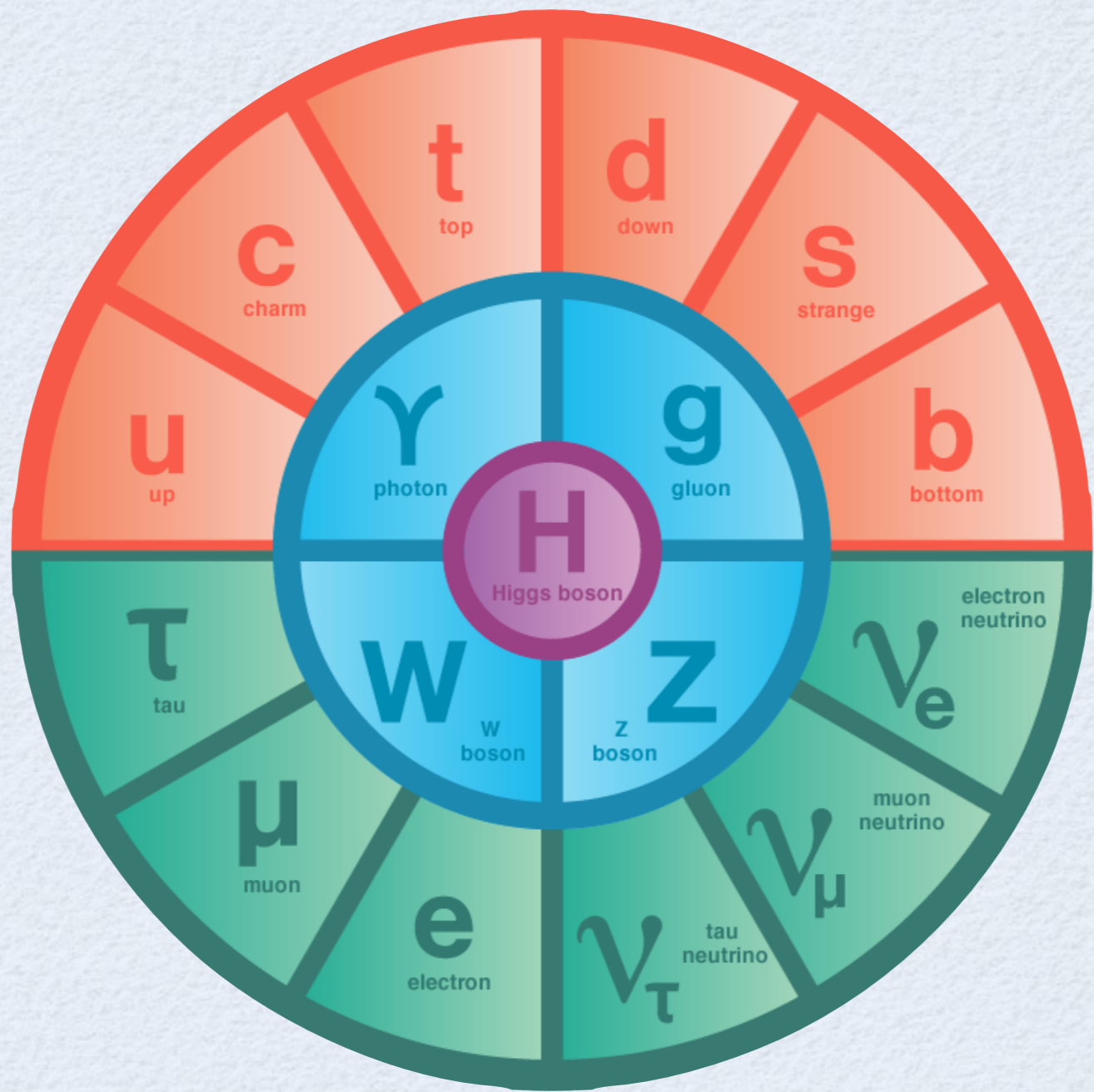


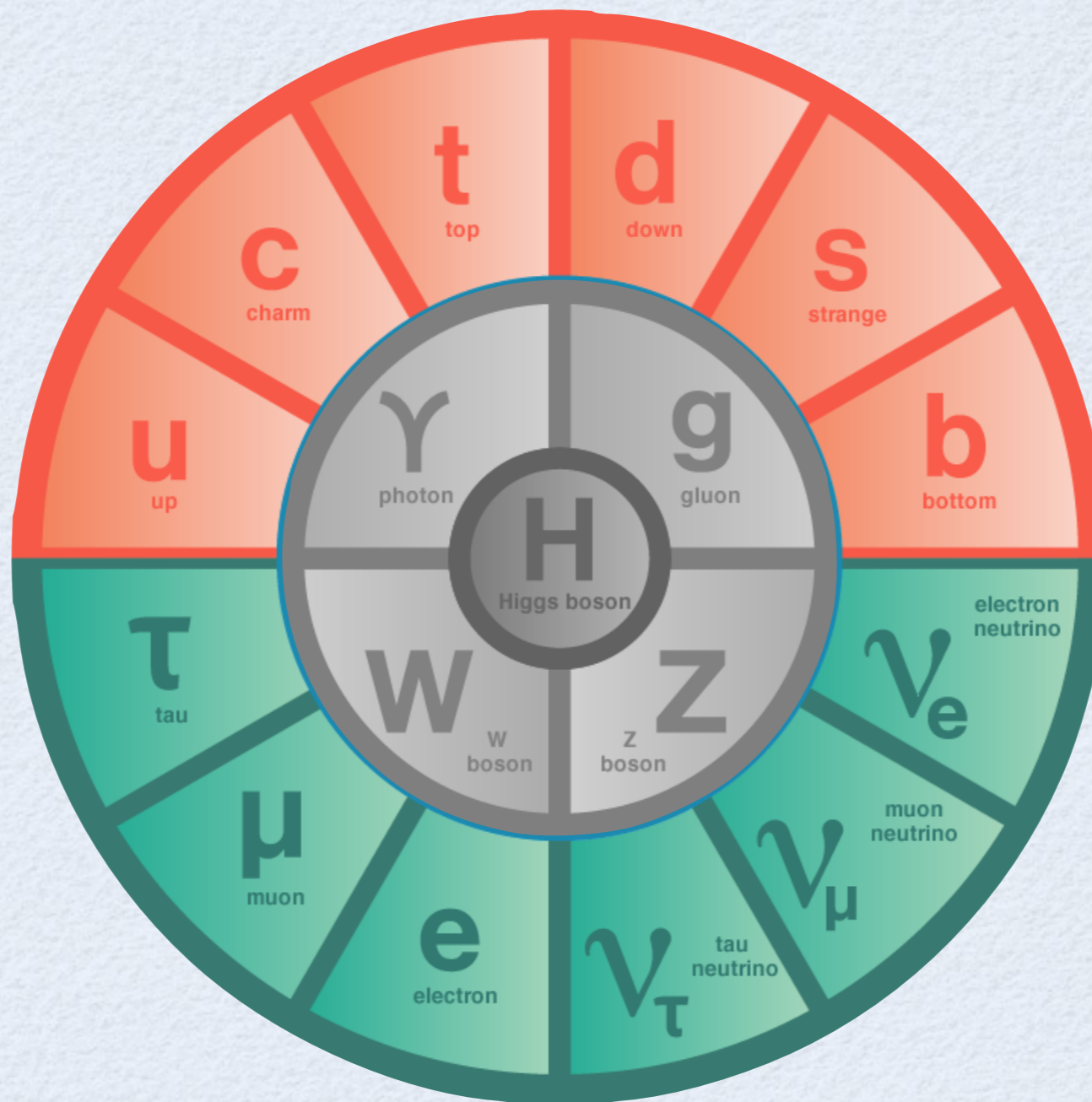


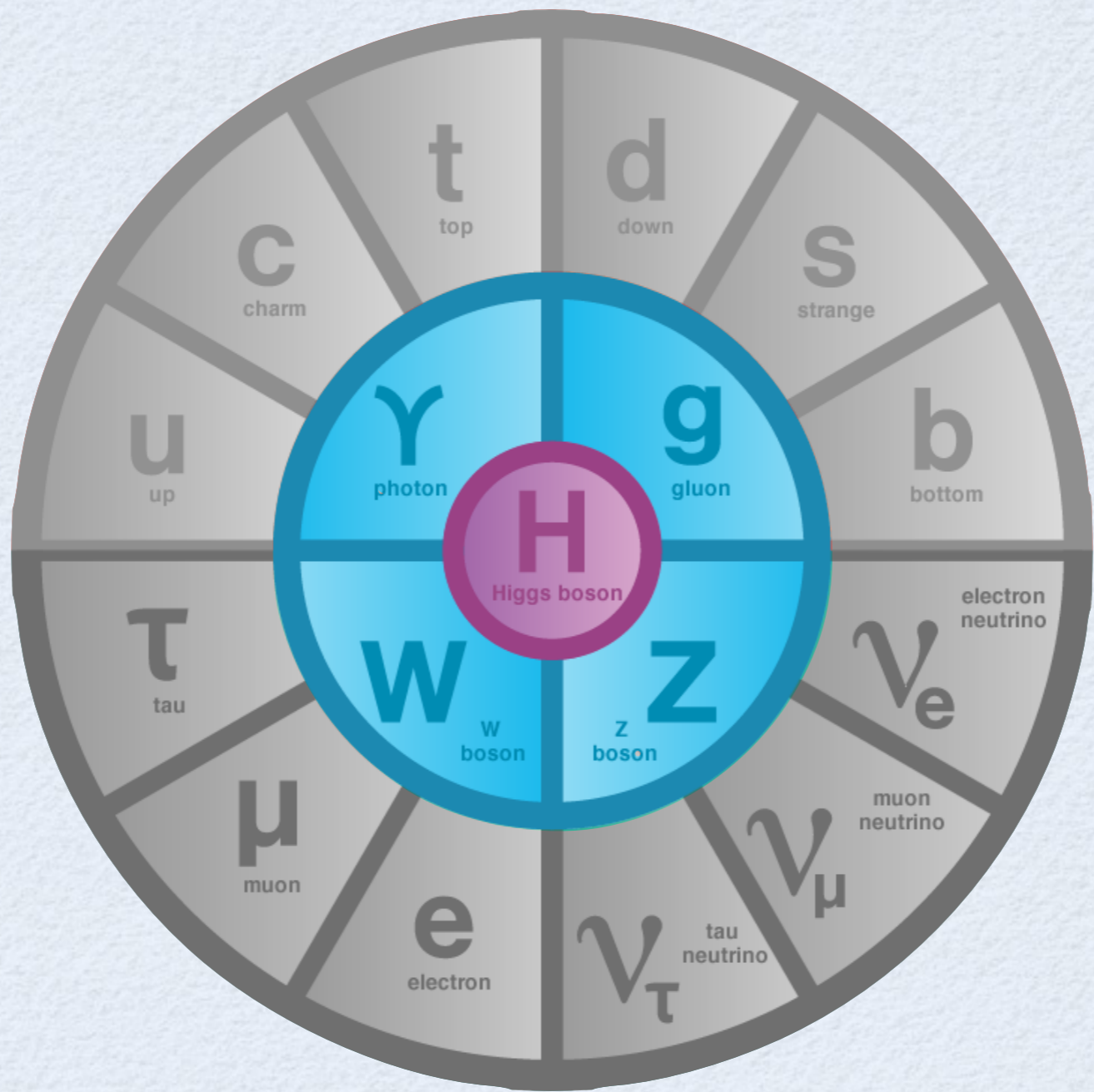
PROBING NEW PHYSICS ACROSS SCALES

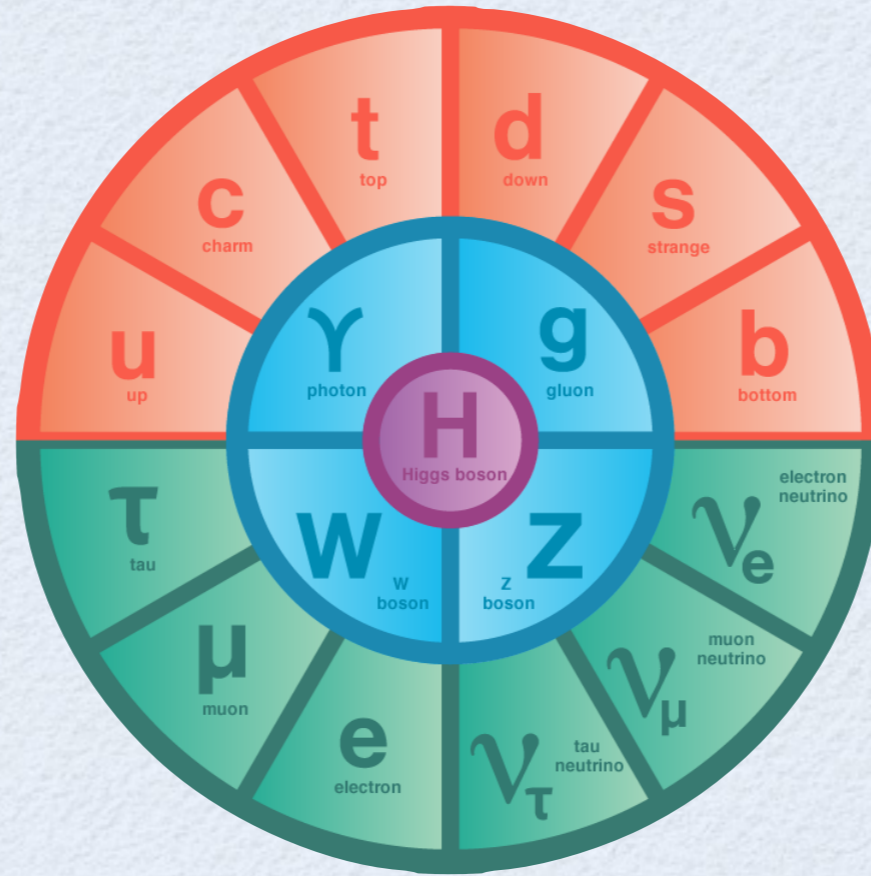
Yotam Soreq

Theoretical Particle Physics Seminar, University of Zurich, Nov. 27, 2018









if the Standard Model is so successful,
why do we look for new physics?

