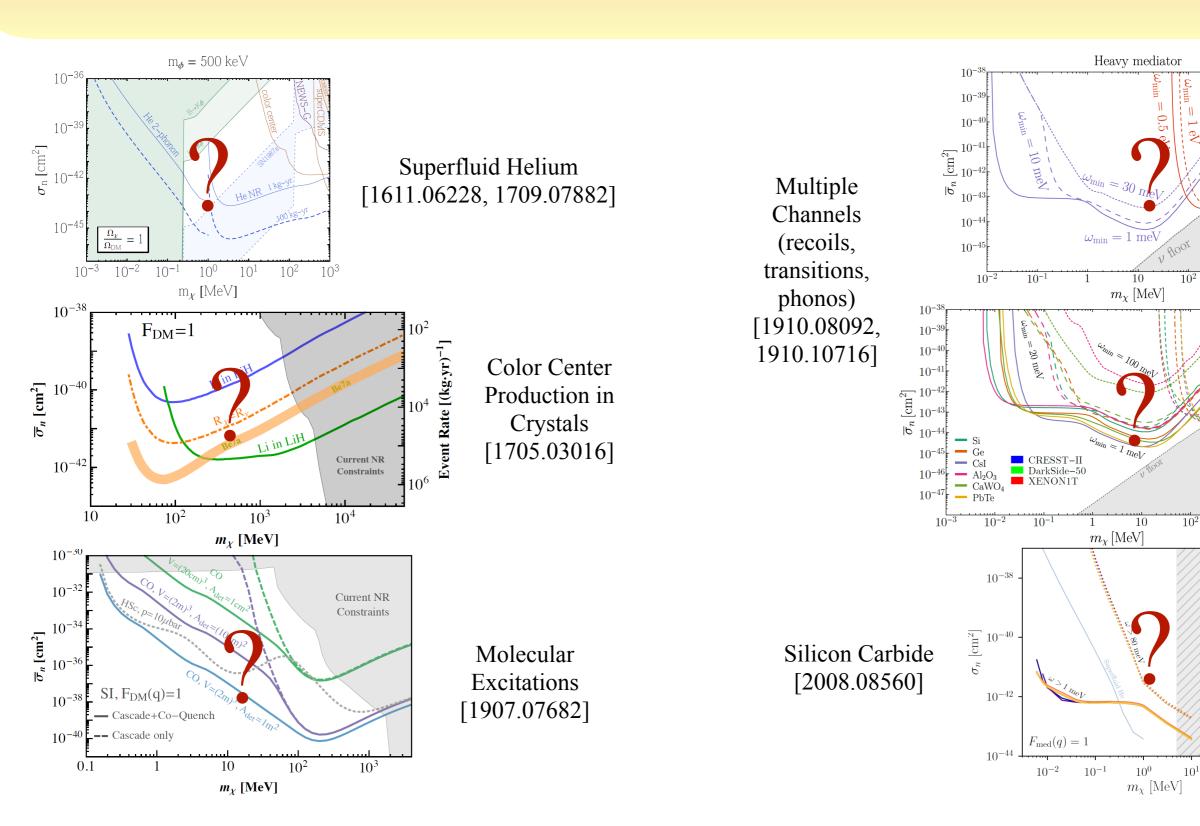
XENON Collaboration PRL 121 (2018) no. 11, 111302

G. Elor

### Sub-GeV Dark Matter



#### Where is the Dark Matter?



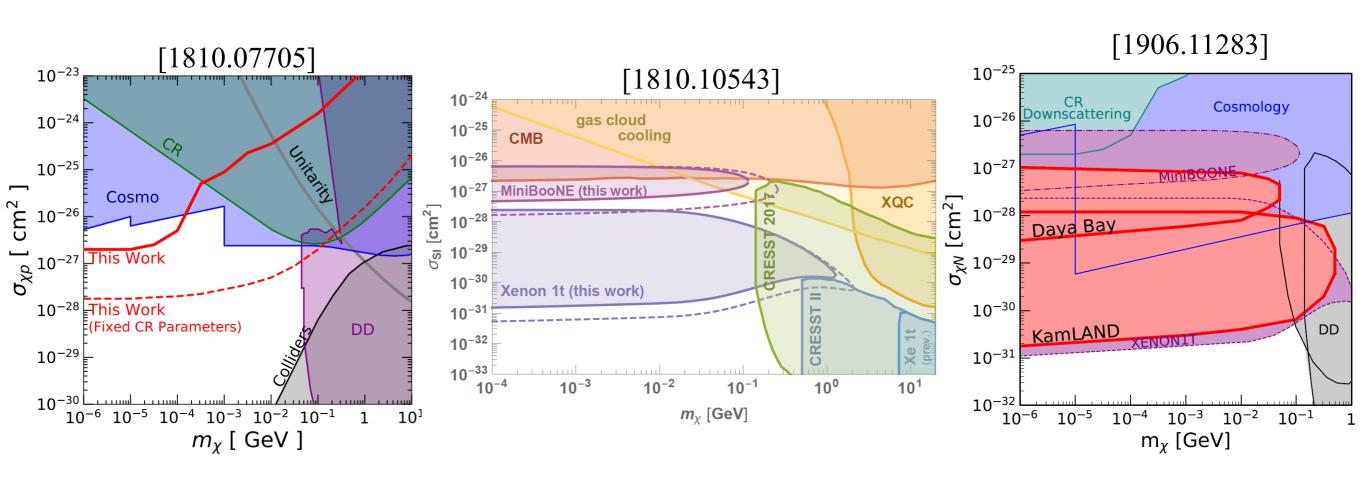
 $\omega_{\rm min} = 500 \ {\rm meV}$ 

 $10^{3}$ 

 $10^{3}$ 

G. Elor

# Bounds from Cosmic Ray Upscattering



Is such a large cross section even feasible?
Is there a maximum cross section for DM nucleon scattering given present day constraints?

## Maximizing Direct Detection

There exists a maximum cross section  $\sigma_{\chi n}^{\rm max}$ . To design experiments targeting larger cross sections is not motivated.

### A Hadrophilic Scalar Mediator

$$\left[ egin{array}{lll} {\cal L} \ \supset \ -m_{\chi}ar{\chi}\chi - y_n\phiar{n}n - y_{\chi}\phiar{\chi}\chi \end{array} 
ight]$$

UV Model: new vector-like quarks at the TeV scale

S. Knapen, T. Lin, K. Zurek [1790.07882]

$$\mathcal{L} \supset \lambda \phi \bar{\psi} \psi \longrightarrow \frac{\alpha_s}{\Lambda} \phi G^{\mu\nu} G_{\mu\nu} \qquad \frac{1}{\Lambda} = \frac{\lambda}{M_{\psi}} \leftrightarrow \frac{y_n}{m_n}$$

### A Hadrophilic Scalar Mediator

$$\left[ egin{array}{lll} {\cal L} \ \supset \ -m_{\chi}ar{\chi}\chi - y_n\phiar{n}n - y_{\chi}\phiar{\chi}\chi \end{array} 
ight]$$

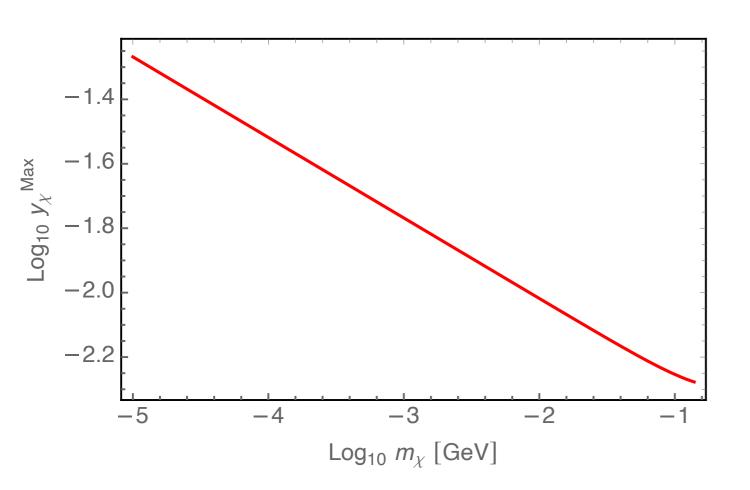
UV Model: new vector-like quarks at the TeV scale

S. Knapen, T. Lin, K. Zurek [1790.07882]

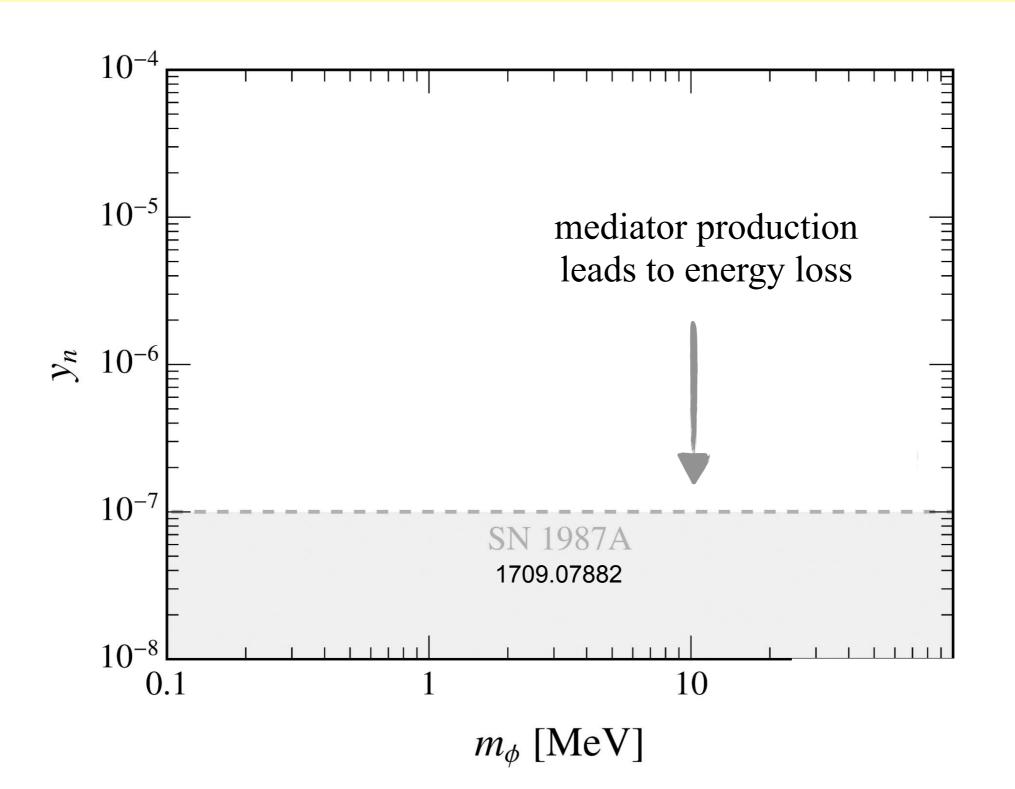
$$\mathcal{L} \supset^{\lambda \phi \bar{\psi} \psi} m_{\chi} \bar{\chi} \bar{\chi} \bar{\chi}^{\alpha_s} \phi G^{\mu\nu} G_{\mu} \phi \bar{n} \bar{n} = \bar{\lambda}_{M} \chi \phi \bar{\chi}^{n} \bar{\chi}$$