WHY NEW PHYSICS?

experimental evidences:

- Baryogenesis
- ν oscillations
- dark matter

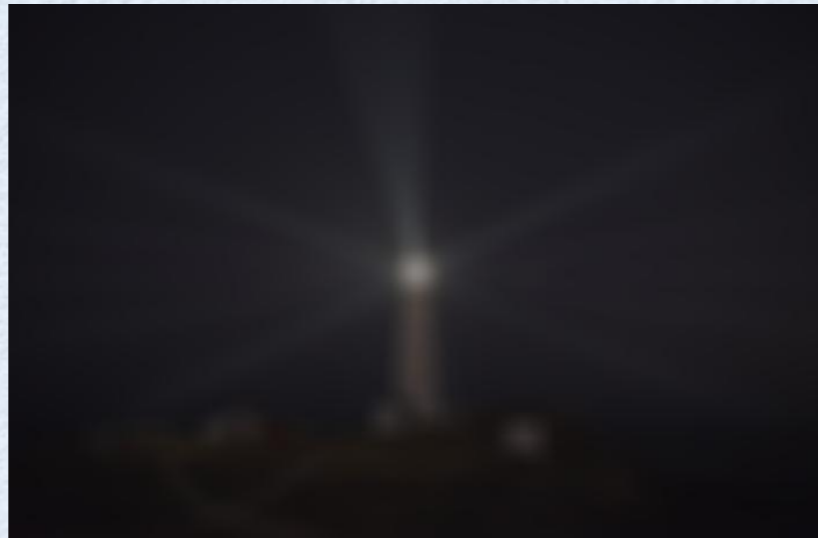
theory issues:

- tuning of the Higgs mass
- the strong CP problem
- the flavor puzzle



the quest for new physics

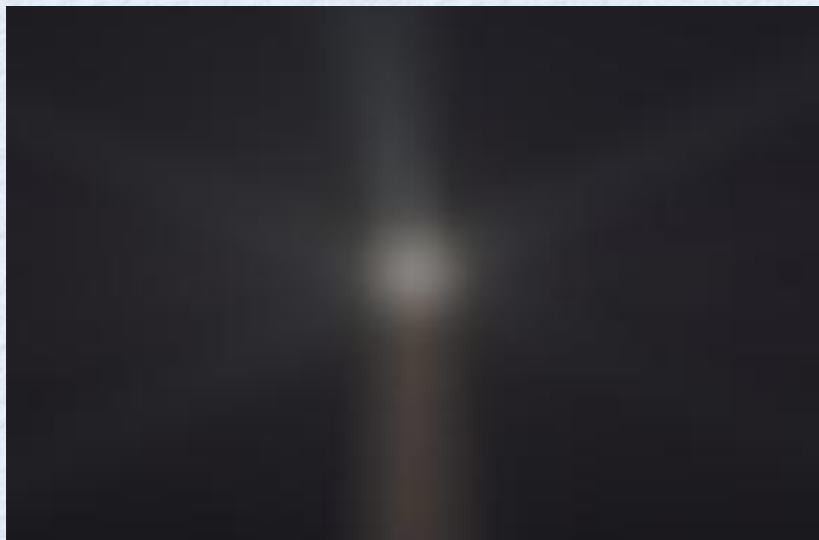
THE QUEST FOR NEW PHYSICS



THE QUEST FOR NEW PHYSICS



zoom

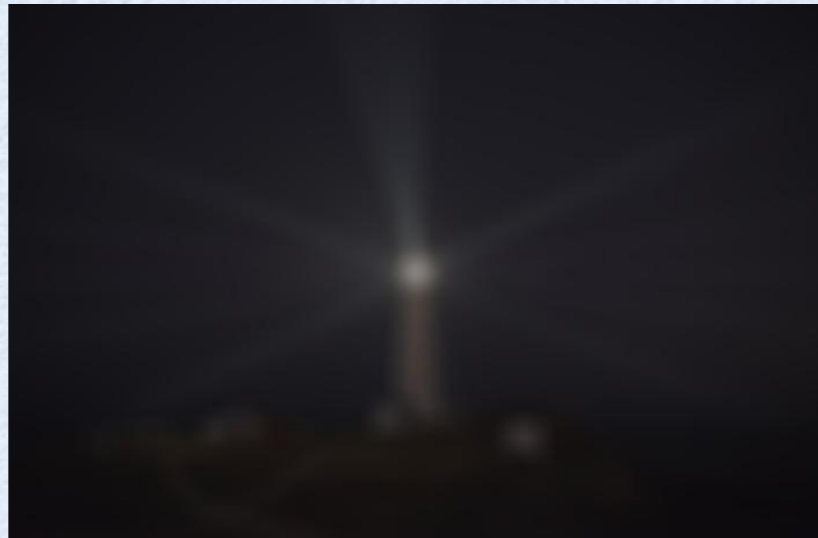


energy
frontier

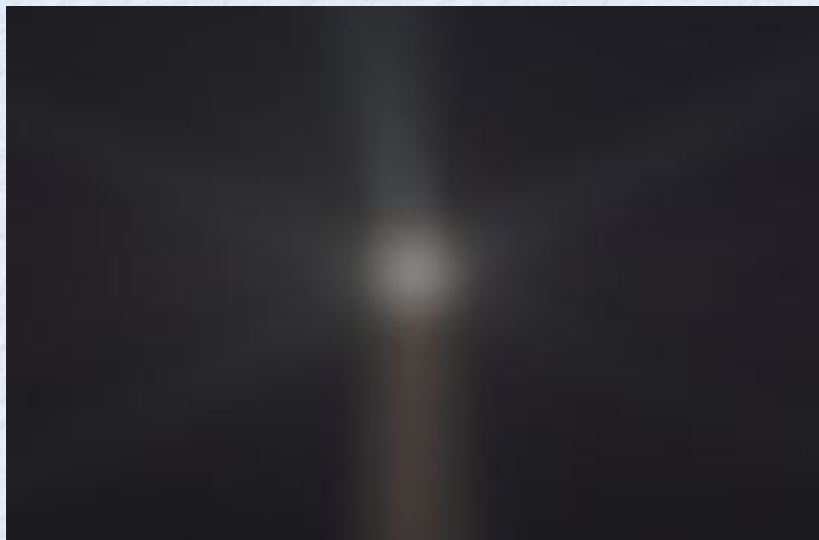
THE QUEST FOR NEW PHYSICS



zoom



brightness



energy
frontier



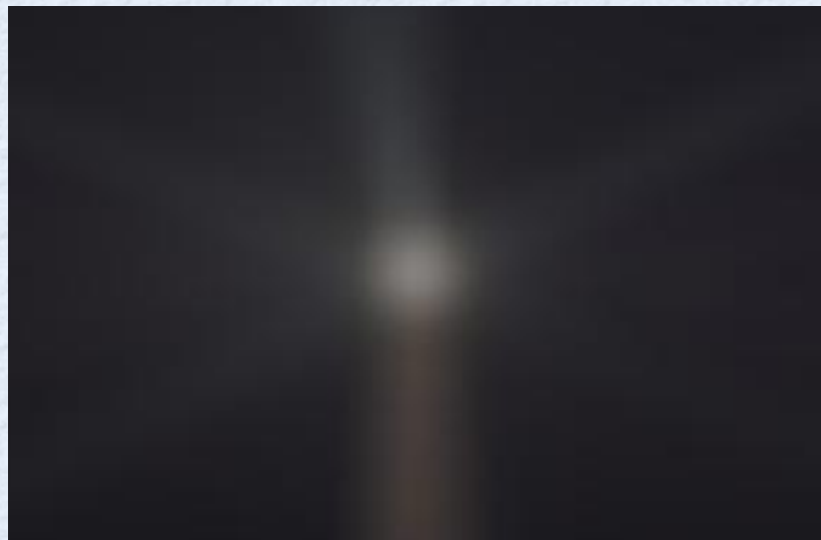
THE QUEST FOR NEW PHYSICS



zoom



brightness



energy
frontier

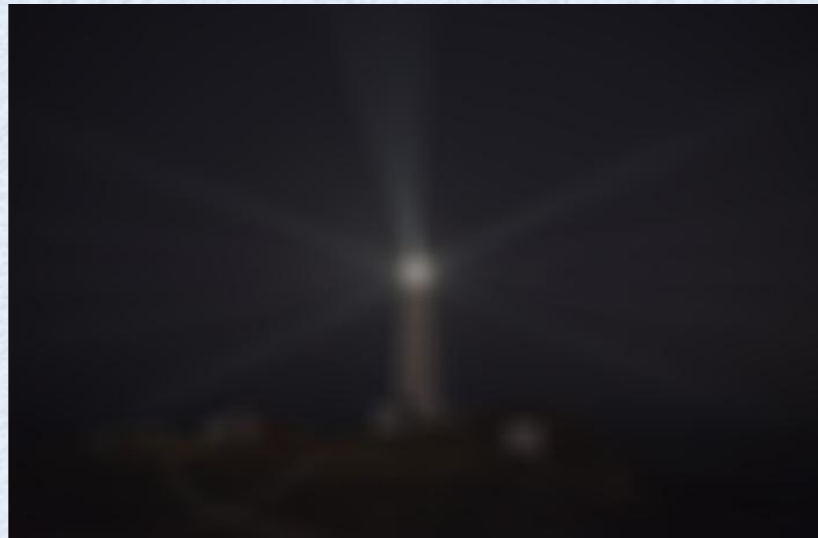


intensity
frontier

THE QUEST FOR NEW PHYSICS

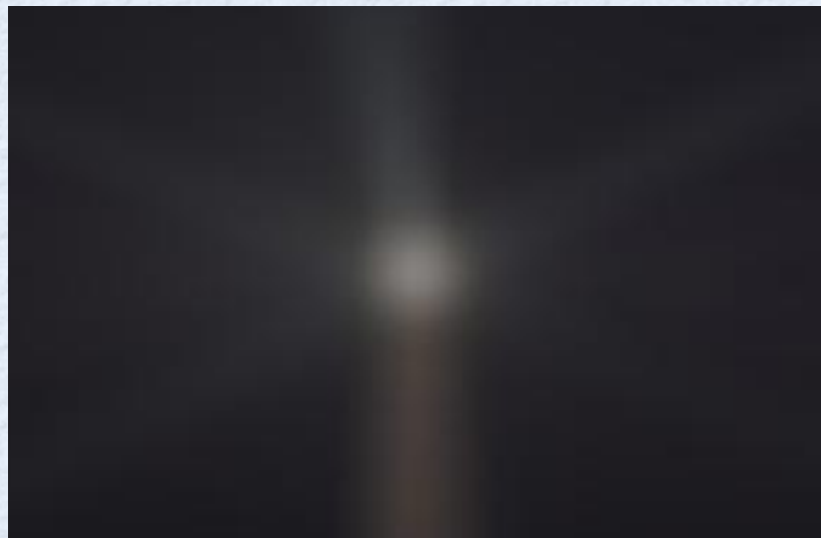


zoom



brightness

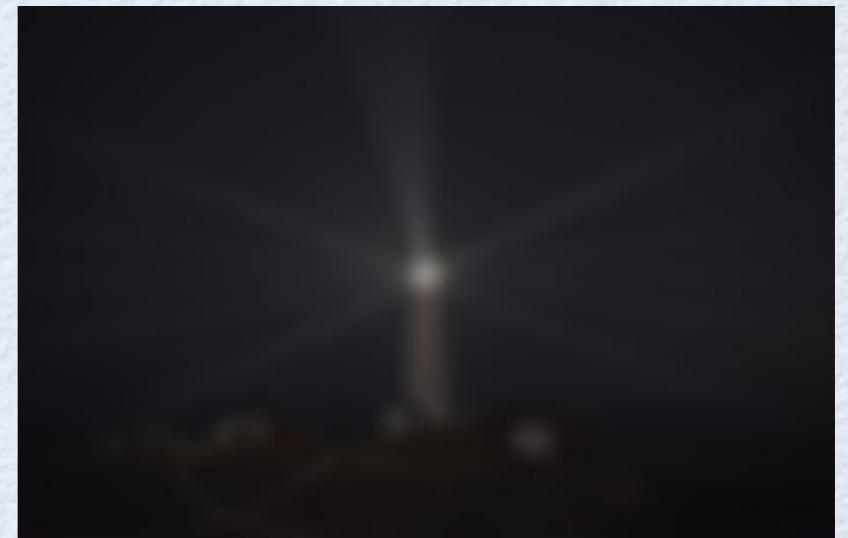
focus



energy
frontier



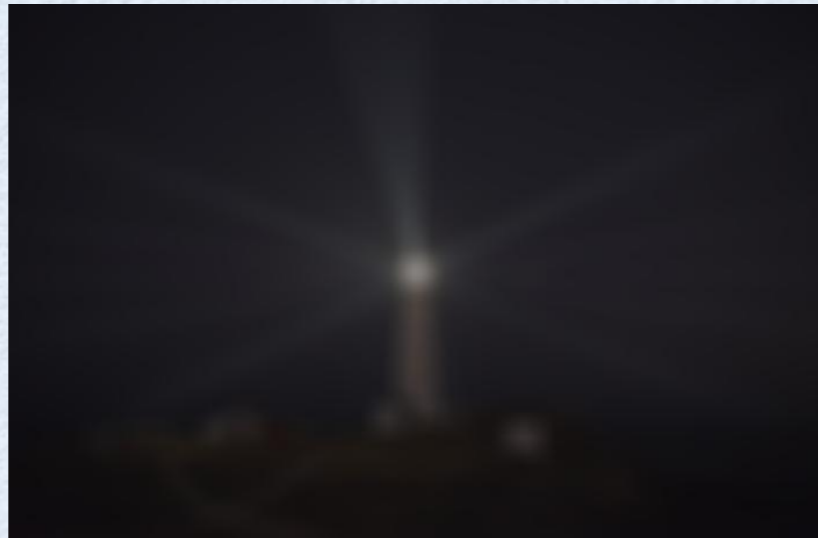
intensity
frontier



THE QUEST FOR NEW PHYSICS

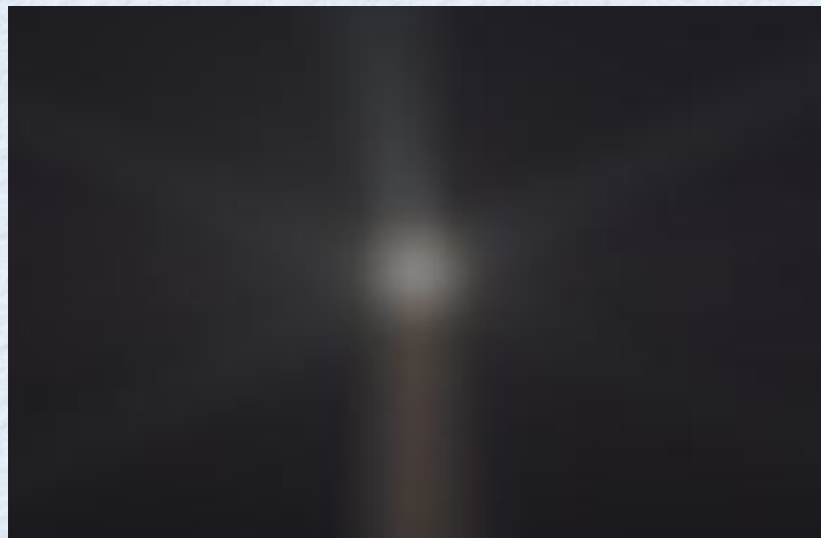


zoom



brightness

focus



energy
frontier



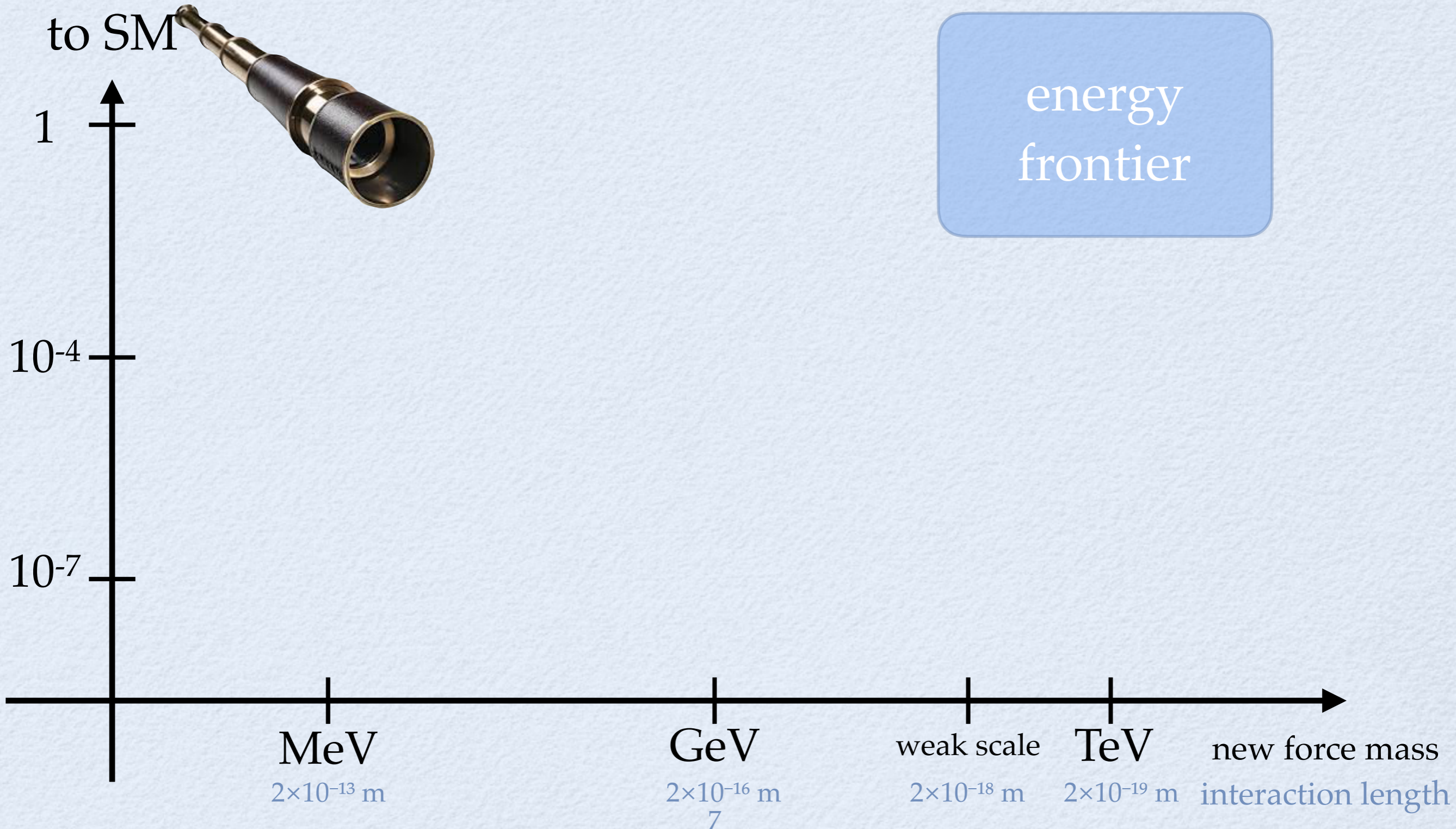
intensity
frontier



precision
frontier

THE QUEST FOR NEW PHYSICS

coupling
to SM



THE QUEST FOR NEW PHYSICS

coupling
to SM



1

10^{-4}

10^{-7}

energy
frontier

intensity
frontier

indirect

MeV

2×10^{-13} m

GeV

2×10^{-16} m
7

weak scale

2×10^{-18} m

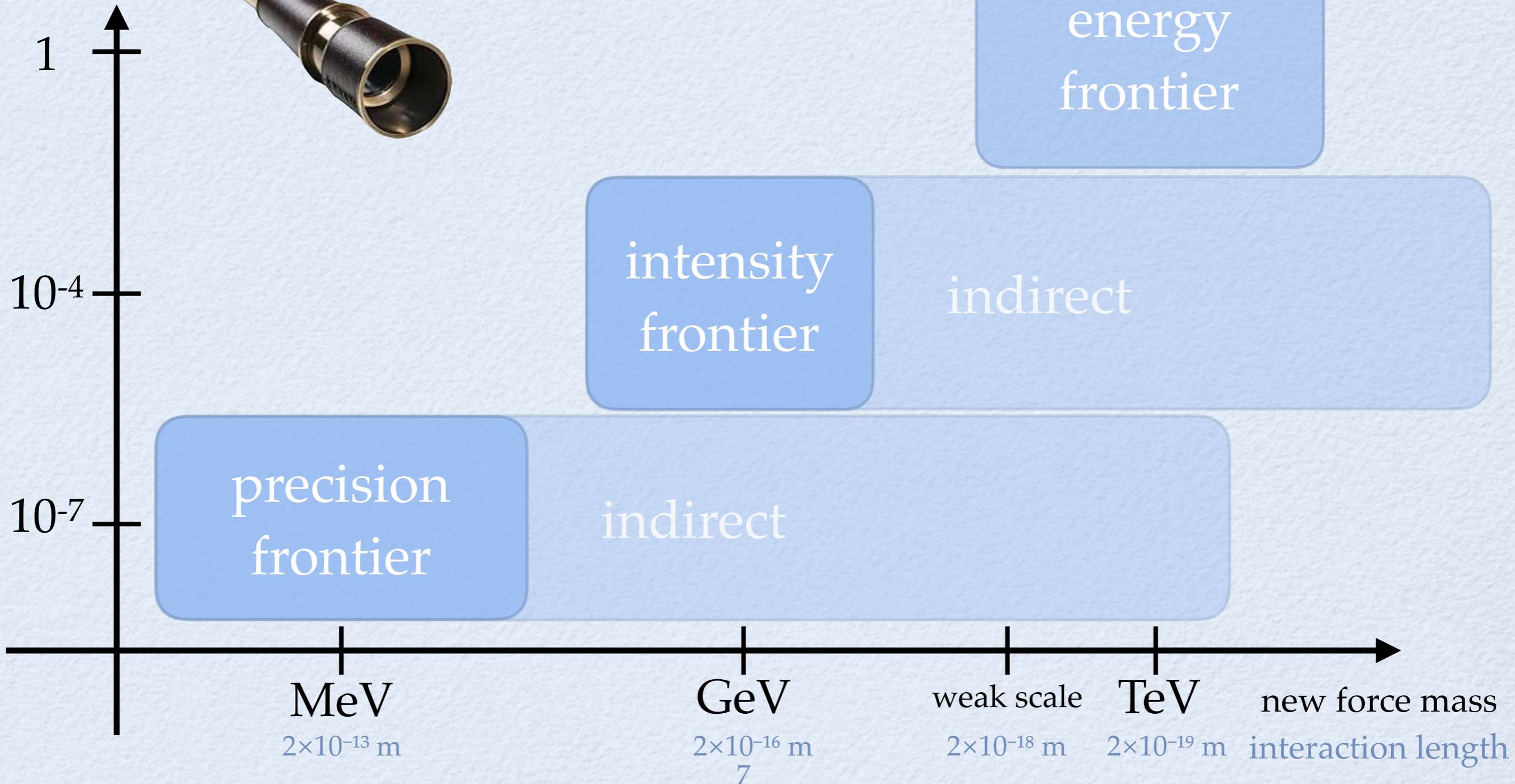
TeV

2×10^{-19} m

new force mass
interaction length

THE QUEST FOR NEW PHYSICS

coupling
to SM



THE QUEST FOR NEW PHYSICS

coupling
to SM



1

energy
frontier

10^{-4}

intensity
frontier

2. *dark photons at LHCb*
3. *ALP at the GeV Scale*

10^{-7}

precision
frontier

1. *probing long range force carriers*

MeV

2×10^{-13} m

GeV

2×10^{-16} m

weak scale

2×10^{-18} m

TeV

2×10^{-19} m

new force mass

interaction length

THE QUEST FOR NEW PHYSICS

coupling
to SM



1

10^{-4}

10^{-7}

precision
frontier

1. *probing long range
force carriers*

MeV

2×10^{-13} m

GeV

2×10^{-16} m

8

weak scale

2×10^{-18} m

TeV

2×10^{-19} m

new force mass

interaction length

Higgs





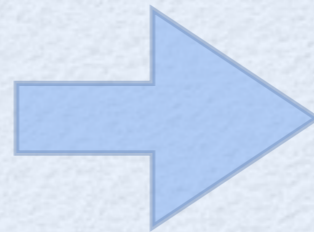
new force couples to matter

$$\frac{y_e y_A}{4\pi} \sin \theta \frac{e^{-m_\phi r}}{r}$$

PROBING NEW SPIN INDEPENDENT INTERACTIONS

ϕ - a new force carrier (spin 0, 1 or 2), mass m_ϕ

at atomic systems



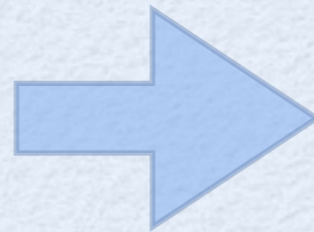
effective Yukawa like potential

interaction length $\sim 1/m_\phi$

PROBING NEW SPIN INDEPENDENT INTERACTIONS

ϕ - a new force carrier (spin 0, 1 or 2), mass m_ϕ

at atomic systems



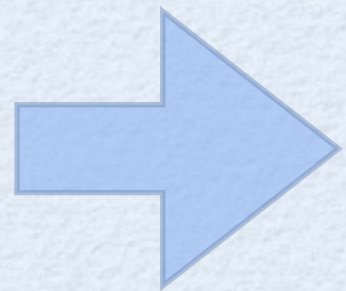
effective Yukawa like potential

$$\frac{y_e(y_p Z + (A - Z)y_n)}{4\pi} \frac{e^{-m_\phi r}}{r}$$

electron-nucleus
interaction

$$\frac{y_e^2}{4\pi} \frac{e^{-m_\phi r_{12}}}{r_{12}}$$

electron-electron
interaction



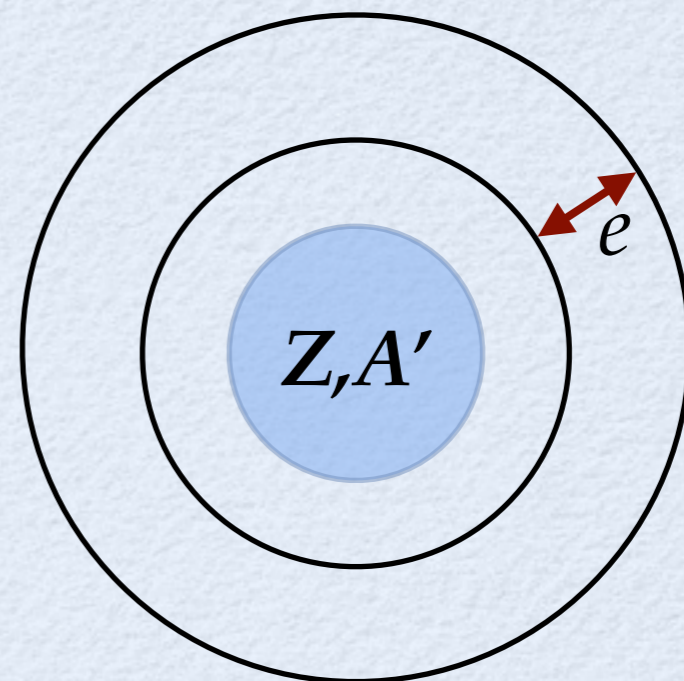
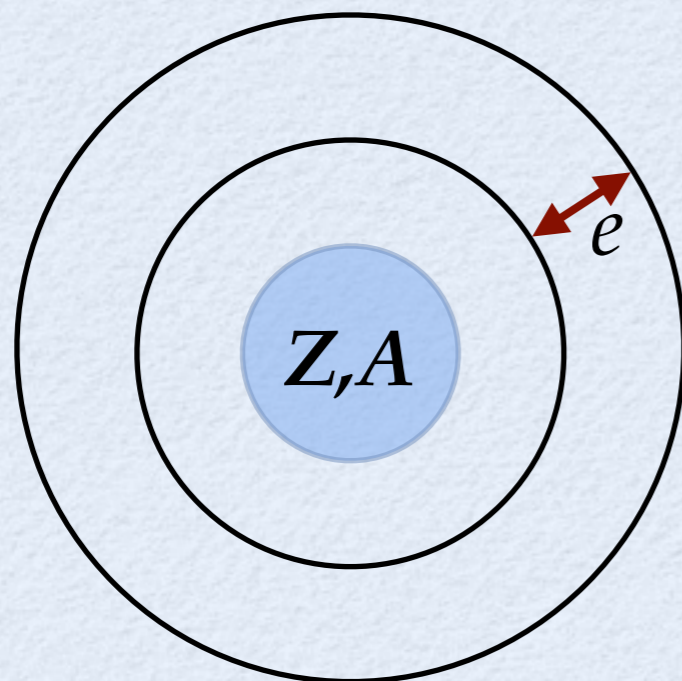
modify the electronic transition frequencies