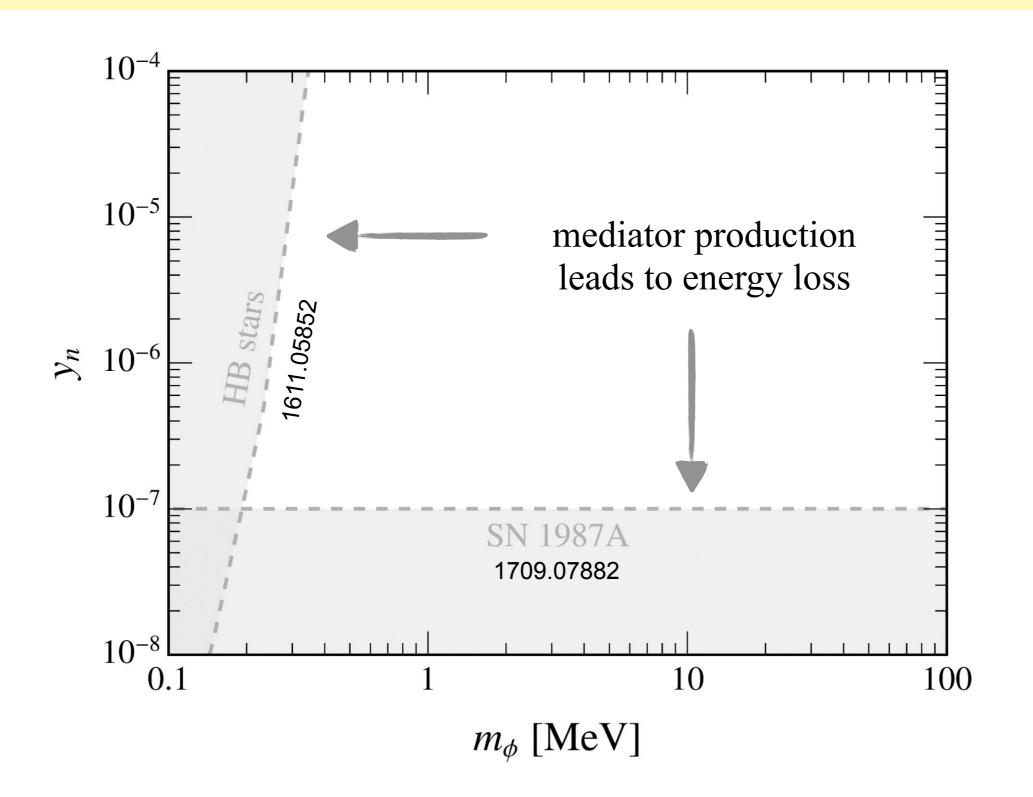
$$\sigma_{\chi n}^{ ext{max}} \equiv rac{\left(y_n^{ ext{max}}y_\chi^{ ext{max}}
ight)^2}{4\pi} rac{\mu_{\chi n}^2}{\left[\left(m_\phi^{ ext{min}}
ight)^2 + v_{ ext{DM}}^2 m_\chi^2
ight]^2}$$

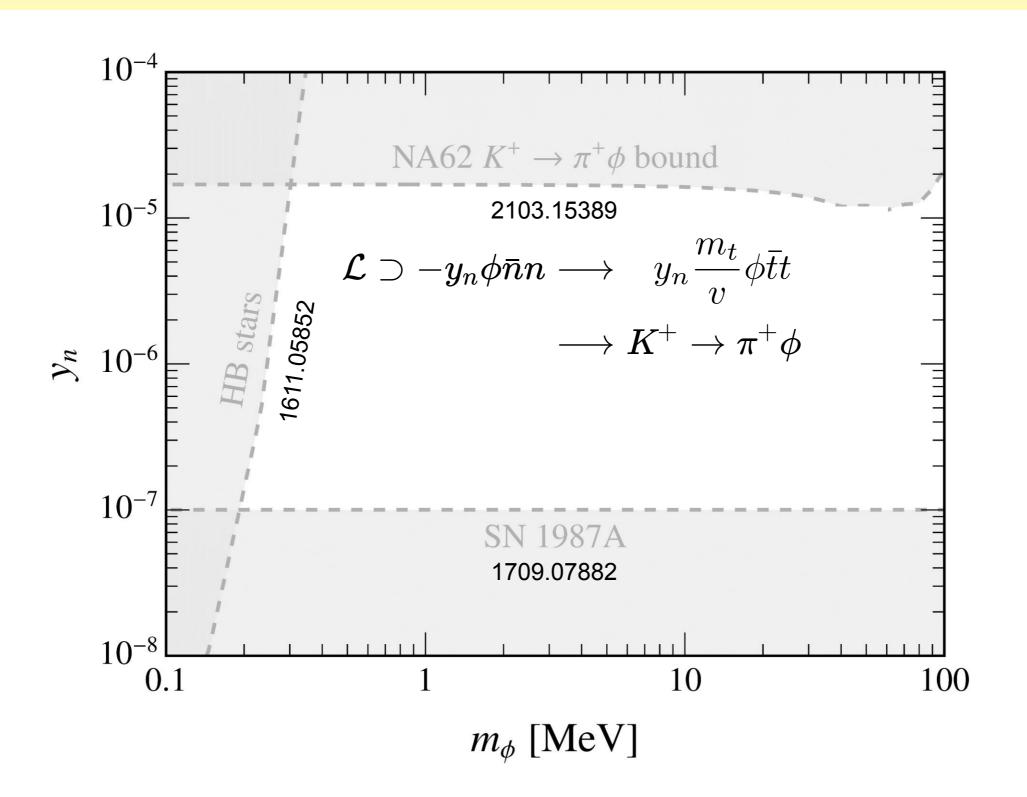
Dark Matter Self Interactions: 
$$\sigma_T^{\rm born} \approx \frac{8\pi\alpha_\chi^2}{m_\chi^2 v^4} \left[ \log(1+R^2) - R^2/(1+R^2) \right]$$
  $R \equiv m_\chi v/m_\phi$ 