interaction length $\sim 1/m_\phi$

ISOTOPE SHIFT

basic idea

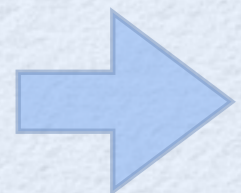
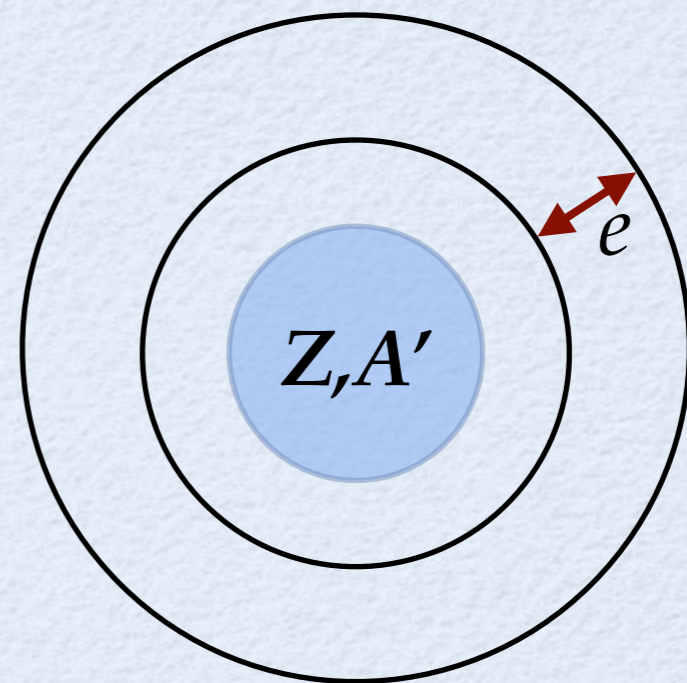
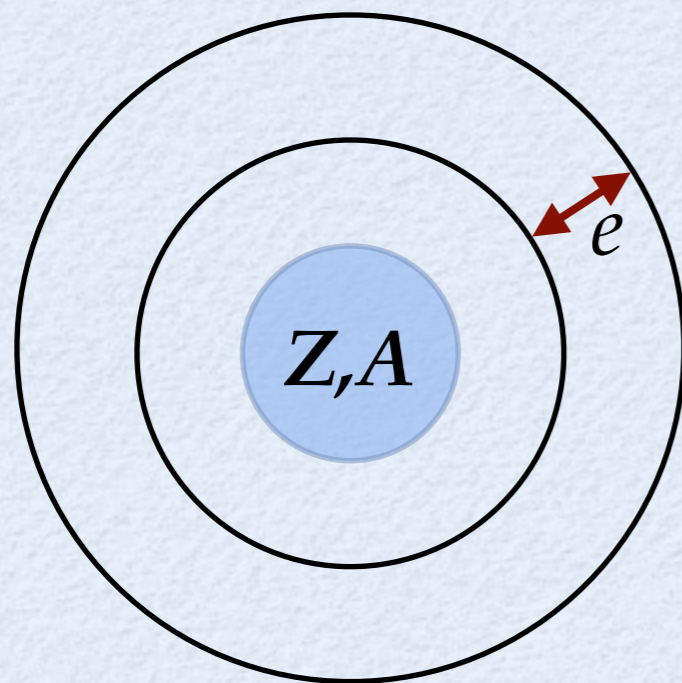
measure the same electronic transition in different isotopes



ISOTOPE SHIFT

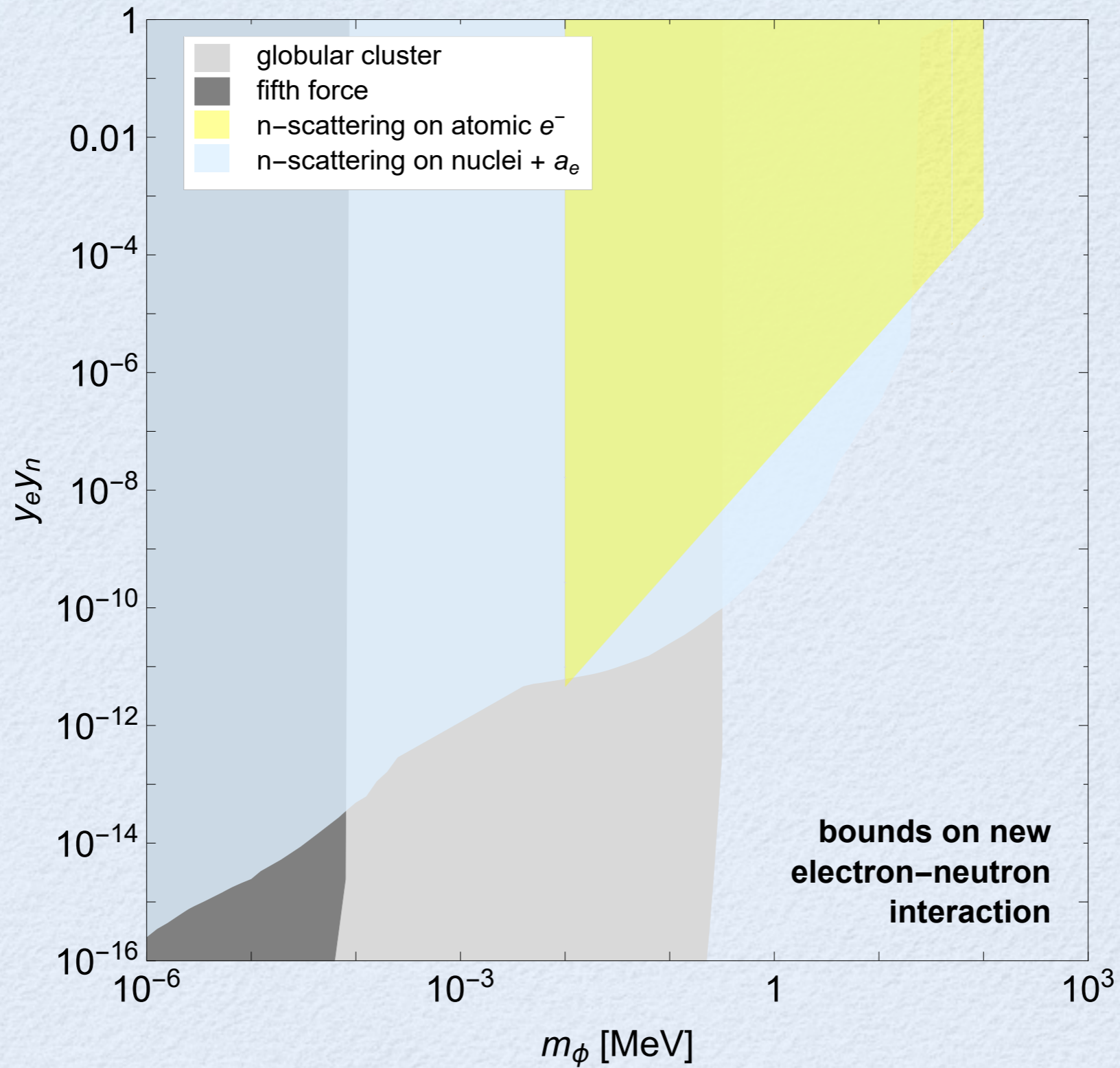
basic idea

measure the same electronic transition in different isotopes



probe the electron-neutron interaction ($\gamma_e \gamma_n$)

CURRENT BOUNDS: E-N INTERACTION



to maximize the sensitivity for new physics

to maximize the sensitivity for new physics



theory uncertainty
smaller than
experimental one

few electrons systems
(hydrogen, helium)

to maximize the sensitivity for new physics



theory uncertainty
smaller than
experimental one

few electrons systems
(hydrogen, helium)

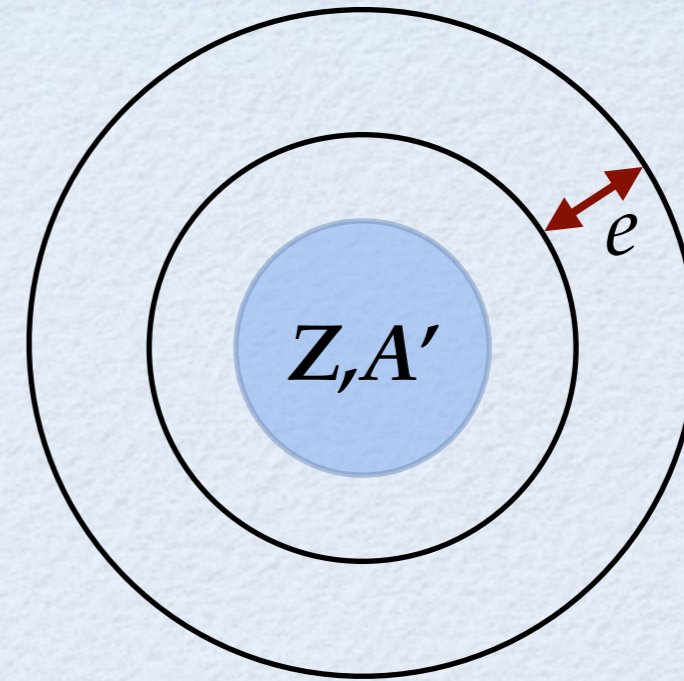
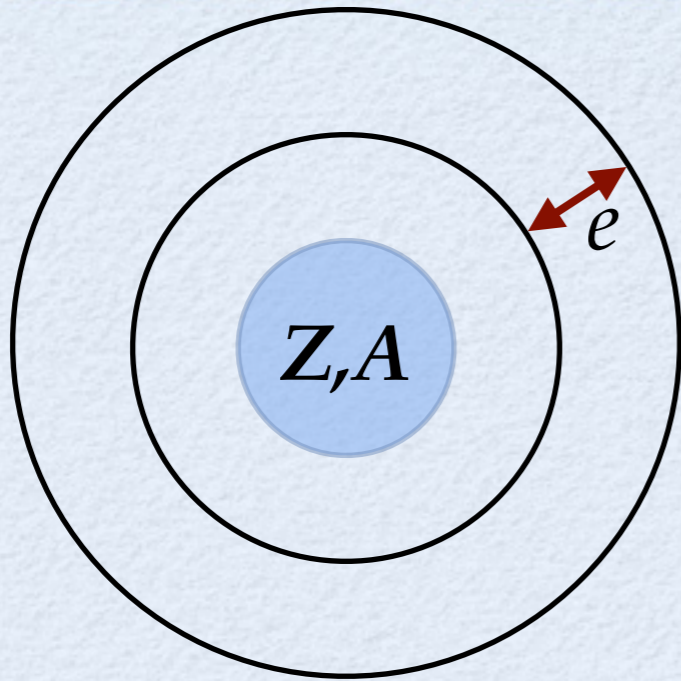


unique observable
which are insensitive to
theory error

heavy elements
(calcium, strontium, ytterbium)

few electrons systems

FEW ELECTRONS SYSTEMS - ISOTOPE SHIFT



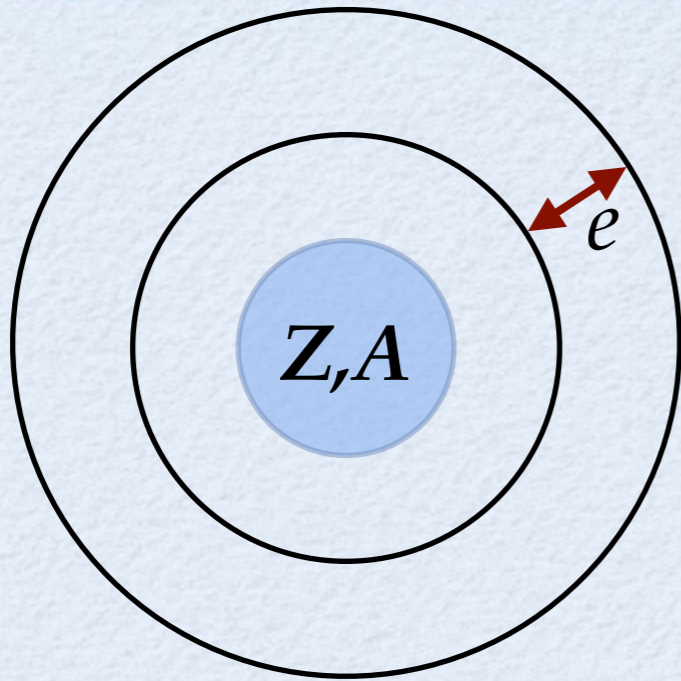
standard model

$$\nu_i^{A, A'} = \nu_i^A - \nu_i^{A'} = \nu_{i,0}^{A, A'} + F_i \langle r^2 \rangle_{A, A'}$$

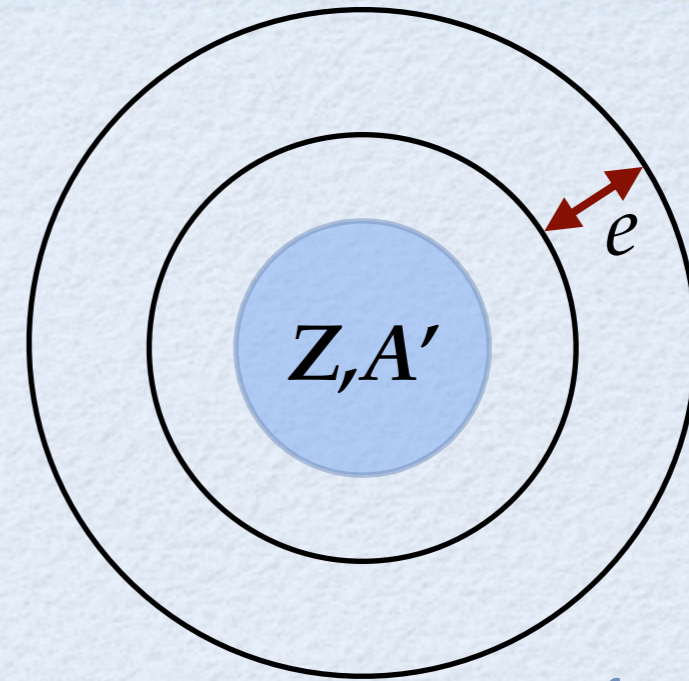
point like

finite radius

FEW ELECTRONS SYSTEMS - ISOTOPE SHIFT



standard model



new force

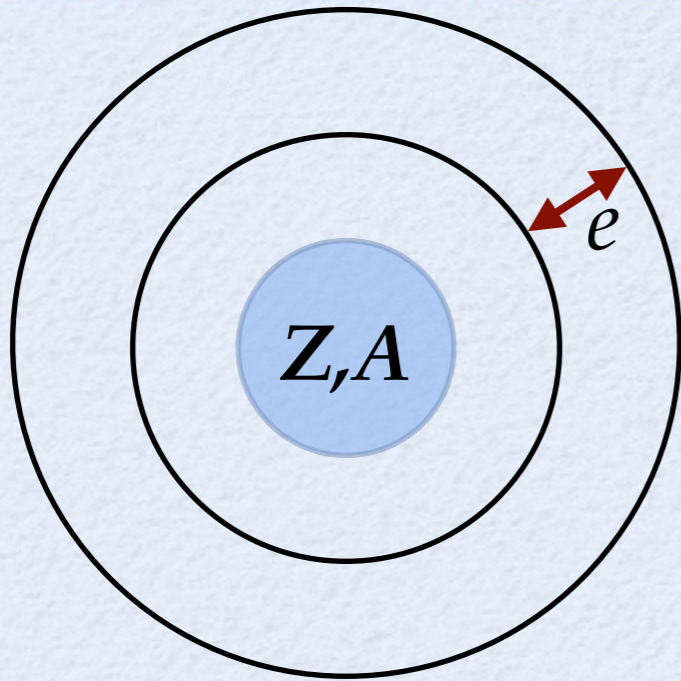
$$\nu_i^{A,A'} = \nu_i^A - \nu_i^{A'} = \nu_{i,0}^{A,A'} + F_i \langle r^2 \rangle_{A,A'} + y_e y_n X_i (A - A')$$

point like

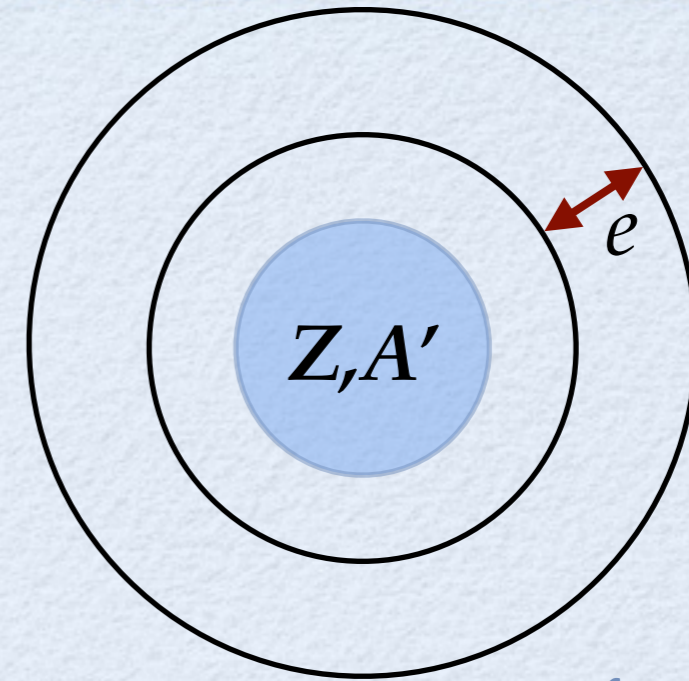
finite radius

e - n interaction

FEW ELECTRONS SYSTEMS - ISOTOPE SHIFT



standard model



new force

$$\nu_i^{A,A'} = \nu_i^A - \nu_i^{A'} = \nu_{i,0}^{A,A'} + F_i \langle r^2 \rangle_{A,A'} + y_e y_n X_i (A - A')$$

point like

finite radius

e - n interaction

the charged radius is the dominant error (from e -scattering)

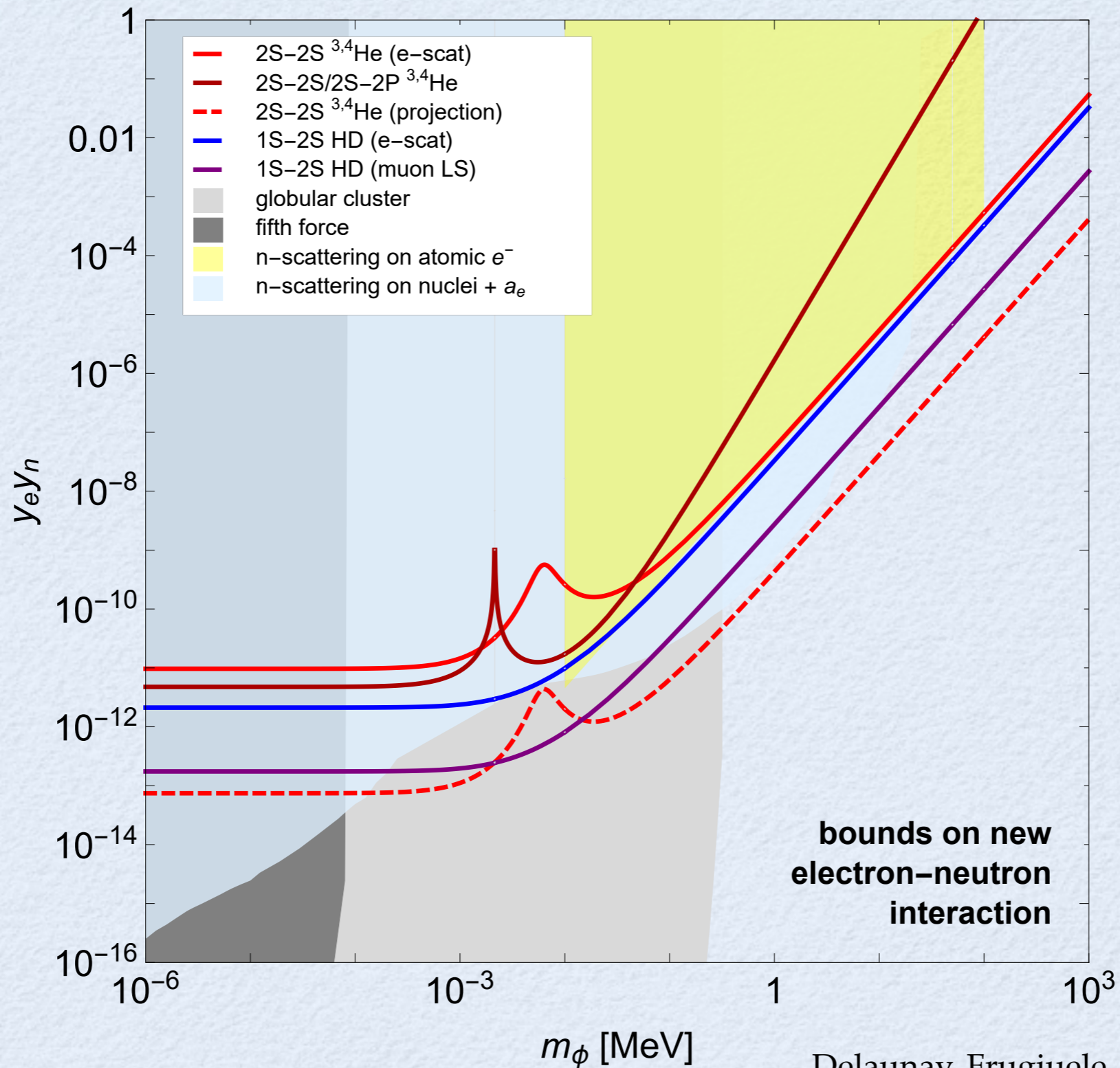


combing two transitions



no need for charge radius from scattering

FEW ELECTRONS SYSTEMS - ISOTOPE SHIFT



heavy elements

HEAVY ELEMENTS - ISOTOPE SHIFT

theory error \gg experimental error

$$\nu_i^{AA'} \equiv \nu_i^A - \nu_i^{A'}$$

HEAVY ELEMENTS - ISOTOPE SHIFT

theory error \gg experimental error

factorization of

**electronic
effects**

**nucleus
effects**

$$\nu_i^{AA'} \equiv \nu_i^A - \nu_i^{A'} = K_i \mu_{AA'} + F_i \delta \langle r^2 \rangle_{AA'} + \dots$$

Mass Shift

Field Shift
(short distance)

$$\mu_{AA'} \equiv \frac{1}{m_A} - \frac{1}{m_{A'}}$$

HEAVY ELEMENTS - ISOTOPE SHIFT

theory error \gg experimental error

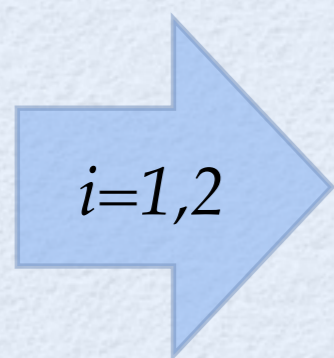
factorization of

$$\nu_i^{AA'} \equiv \nu_i^A - \nu_i^{A'} = K_i \mu_{AA'} + F_i \delta \langle r^2 \rangle_{AA'} + \dots$$

electronic effects nucleus effects

Mass Shift Field Shift
(short distance)

$$\mu_{AA'} \equiv \frac{1}{m_A} - \frac{1}{m_{A'}}$$



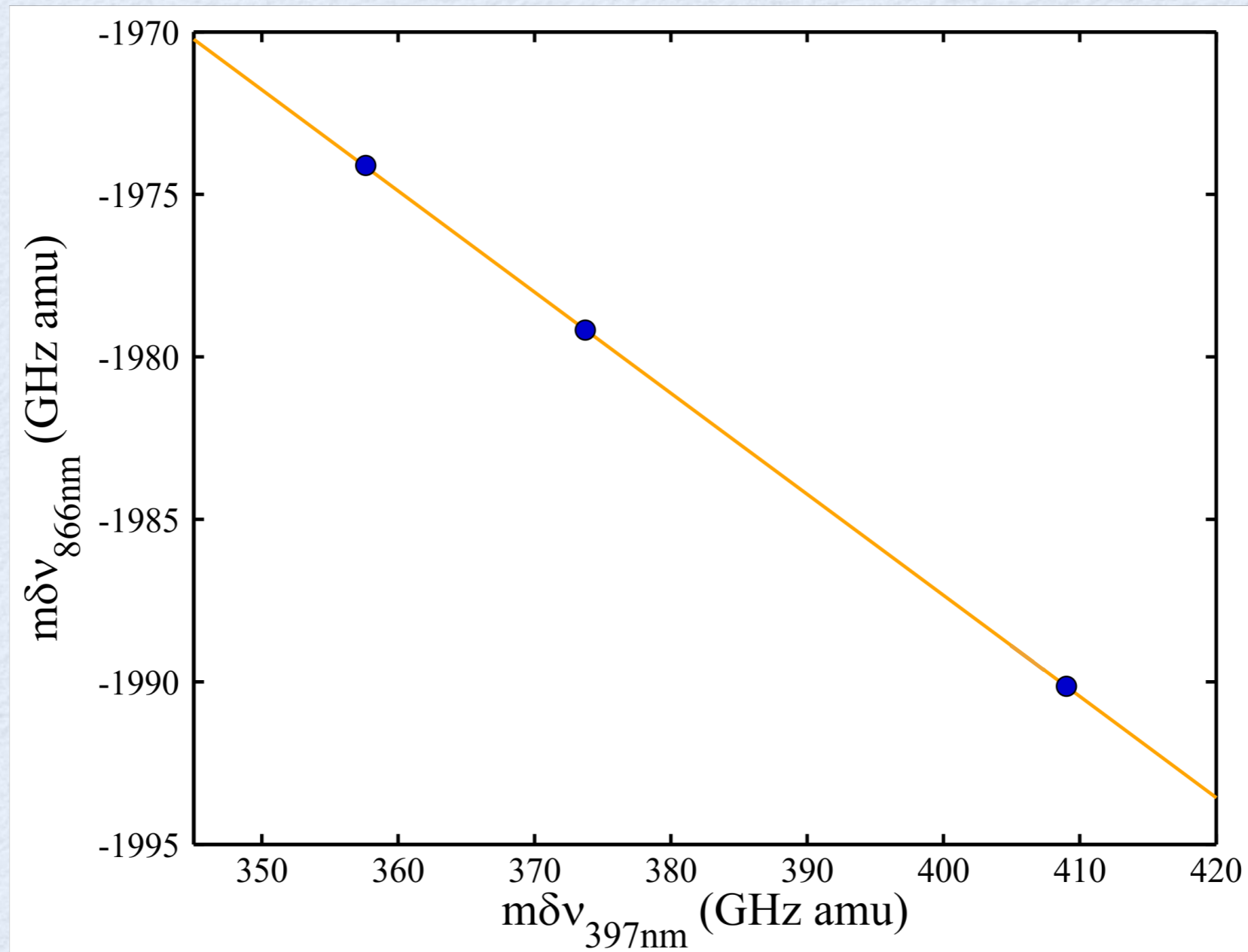
$$\left(\frac{\nu_2^{AA'}}{\mu_{AA'}} \right) = K_{21} + F_{21} \left(\frac{\nu_1^{AA'}}{\mu_{AA'}} \right)$$