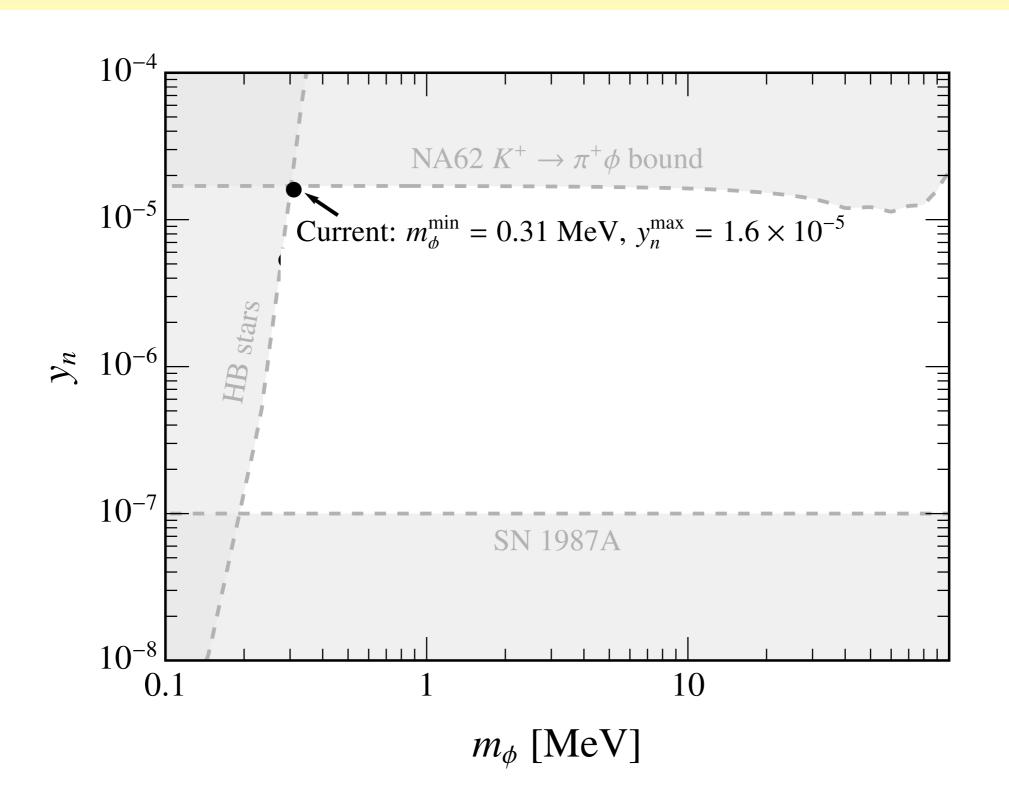


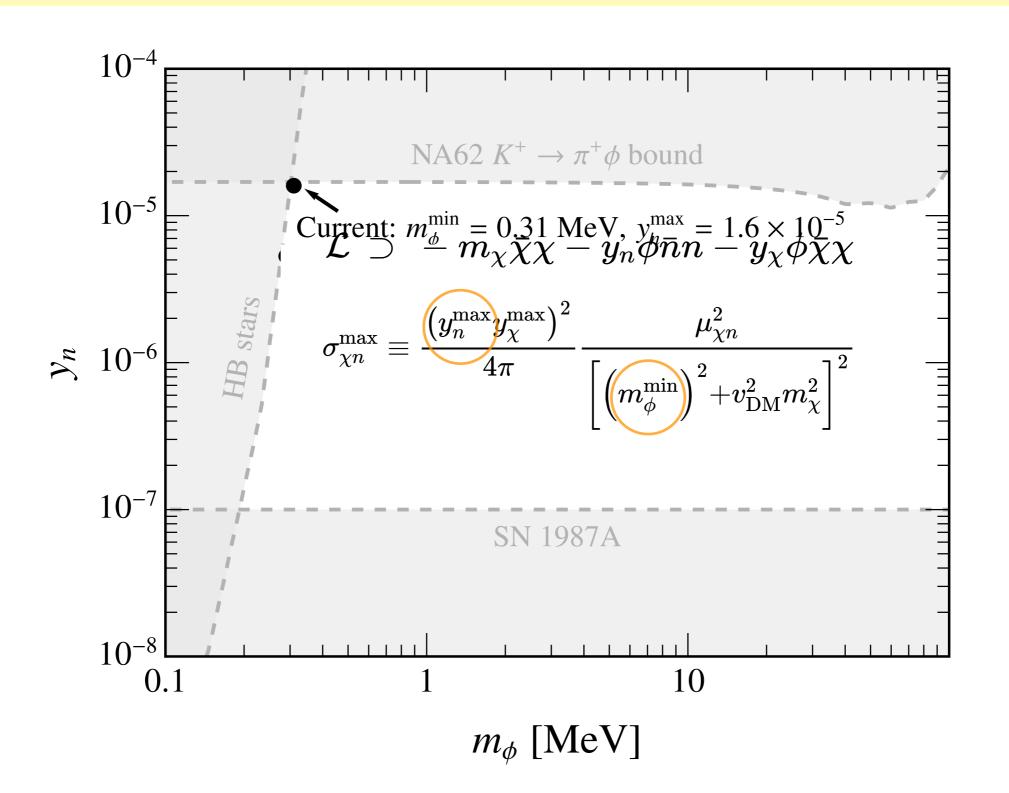
$$egin{aligned} \sigma_{\chi\chi}/m_\chi &\lesssim 1~\mathrm{cm}^2/\mathrm{g} \ \mathrm{at}~v \sim &10^{-3} \end{aligned} \ \mathcal{L}~\supset~-m_\chiar\chi\chi - y_n\phiar n n - y_\chi\phiar\chi\chi \ \sigma_{\chi n}^\mathrm{max} &\equiv rac{\left(y_n^\mathrm{max}y_\chi^\mathrm{max}
ight)^2}{4\pi} rac{\mu_{\chi n}^2}{\left[\left(m_\phi^\mathrm{min}
ight)^2 + v_\mathrm{DM}^2 m_\chi^2
ight]^2} \end{aligned}$$