linear relation in the SM

HEAVY ELEMENTS - ISOTOPE SHIFT

isotope shift of Ca^+

$4D \rightarrow 4P_{1/2}$



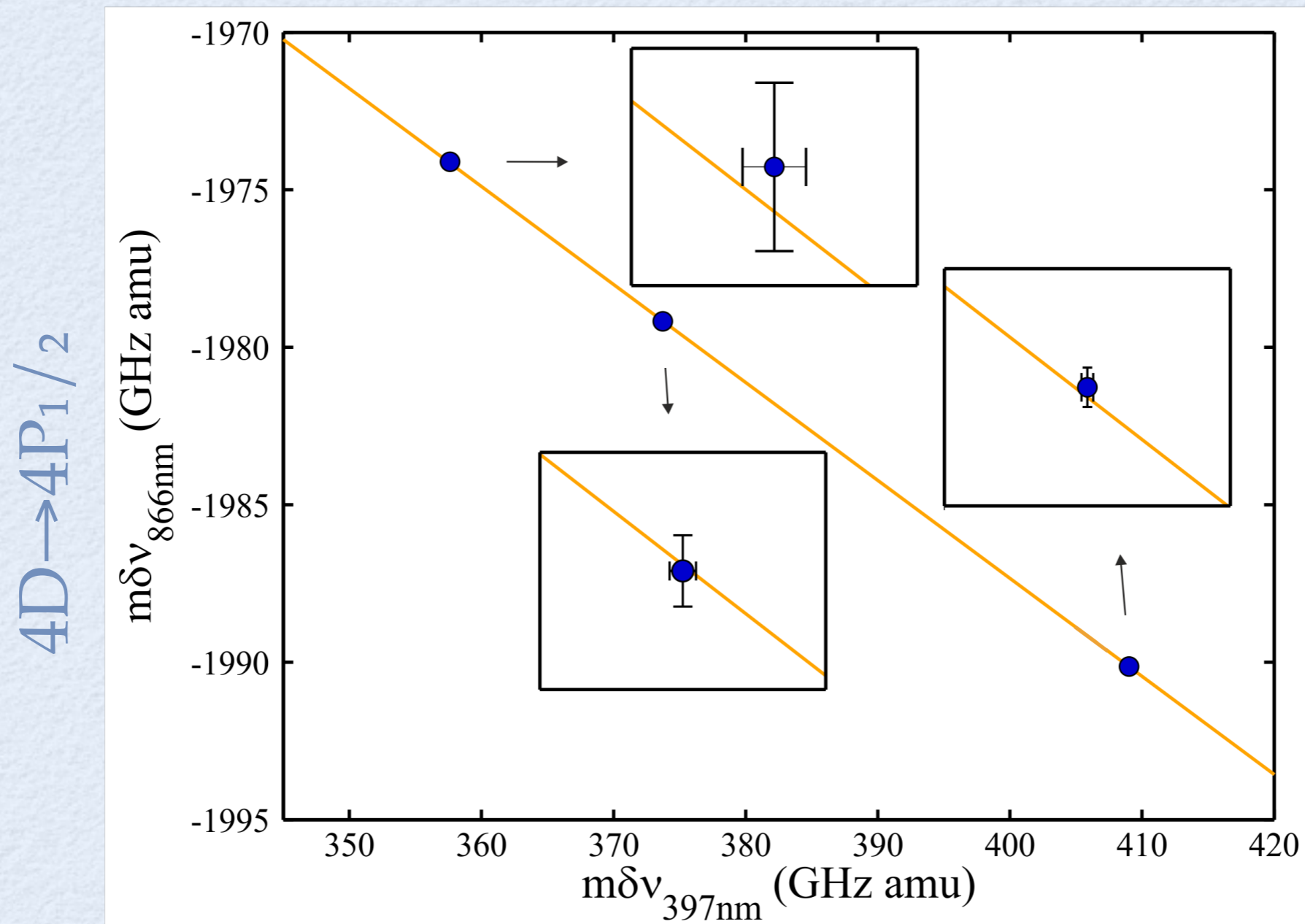
Gebert et al. 2015

$4S \rightarrow 4P_{1/2}$

$O(100\text{kHz})$ error

HEAVY ELEMENTS - ISOTOPE SHIFT

isotope shift of Ca^+



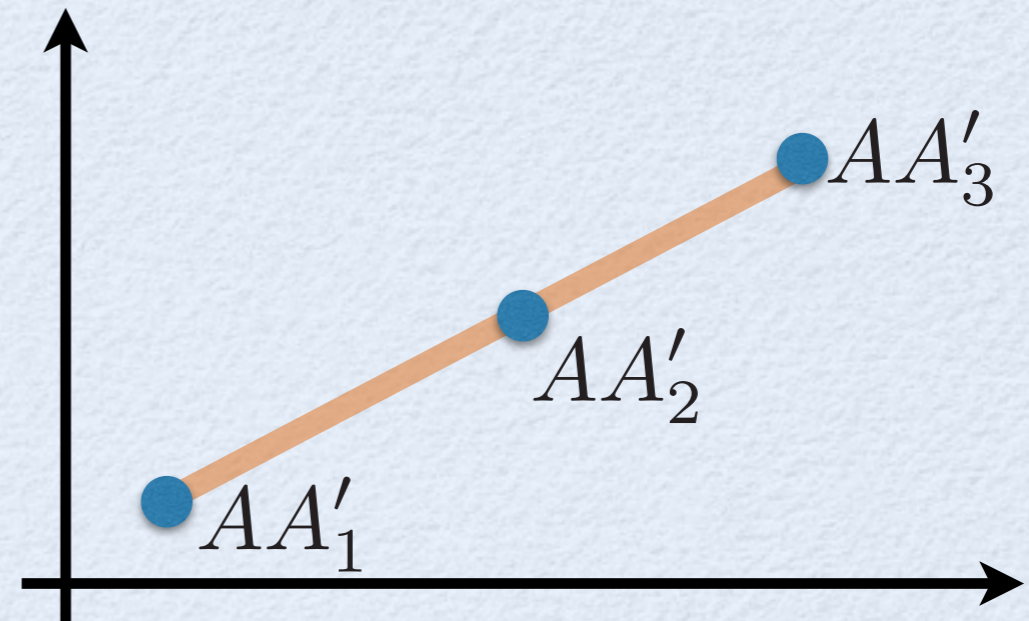
$4S \rightarrow 4P_{1/2}$

$O(100\text{kHz})$ error

HEAVY ELEMENTS - ISOTOPE SHIFT

$$\nu_i^{AA'} = K_i \mu_{A,A'} + F_i \langle r^2 \rangle_{A,A'} + y_e y_n X_i (A - A')$$

new force
(the *only* theory inputs)



Delaunay, Ozeri, Perez, **YS**, PRD 1601.05087

Berengut, Budker, Delaunay, Flambaum, Frugiuele,
Fuchs, Grojean, Harnik, Ozeri, Perez, **YS**, 1704.05068

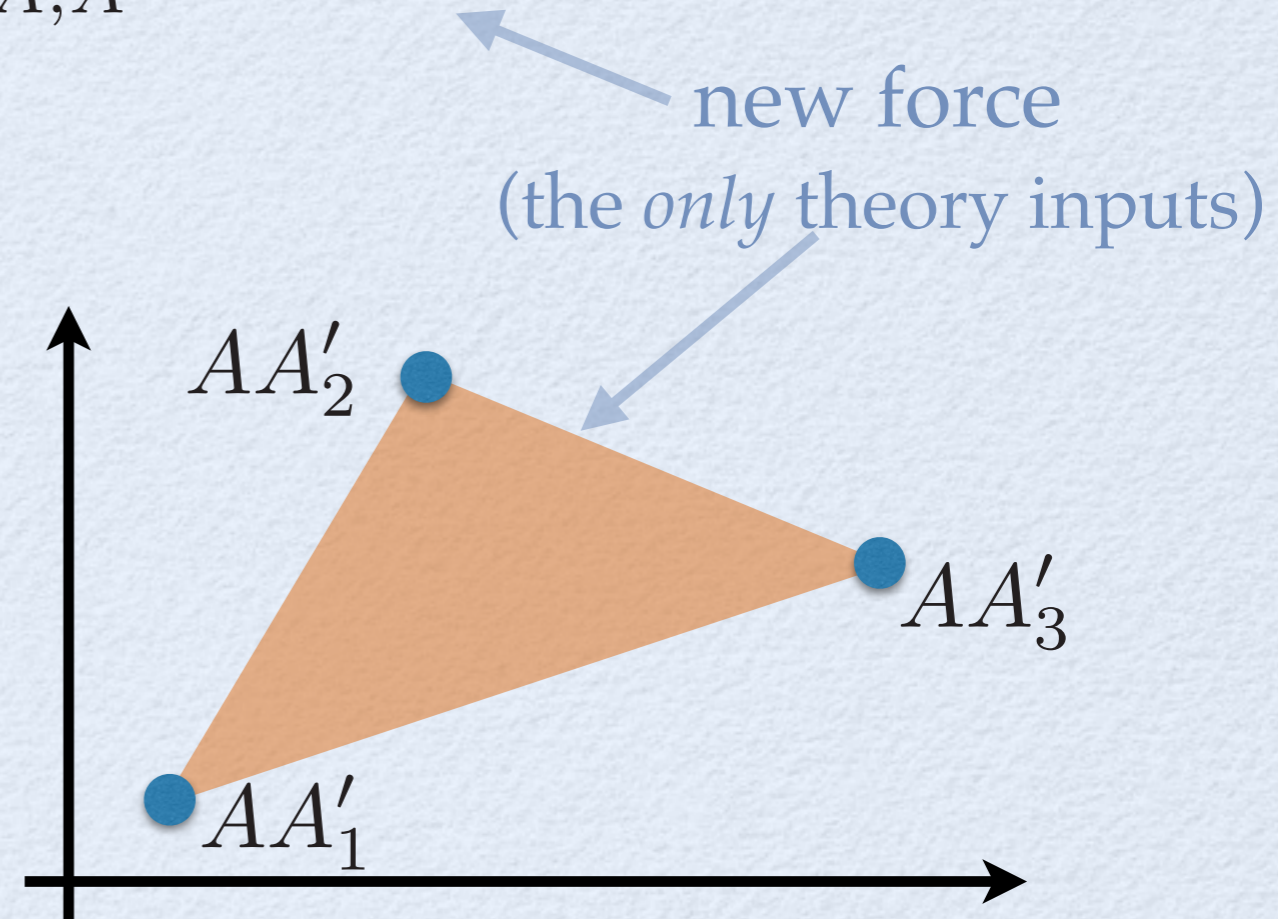
HEAVY ELEMENTS - ISOTOPE SHIFT

$$\nu_i^{AA'} = K_i \mu_{A,A'} + F_i \langle r^2 \rangle_{A,A'} + y_e y_n X_i (A - A')$$

- long range new physics
- misalignment between the new physics and the isotope shift



nonlinear King plot from
new physics



similar to data driven background estimation at the LHC

Delaunay, Ozeri, Perez, **YS**, PRD 1601.05087

Berengut, Budker, Delaunay, Flambaum, Frugiuele,
Fuchs, Grojean, Harnik, Ozeri, Perez, **YS**, 1704.05068

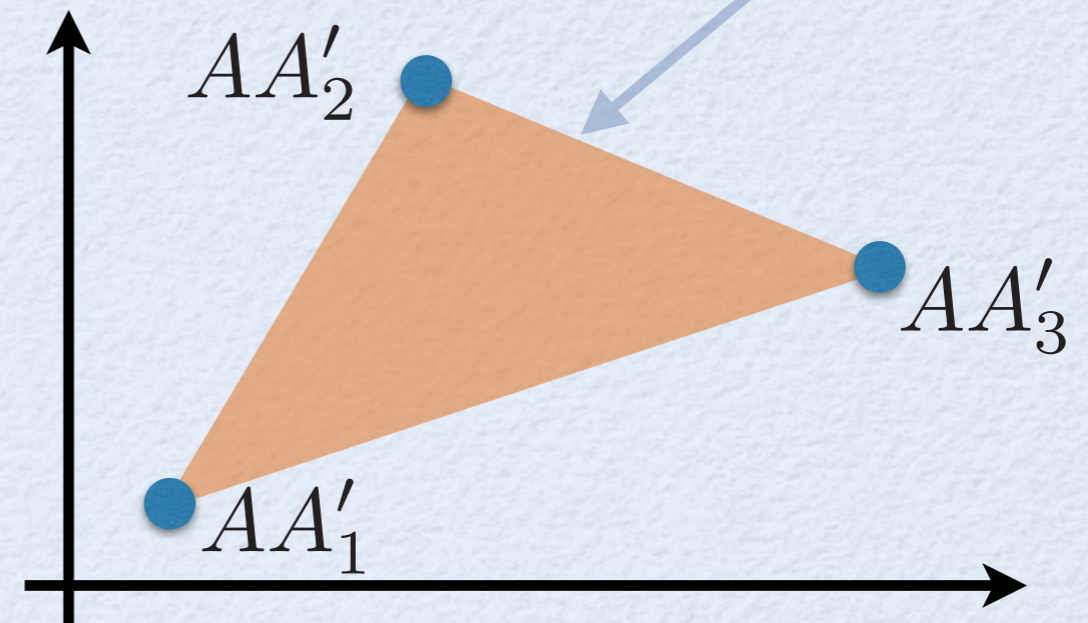
HEAVY ELEMENTS - ISOTOPE SHIFT

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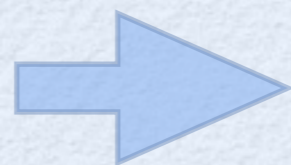


nonlinear King plot from
new physics



similar to data driven background estimation at the LHC

data consistent
with linearity

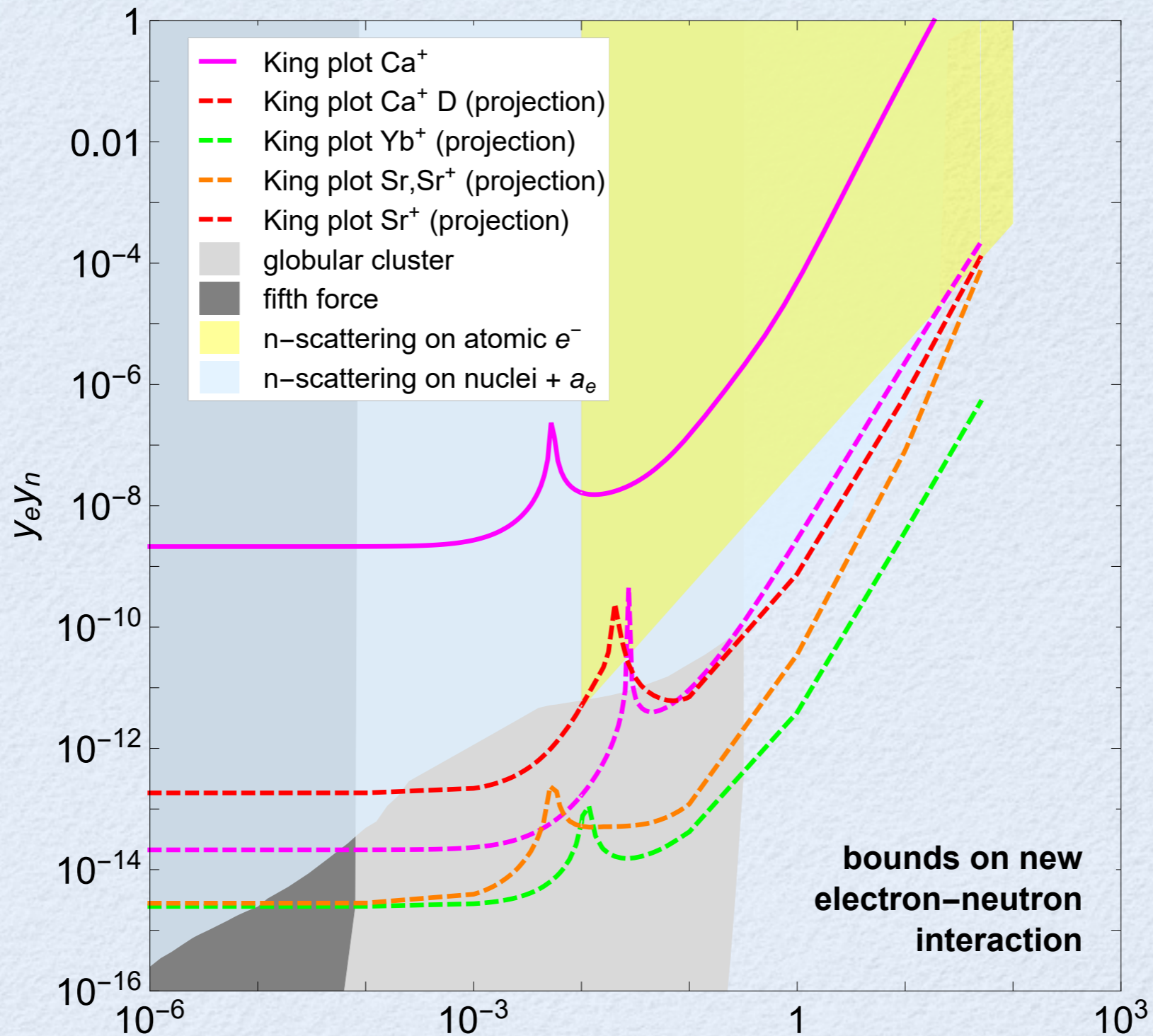


constrain NP

Delaunay, Ozeri, Perez, **YS**, PRD 1601.05087

Berengut, Budker, Delaunay, Flambaum, Frugiuele,
Fuchs, Grojean, Harnik, Ozeri, Perez, **YS**, 1704.05068