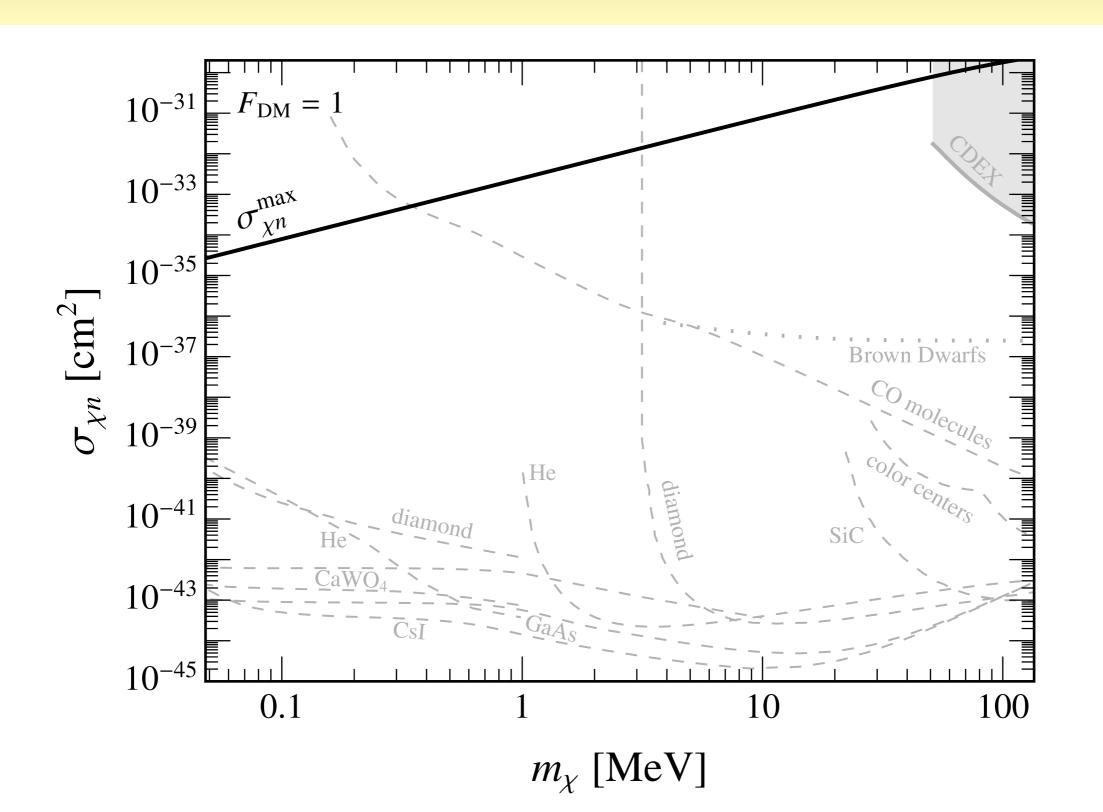




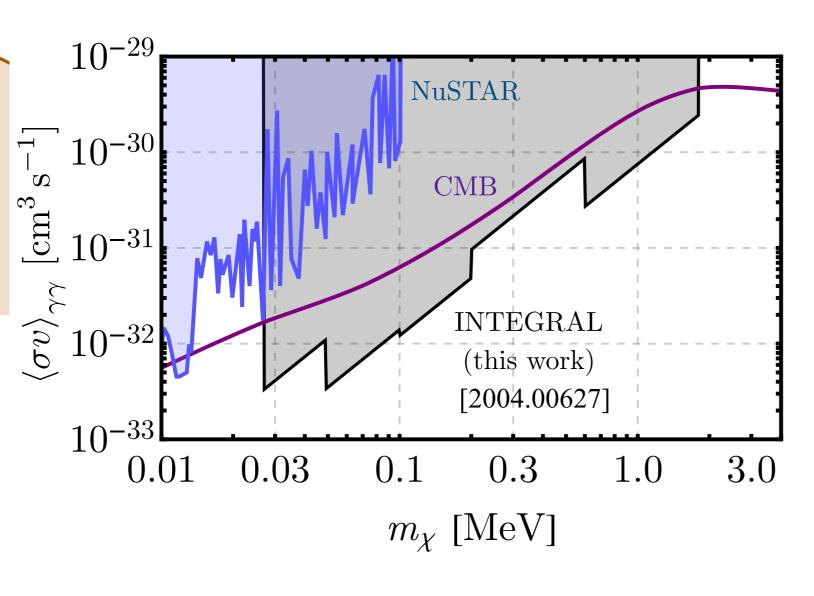


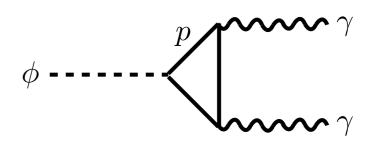






### Indirect Detection $\chi \bar{\chi} \rightarrow \gamma \gamma$

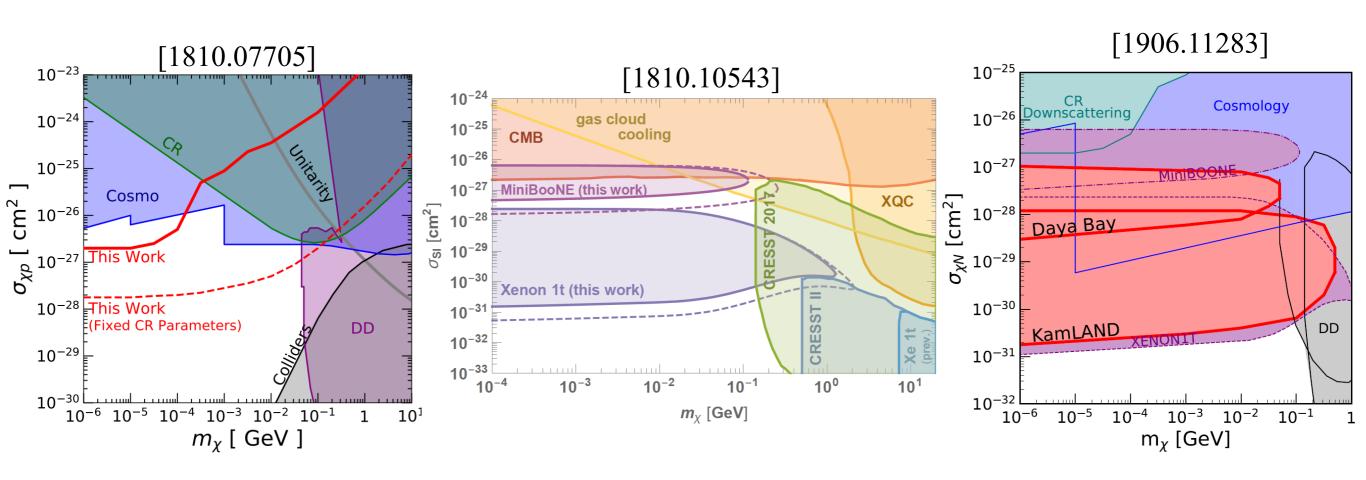




$$\mathcal{L} \supset \frac{\alpha y_n}{6\pi m_p} \phi F_{\mu\nu} F^{\mu\nu}$$

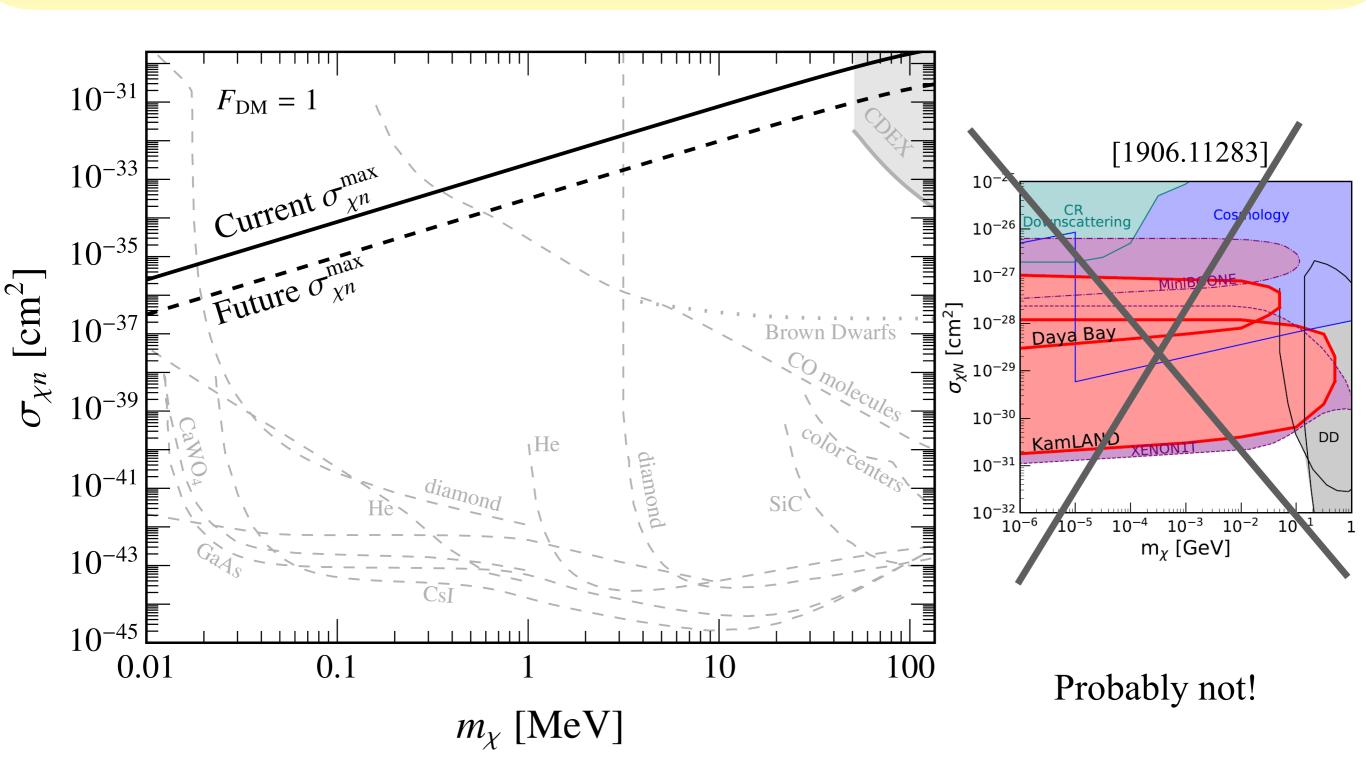
$$\sigma v_{\rm ann} = \frac{1}{32\pi} \left( \frac{2\alpha y_n^{\rm max} y_\chi^{\rm max}}{3\pi m_p} \right)^2 \frac{s \left( s - 4m_\chi^2 \right)}{\left( s - (m_\phi^{\rm min})^2 \right)^2} \sim 10^{-44} \rm cm^3 s^{-1}$$

# Bounds from Cosmic Ray Upscattering



Is such a large cross section even feasible?

# Bounds from Cosmic Ray Upscattering



# Robustness of $\sigma_{n\chi}^{\text{max}}$ ?

Is  $\sigma_{n\chi}^{\text{max}}$  for the Hydrophilic scalar model the  $\sigma_{n\chi}^{\text{max}}$ ?

$$\mathcal{L} \supset \lambda \phi \bar{\psi} \psi$$
  $\longrightarrow$   $\frac{\alpha_s}{\Lambda} \phi G^{\mu\nu} G_{\mu\nu}$   $\longrightarrow$   $\mathcal{L} \supset -m_{\chi} \bar{\chi} \chi - y_n \phi \bar{n} n - y_{\chi} \phi \bar{\chi} \chi$ 

- Hadrophilic scalar with different UV completion e.g. mediator couples directly to quarks  $\longrightarrow$  Meson bounds are more constraining  $\longrightarrow$  smaller  $\sigma_{n\chi}^{\text{max}}$ .
- Visibly decaying dark photon? Beam dump and collider constraints make  $\sigma_{n\chi}^{\rm max}$  smaller.

# Achieving $\sigma_{n\chi}^{\text{max}}$ ?

Is there a sub-GeV dark matter candidate that:

1) may be detected at proposed experiments?

$$\mathcal{L} \supset -m_{\chi}^{2} \chi^{\text{may}} \eta_{n} \psi_{n}^{\text{hay}} \eta_{n}^{\text{uch}} y_{\chi}^{\text{large}} \text{ cross section?}$$

$$\sigma_{\chi n}^{ ext{max}} \equiv rac{\left(y_n^{ ext{max}}y_\chi^{ ext{max}}
ight)^2}{4\pi} rac{\mu_{\chi n}^2}{\left[\left(m_{\star}^{ ext{min}}
ight)^2 + v_{ ext{DM}}^2 m_{\star}^2
ight]^2} \qquad ext{and} \qquad \Omega_\chi h^2 = 0.11$$

# Achieving $\sigma_{n\chi}^{\text{max}}$ ?

Is there a sub-GeV dark matter candidate that:

1) may be detected at proposed experiments?

$$\mathcal{L} \supset -m_{\chi}^{2} \chi^{\text{may}} \eta_{n} \psi_{n}^{\text{may}} \eta_{n}^{\text{may}} \eta_{n}^{\text{such}} \eta_{\chi}^{\text{large}} \text{ cross section?}$$

$$\sigma_{\chi n}^{
m max} \equiv rac{\left(y_n^{
m max}y_\chi^{
m max}
ight)^2}{4\pi} rac{\mu_{\chi n}^2}{\left[\left(m_\phi^{
m min}
ight)^2 + v_{
m DM}^2 m_\chi^2
ight]^2} \qquad ext{and} \qquad \Omega_\chi h^2 = 0.11$$

- Large couplings could over-annihilate in the early Universe:  $\chi \bar{\chi} \to \phi \phi$ , leading to  $\Omega_{\gamma} h^2 < 0.1$
- BBN and CMB constrain sub-MeV dark matter with large cross sections.
- Dark matter (and mediators) with MeV mass and large interactions could thermalize the bath and lead to  $N_{\text{eff}}$  constraints.

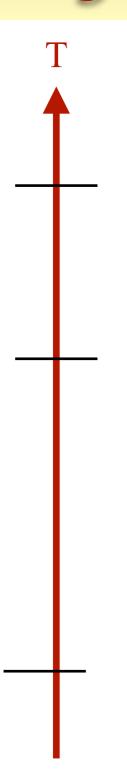
### Maximizing Direct Detection

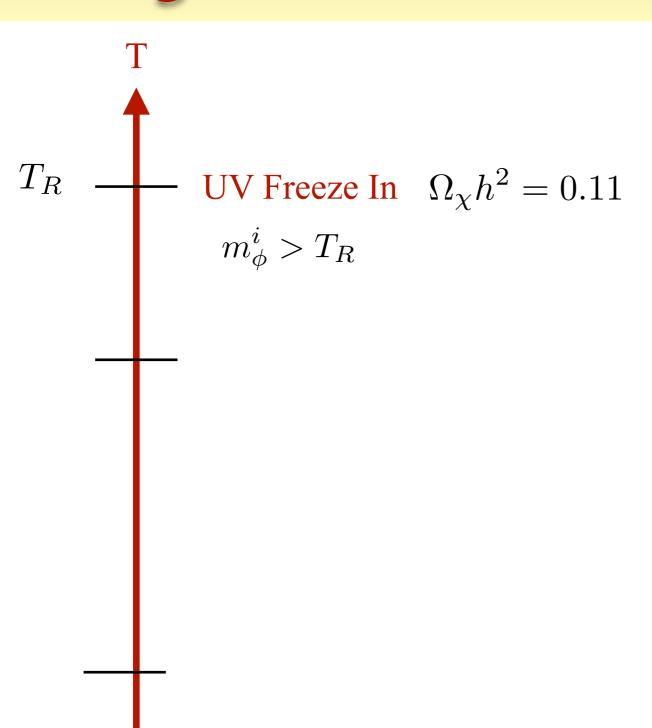
There exists a maximum cross section  $\sigma_{\chi n}^{\rm max}$ .

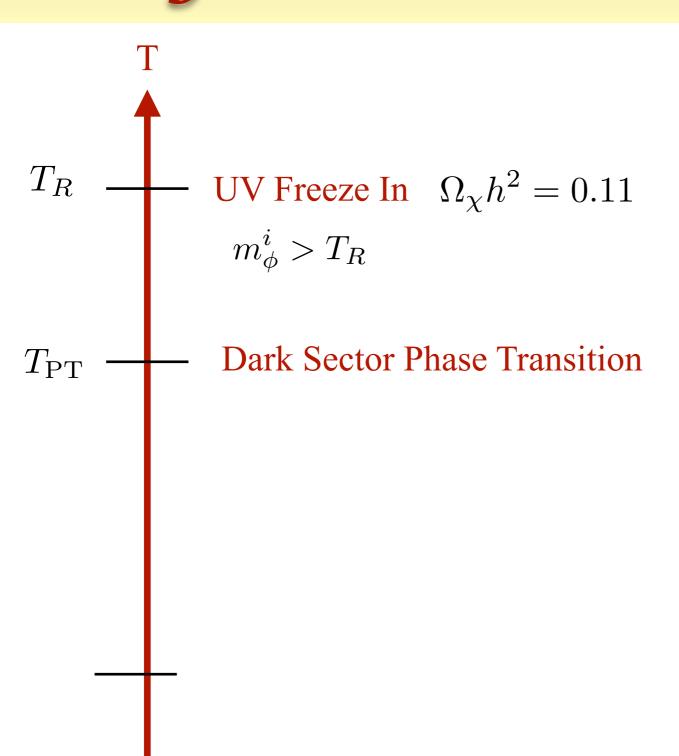
To design experiments targeting larger cross sections is not motivated.

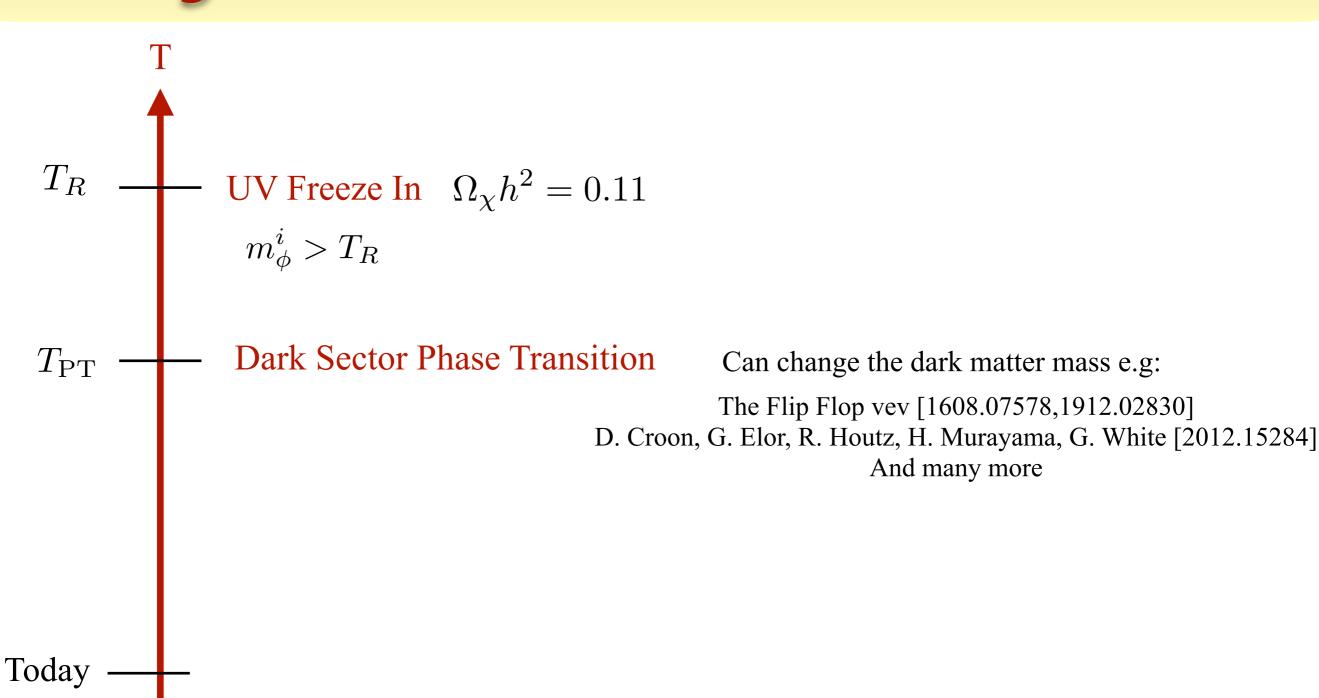
### with HYPER Dark Matter

There exists a model of dark matter that can achieve  $\sigma_{\chi n}^{\max}$ , and generally lives in a parameter space upcoming experiments will target.

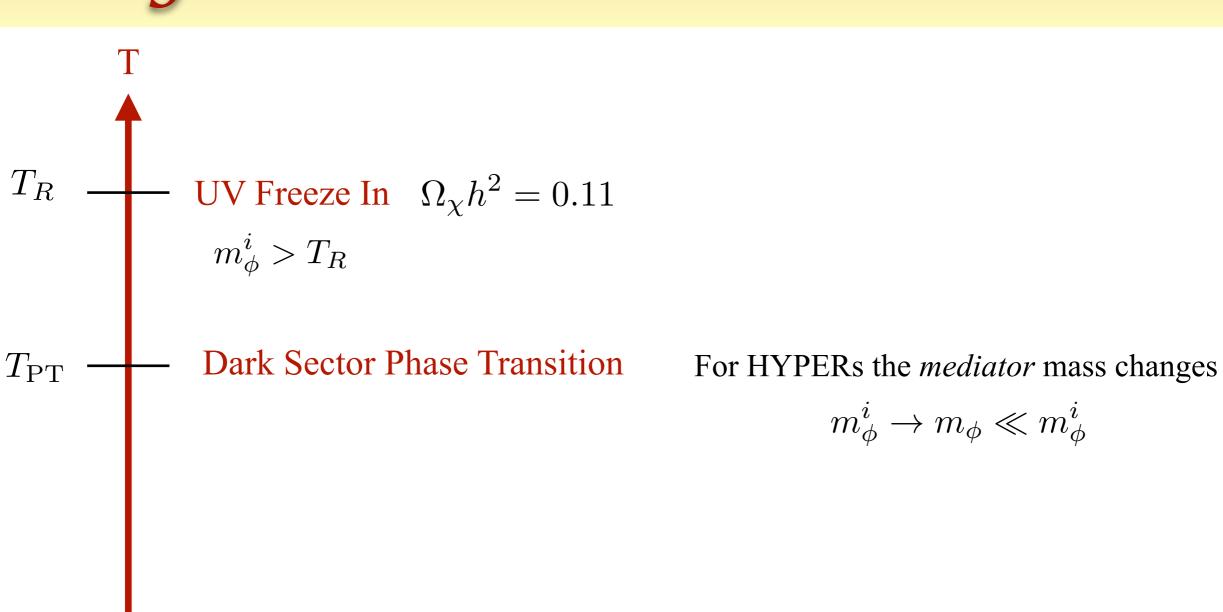




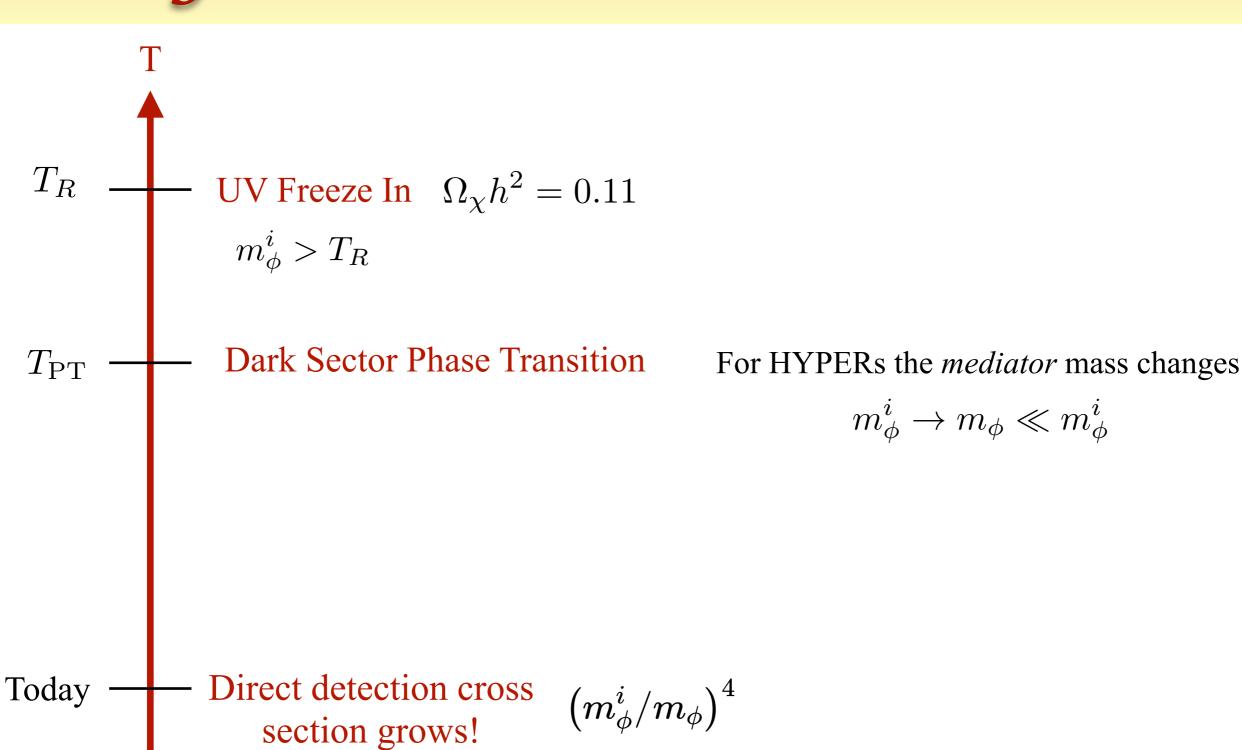


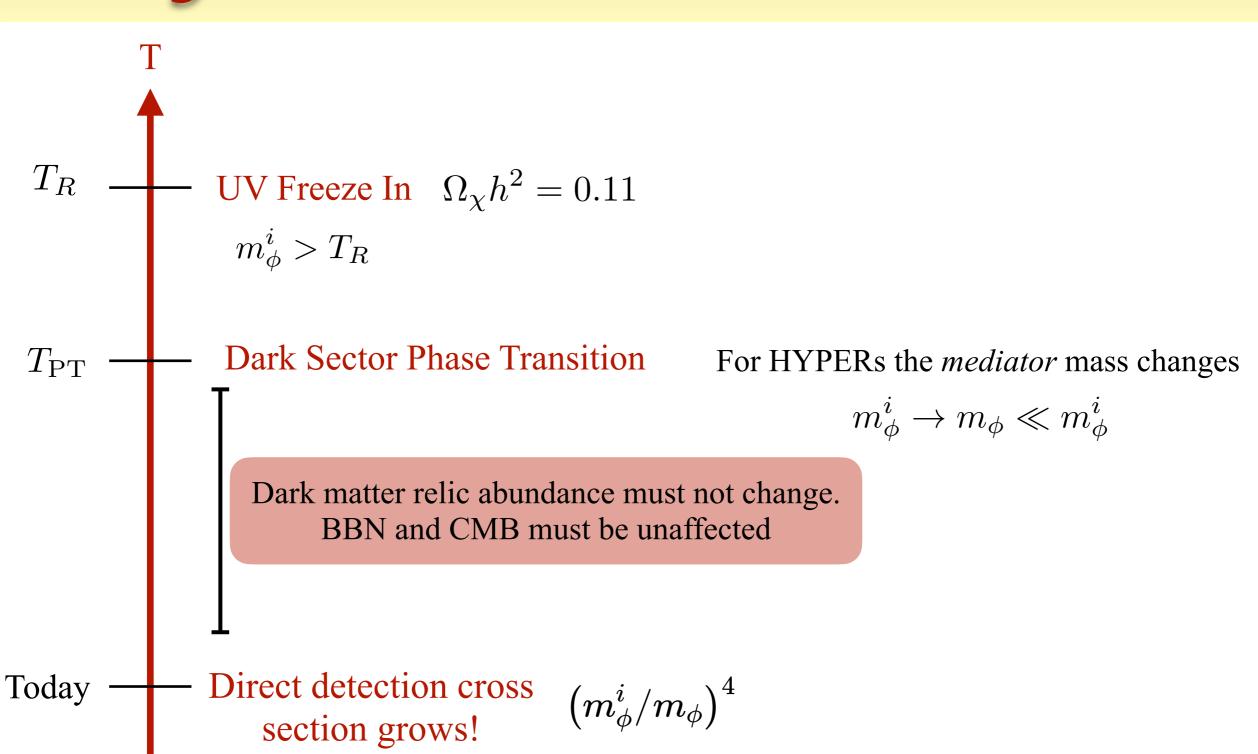


### HighlY interactive ParticlE Relics



Today





#### UV Freeze-In

At high scales, integrating out initially heavy mediator and heavy vector like quarks:

$$\mathcal{L} \supset \lambda \phi \bar{\psi} \psi \longrightarrow \frac{\alpha_s}{\Lambda} \phi G^{\mu\nu} G_{\mu\nu} \longrightarrow \left( \frac{\alpha_s y_{\chi} y_n}{2.6 \, m_n (m_{\phi}^i)^2} \chi \bar{\chi} G^{a, \, \mu\nu} G^a_{\mu\nu} \right)$$

#### UV Freeze-In

At high scales, integrating out initially heavy mediator and heavy vector like quarks:

$$\mathcal{L} \supset \lambda \phi \bar{\psi} \psi \longrightarrow \frac{\alpha_s}{\Lambda} \phi G^{\mu\nu} G_{\mu\nu} \longrightarrow \left( \frac{\alpha_s y_{\chi} y_n}{2.6 \, m_n (m_{\phi}^i)^2} \chi \bar{\chi} G^{a, \, \mu\nu} G^a_{\mu\nu} \right)$$

UV Freeze In: 
$$T_R < \text{Min} \left[ m_{\phi}^i / 20, m_{\psi} / 20 \right]$$
  $Y_{\chi} \simeq 4 \times 10^{-5} \left( \frac{y_n y_{\chi} \alpha_s}{m_n (m_{\phi}^i)^2} \right)^2 \frac{M_{\rm Pl} T_R^5}{g_{s,*} \sqrt{g_*}}$ 

We can adjust  $T_{\rm R}$  and  $m_{\phi}^i$  to yield the correct relic abundance.

Both have no impact on  $\sigma_{\chi n}^{\rm max}$ .

#### UV Freeze-In

At high scales, integrating out initially heavy mediator and heavy vector like quarks:

$$\mathcal{L} \supset \lambda \phi \bar{\psi} \psi \longrightarrow \frac{\alpha_s}{\Lambda} \phi G^{\mu\nu} G_{\mu\nu} \longrightarrow \left( \frac{\alpha_s y_{\chi} y_n}{2.6 \, m_n (m_{\phi}^i)^2} \chi \bar{\chi} G^{a, \, \mu\nu} G^a_{\mu\nu} \right)$$

UV Freeze In: 
$$T_R < \text{Min} \left[ m_{\phi}^i / 20, m_{\psi} / 20 \right]$$
  $Y_{\chi} \simeq 4 \times 10^{-5} \left( \frac{y_n y_{\chi} \alpha_s}{m_n (m_{\phi}^i)^2} \right)^2 \frac{M_{\rm Pl} T_R^5}{g_{s,*} \sqrt{g_*}}$ 

We can adjust  $T_R$  and  $m_{\phi}^i$  to yield the correct relic abundance. Both have no impact on  $\sigma_{\chi n}^{\text{max}}$ .

$$\frac{\lambda}{M_{\psi}} \leftrightarrow \frac{y_n}{m_n} \longrightarrow y_n^{\text{max}} = 1.5 \times 10^{-5} \quad \leftrightarrow \quad m_{\psi}^{\text{max}} \sim 40 \,\text{TeV}$$

$$T_R \lesssim 2 \; {
m TeV}$$

### Dark Sector Phase Transition

$$m_{\phi}^{i} \rightarrow m_{\phi} \ll m_{\phi}^{i}$$

- Dark matter relic abundance must not change.
- BBN and CMB must be unaffected

e.g. after the phase transition we must forbid or suppress processes such as

$$m_{\phi}^{i} \rightarrow m_{\phi} \ll m_{\phi}^{i}$$

- Dark matter relic abundance must not change.
- BBN and CMB must be unaffected

e.g. after the phase transition we must forbid or suppress processes such as

$$\bar{\chi}\chi \to \text{hadrons}$$

hadrons 
$$\rightarrow \bar{\chi}\chi$$

$$\bar{\chi}\chi \to \phi\phi$$

• • • •

$$m_{\phi}^{i} \rightarrow m_{\phi} \ll m_{\phi}^{i}$$

- Dark matter relic abundance must not change.
- BBN and CMB must be unaffected

e.g. after the phase transition we must forbid or suppress processes such as

$$\bar{\chi}\chi \rightarrow \text{hadrons}$$

$$\mathcal{O}(10 \text{ keV}) \lesssim m_{\chi} < m_{\pi^0}$$

hadrons 
$$\rightarrow \bar{\chi}\chi$$

$$\bar{\chi}\chi \to \phi\phi$$

• • •

$$m_{\phi}^{i} \rightarrow m_{\phi} \ll m_{\phi}^{i}$$

- Dark matter relic abundance must not change.
- BBN and CMB must be unaffected

e.g. after the phase transition we must forbid or suppress processes such as

$$\bar{\chi}\chi \to {
m hadrons}$$
  $\mathcal{O}(10~{
m keV}) \lesssim m_{\chi} < m_{\pi^0}$  hadrons  $\to \bar{\chi}\chi$  1 MeV  $\lesssim T_{\rm PT} \lesssim m_{\pi^0}$ 

• • •

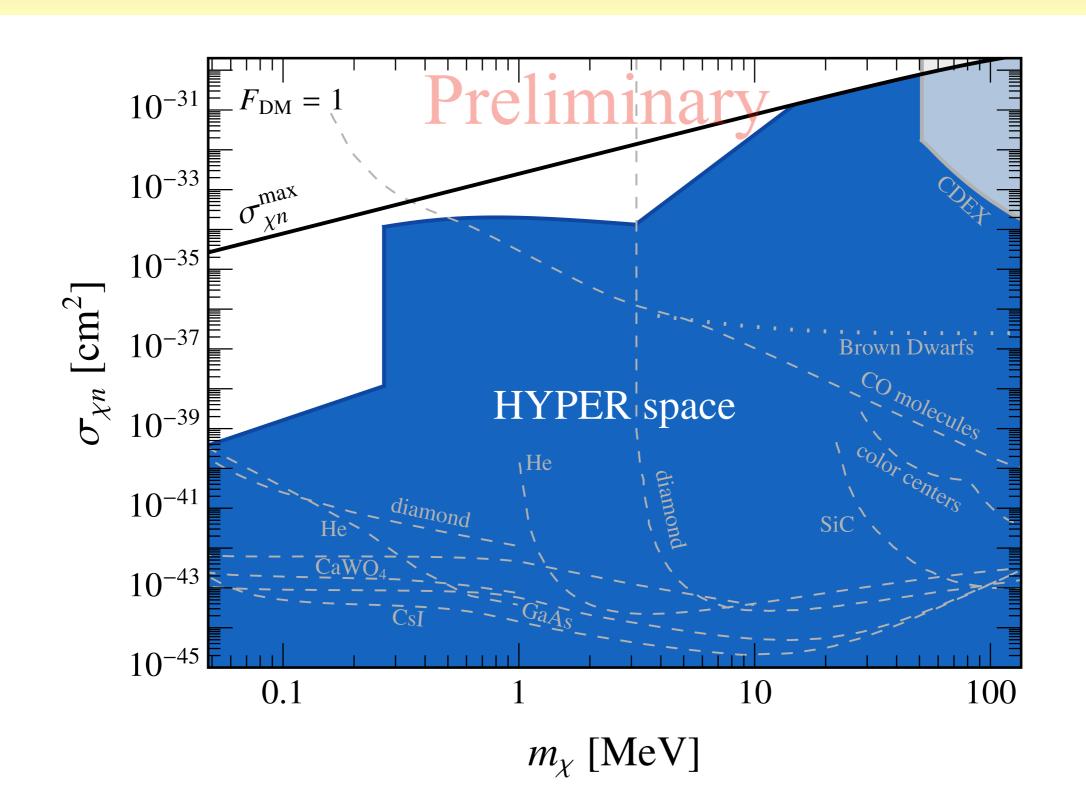
$$m_{\phi}^{i} \rightarrow m_{\phi} \ll m_{\phi}^{i}$$

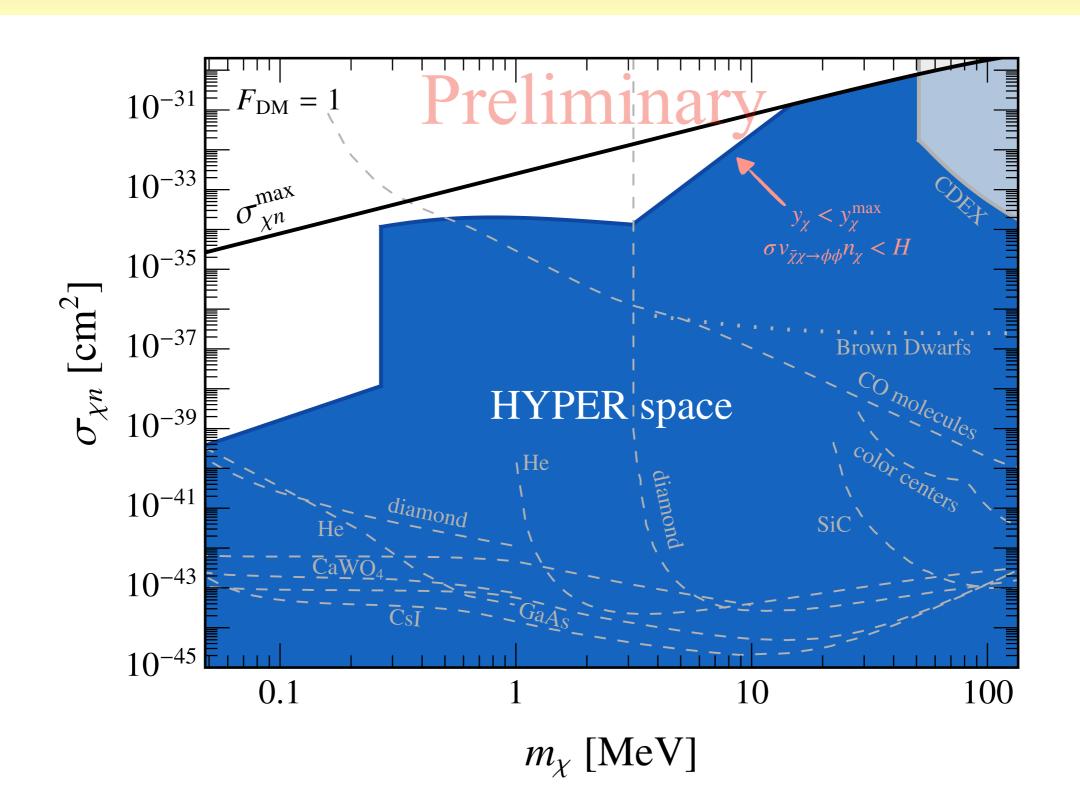
- Dark matter relic abundance must not change.
- BBN and CMB must be unaffected

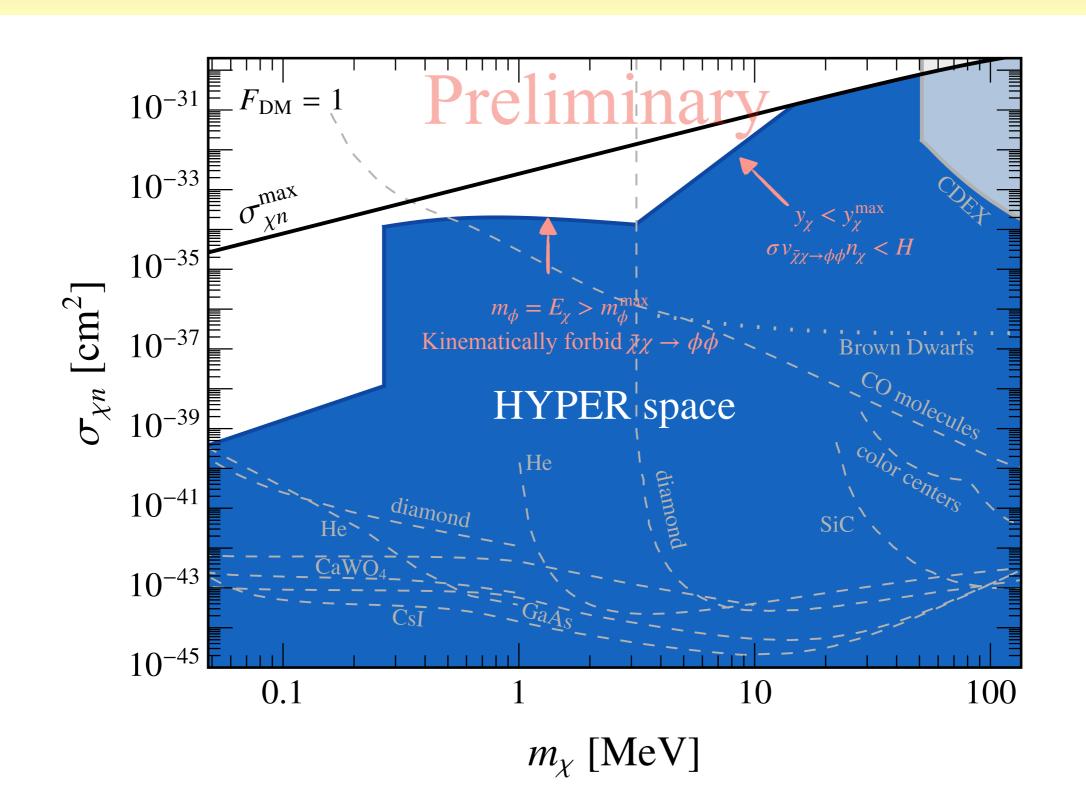
e.g. after the phase transition we must forbid or suppress processes such as

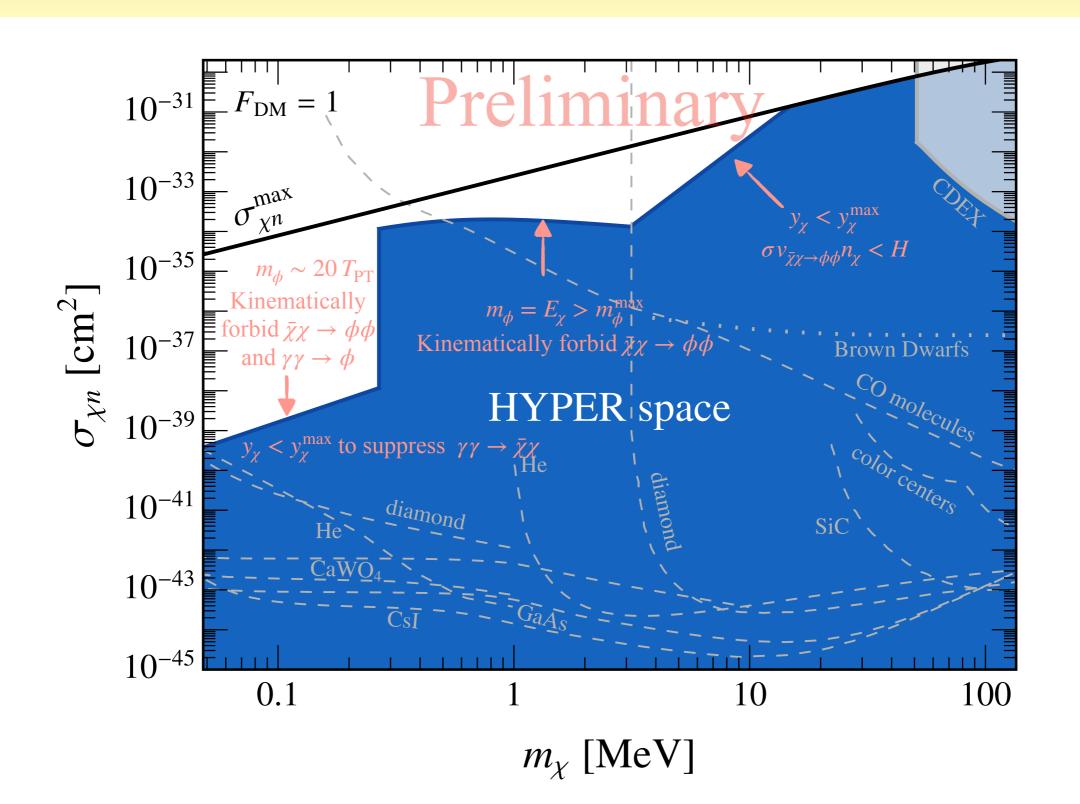
$$ar{\chi}\chi 
ightarrow ext{hadrons}$$
  $\mathcal{O}(10 ext{ keV}) \lesssim m_\chi < m_{\pi^0}$  hadrons  $ightarrow ar{\chi}\chi$  1 MeV  $\lesssim T_{\mathrm{PT}} \lesssim m_{\pi^0}$   $ar{\tau}\chi \gamma \rightarrow \phi\phi$   $\sigma v_{\chi\bar{\chi}\to\phi\phi} n_\chi < H$ 

• • •









## Summary

• Given present day constraints, it is unmotivated to think about cross sections larger than

$$\sigma_{\chi n} \lesssim 10^{-36} - 10^{-30} \,\mathrm{cm}^2$$
 for  $10 \,\mathrm{keV} < m_\chi < 100 \,\mathrm{MeV}$ 

• It is not easy to find a dark matter model that realizes such large cross sections, or in general live in the parameter space of interest to proposed light dark matter direct detection experiments. However, HYPERs is one such candidate.

### Outlook/Future Directions

- Derive  $\sigma_{\chi e}^{\rm max}$  and leptophilic HYPER models! Would likely require  $T_{\rm PT} \lesssim m_e$
- Fully explore the HYPER space of the hadrophilic hyper model. Perhaps considering vector mediators as well.
- Details of the dark sector phase transition.
- And many more

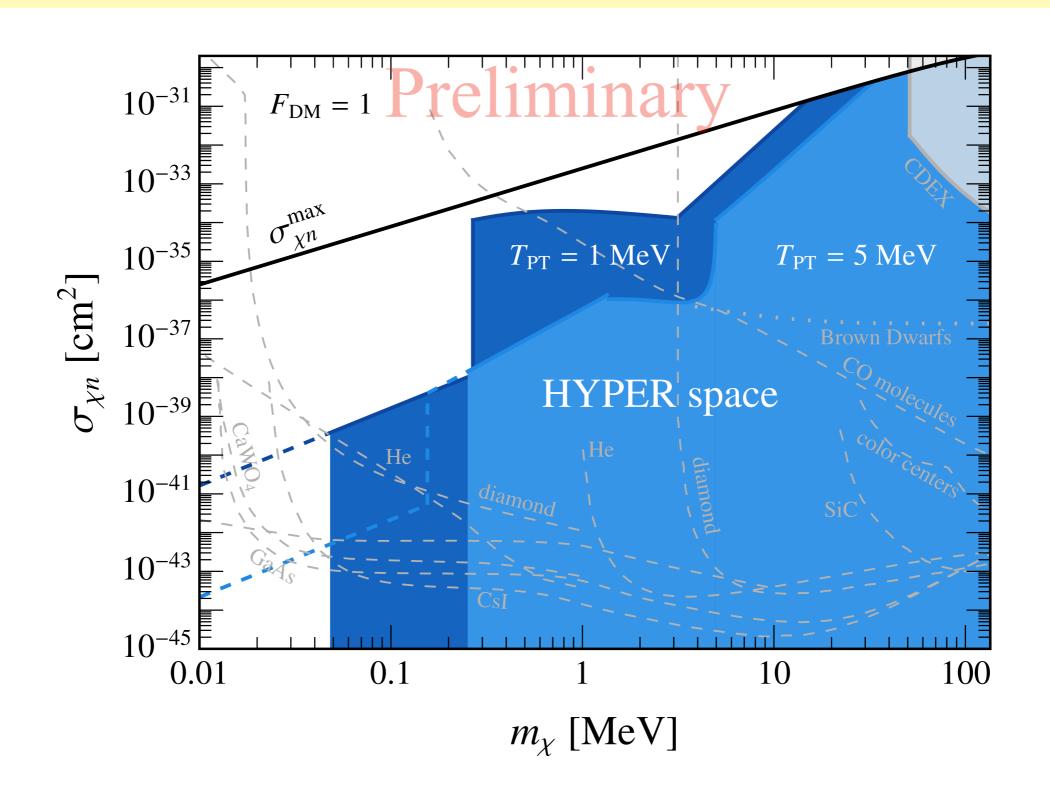
### Outlook/Future Directions

- Derive  $\sigma_{\chi e}^{\rm max}$  and leptophilic HYPER models! Would likely require  $T_{\rm PT} \lesssim m_e$
- Fully explore the HYPER space of the hadrophilic hyper model. Perhaps considering vector mediators as well.
- Details of the dark sector phase transition.
- And many more

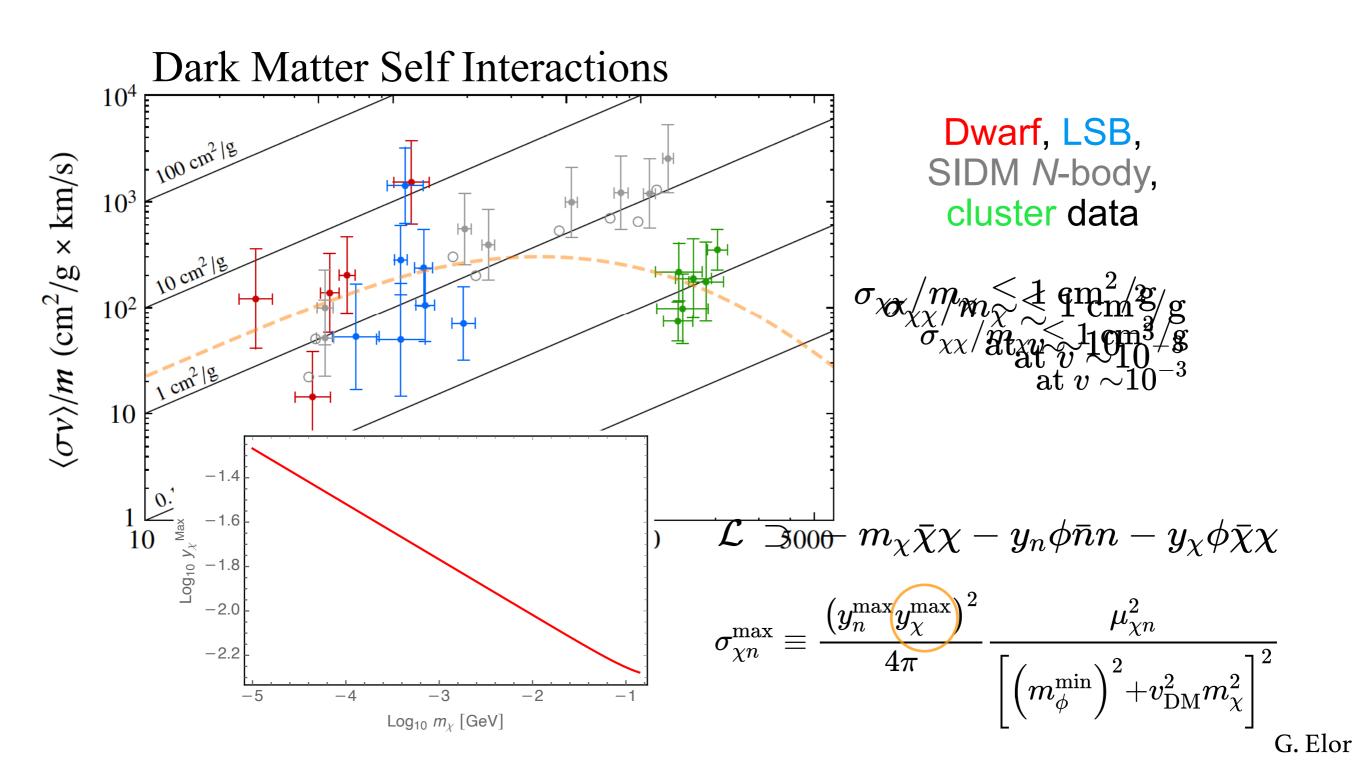
#### Thanks!

### Back ups

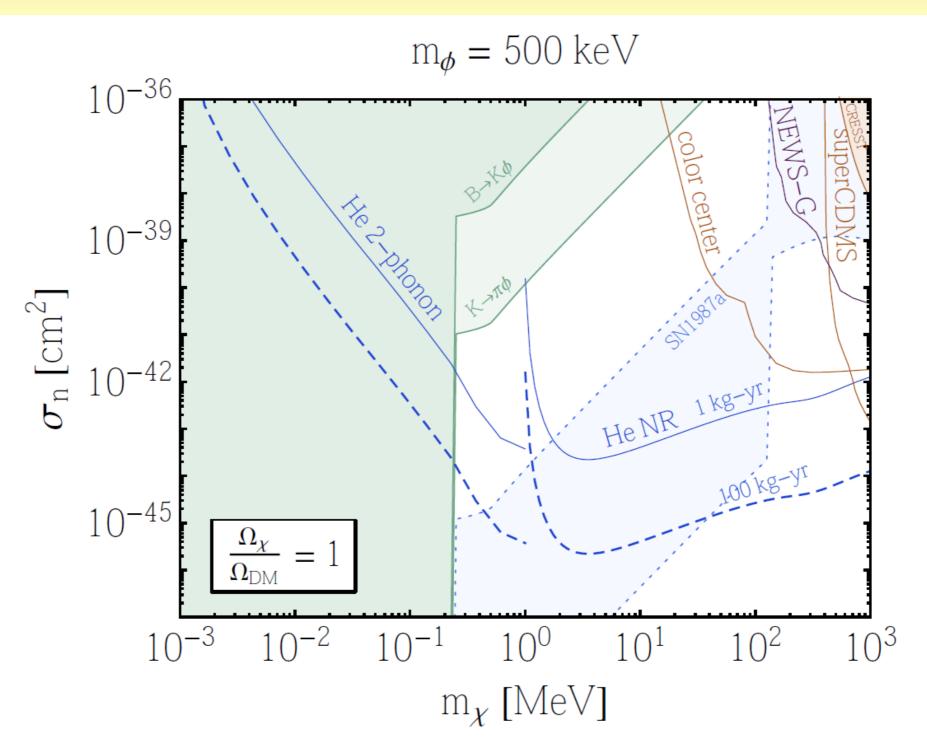
### Higher Temperatures



### Estimating $\sigma_{n\chi}^{\text{max}}$



#### Sub-GeV DM Detetors



Superfluid Helium [1611.06228, 1709.07882]