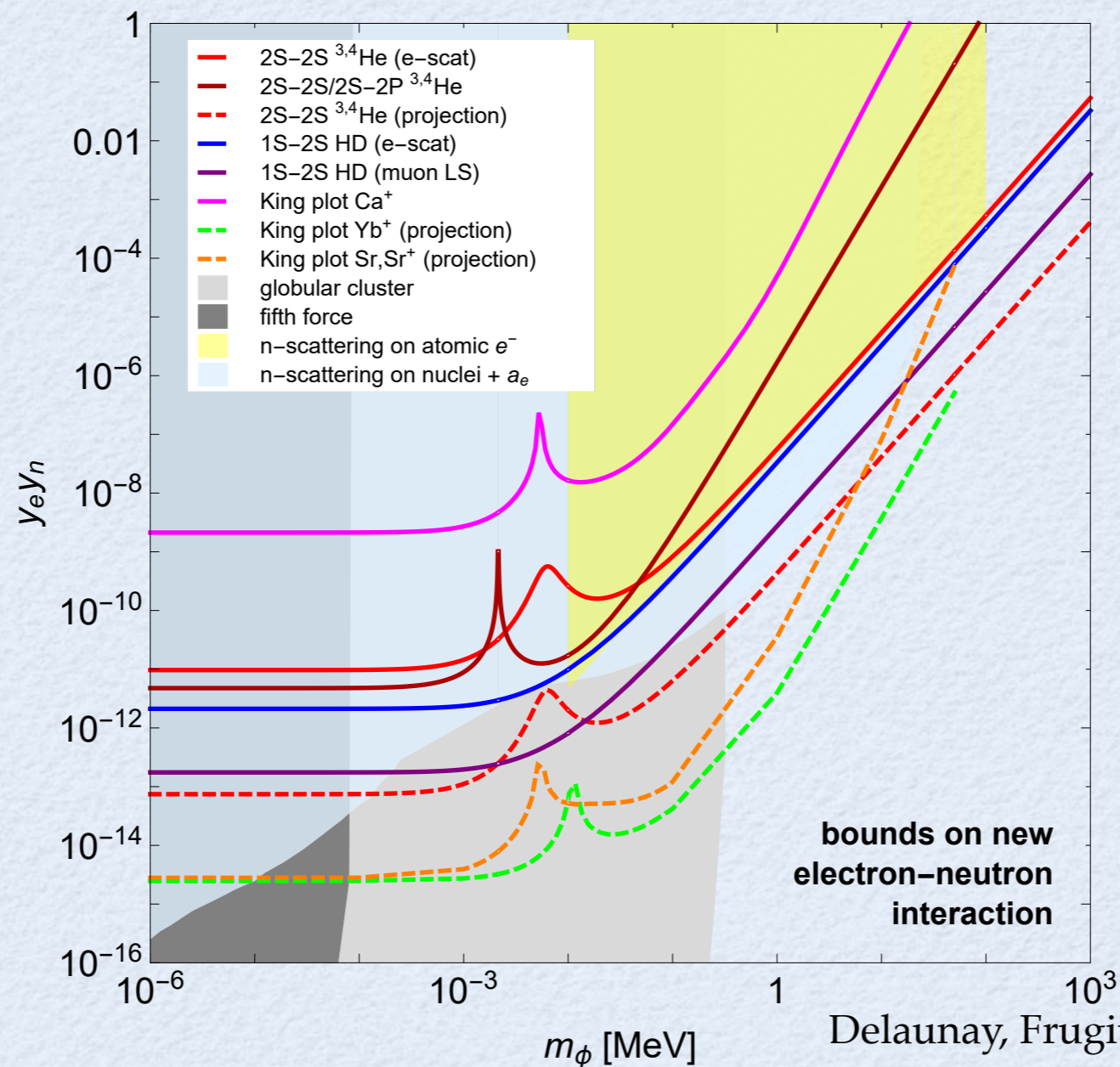
HEAVY ELEMENTS - ISOTOPE SHIFT



m_ϕ [MeV]
21

Berengut, Budker, Delaunay, Flambaum, Frugiuele, Fuchs, Grojean, Harnik, Ozeri, Perez, YS, 1704.05068

ISOTOPE SHIFTS: SUMMARY



Delaunay, Frugieuele, Fuchs, YS, PRD 1709.02817

Berengut, Budker, Delaunay, Flambaum, Frugieuele,

Fuchs, Grojean, Harnik, Ozeri, Perez, YS, 1704.05068

THE QUEST FOR NEW PHYSICS

coupling
to SM



1

10^{-4}

10^{-7}

intensity
frontier



MeV

2×10^{-13} m

GeV

2×10^{-16} m
23

weak scale

2×10^{-18} m

TeV

2×10^{-19} m

new force mass
interaction length

THE QUEST FOR NEW PHYSICS

coupling
to SM



1

10^{-4}

10^{-7}

intensity
frontier



2. *dark photons at LHCb*
3. *ALP at the GeV Scale*

MeV

2×10^{-13} m

GeV

2×10^{-16} m
23

weak scale

2×10^{-18} m

TeV

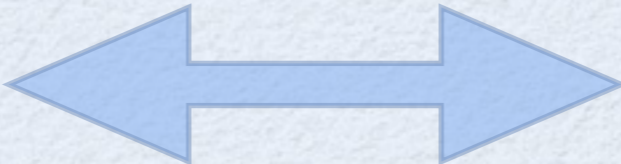
2×10^{-19} m

new force mass

interaction length

dark matter

standard model

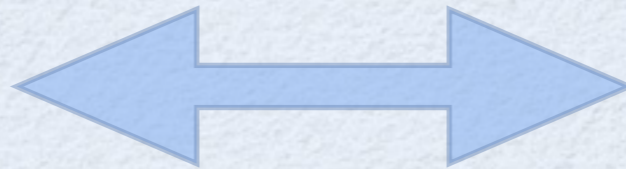


dark matter



a portal

standard model



dark matter



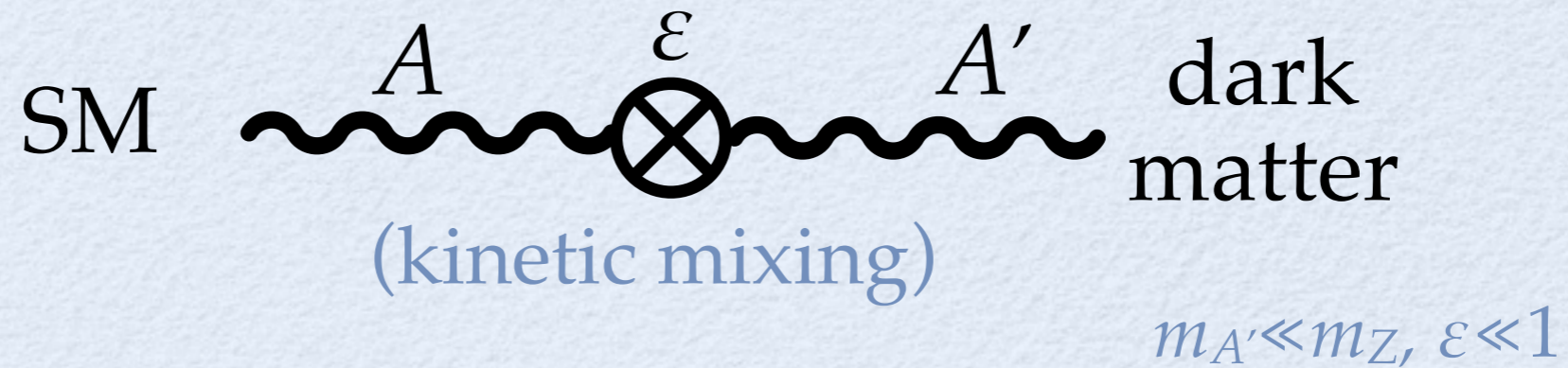
a portal



dark photon - A' / Axion like particles - a
(ALPs)

DARK PHOTON

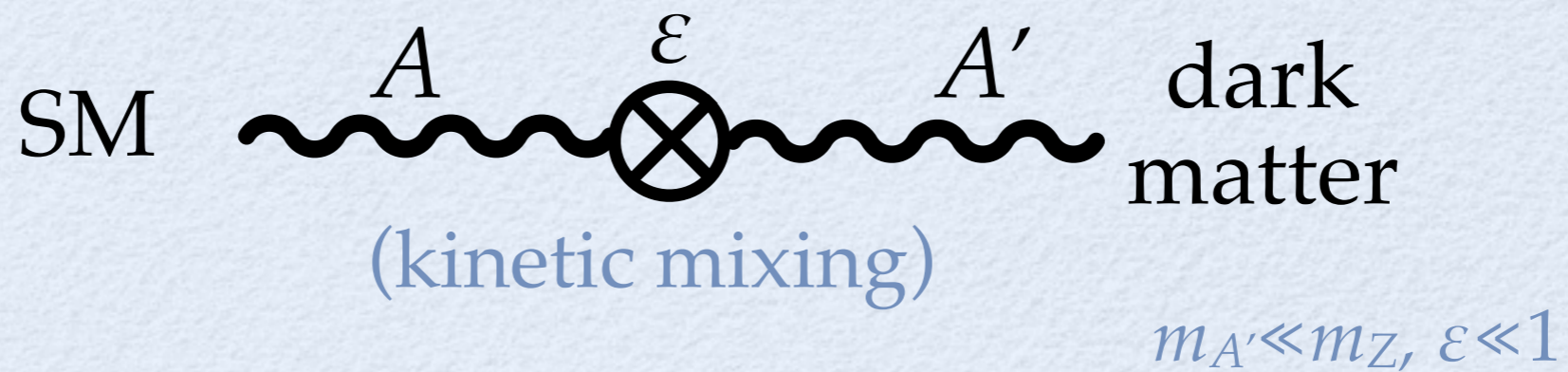
portal between the standard model and dark matter



Holdom, 86'

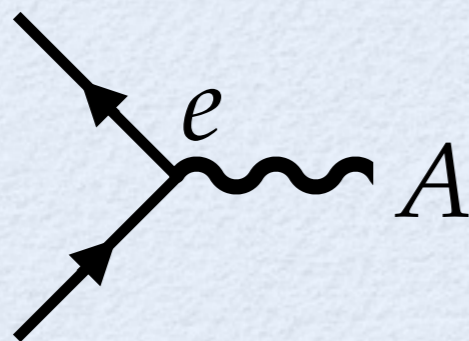
DARK PHOTON

portal between the standard model and dark matter

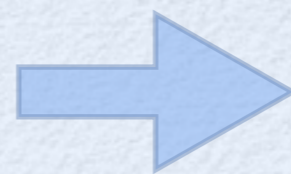


Holdom, 86'

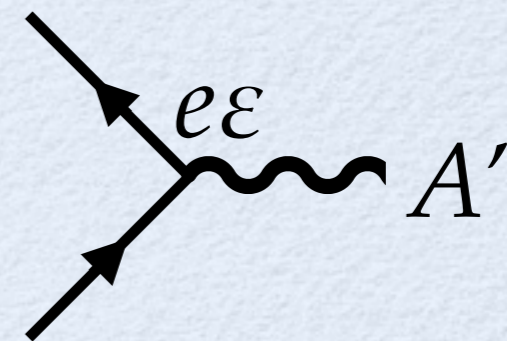
electromagnetic process



$m_A=0$



dark photon process



+ mass effect
($m_{A'} > 0$)

DARK PHOTON SIGNAL

differential relation:

$$\frac{d\sigma_{pp \rightarrow X A' \rightarrow X \mu^+ \mu^-}}{d\sigma_{pp \rightarrow X \gamma^* \rightarrow X \mu^+ \mu^-}} = \epsilon^4 \frac{m_{\mu\mu}^4}{(m_{\mu\mu}^2 - m_{A'}^2)^2 + \Gamma_{A'}^2 m_{A'}^2}$$

DARK PHOTON SIGNAL

differential relation:

$$\frac{d\sigma_{pp \rightarrow X A' \rightarrow X \mu^+ \mu^-}}{d\sigma_{pp \rightarrow X \gamma^* \rightarrow X \mu^+ \mu^-}} = \epsilon^4 \frac{m_{\mu\mu}^4}{(m_{\mu\mu}^2 - m_{A'}^2)^2 + \Gamma_{A'}^2 m_{A'}^2}$$

per mass bin:

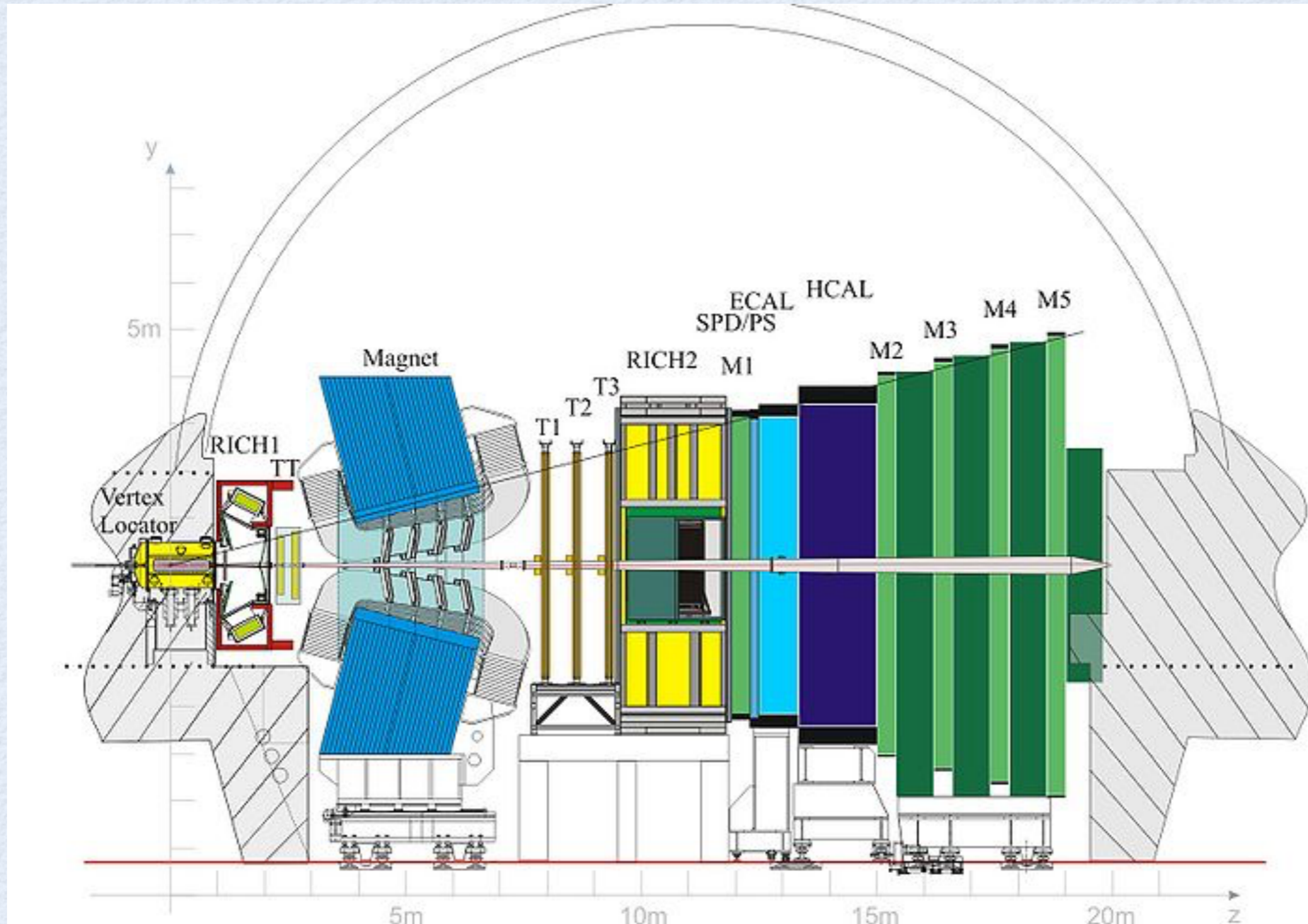
$$\frac{S}{B_{\text{EM}}} \approx \epsilon^4 \frac{\pi}{8} \frac{m_{A'}^2}{\Gamma_{A'} \sigma_{m_{\mu\mu}}} \approx \frac{3\pi}{8} \frac{m_{A'}}{\sigma_{m_{\mu\mu}}} \frac{\epsilon^2}{\alpha_{\text{EM}} (N_\ell + \mathcal{R}_\mu)}$$

$$\frac{\Gamma_{A'}}{m_{A'}} \approx \frac{\epsilon^2 \alpha_{\text{EM}}}{3} (N_\ell + \mathcal{R}_\mu)$$

number of leptons with
mass below $m_{A'}/2$

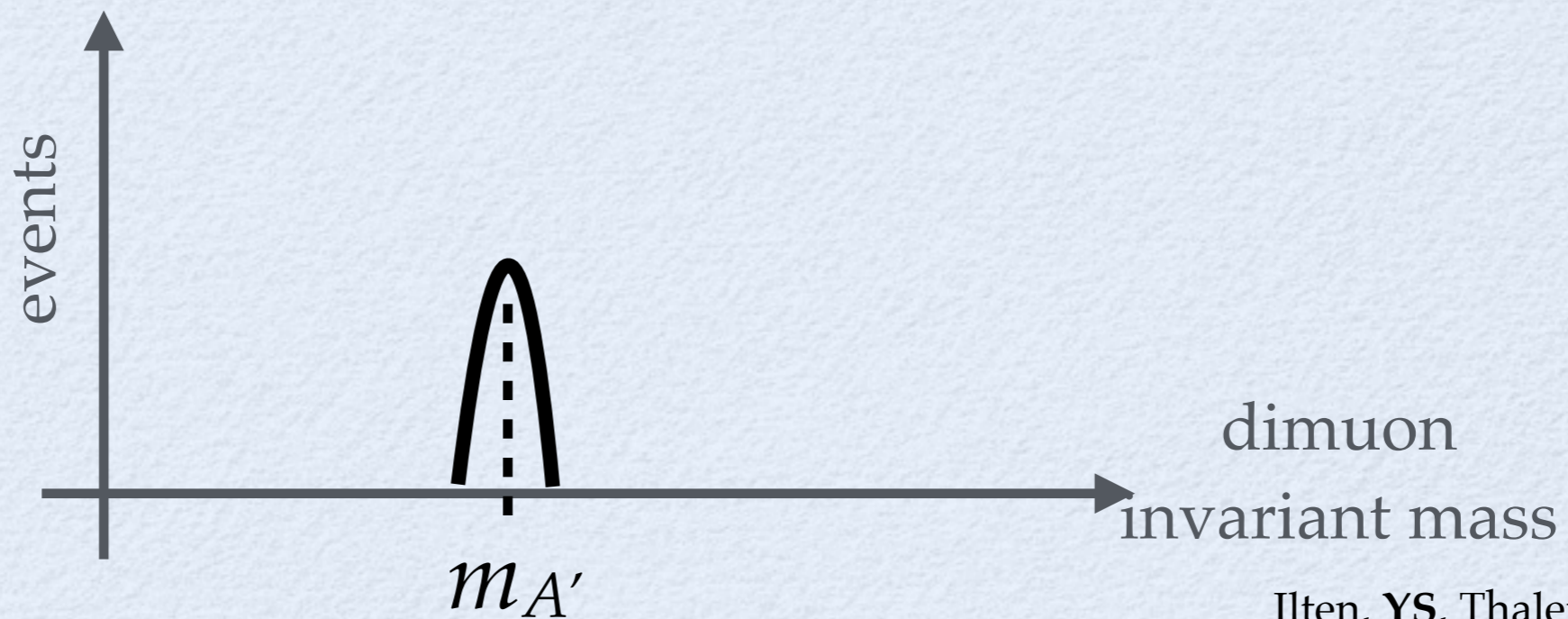
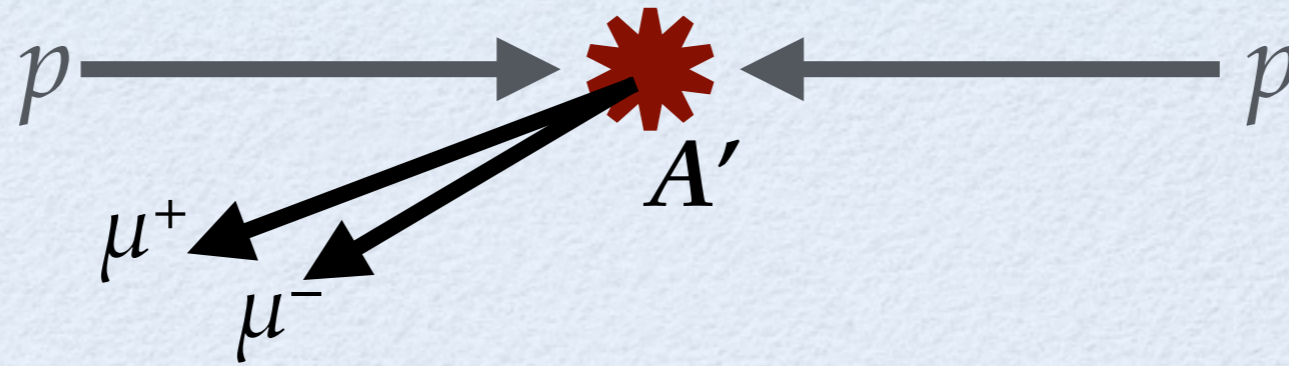
$$\frac{\sigma_{e^+e^- \rightarrow \text{hadrons}}}{\sigma_{e^+e^- \rightarrow \mu^+\mu^-}}$$

DARK PHOTON AT LHCb



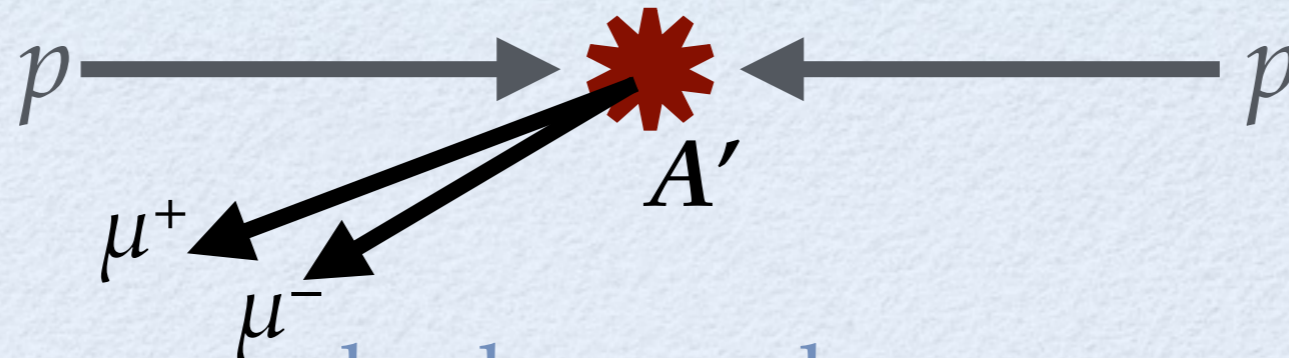
SIGNAL AND BACKGROUNDS

prompt
short lifetime
larger ϵ



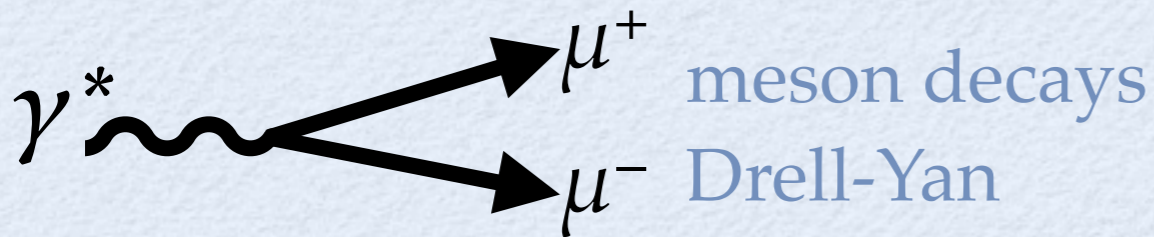
SIGNAL AND BACKGROUNDS

prompt
short lifetime
larger ϵ

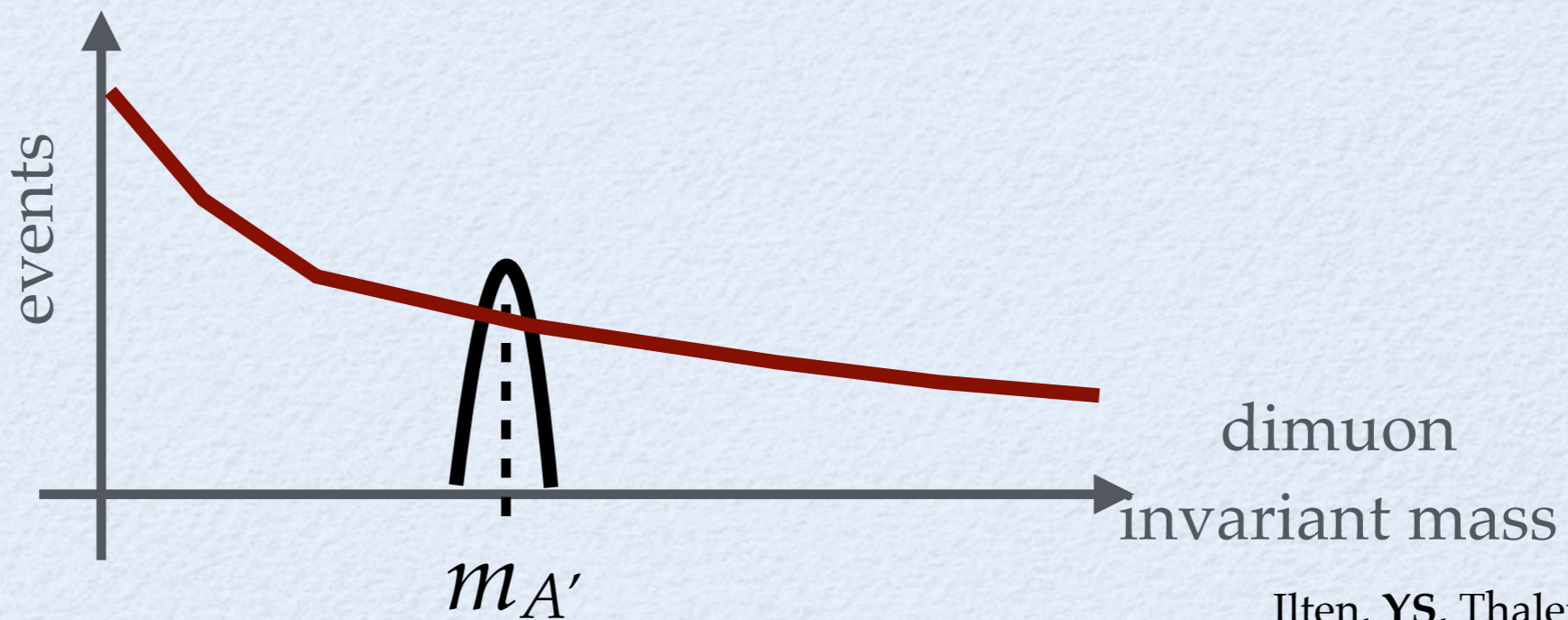


“good” backgrounds

$$S_{A'} \propto B_{EM}$$



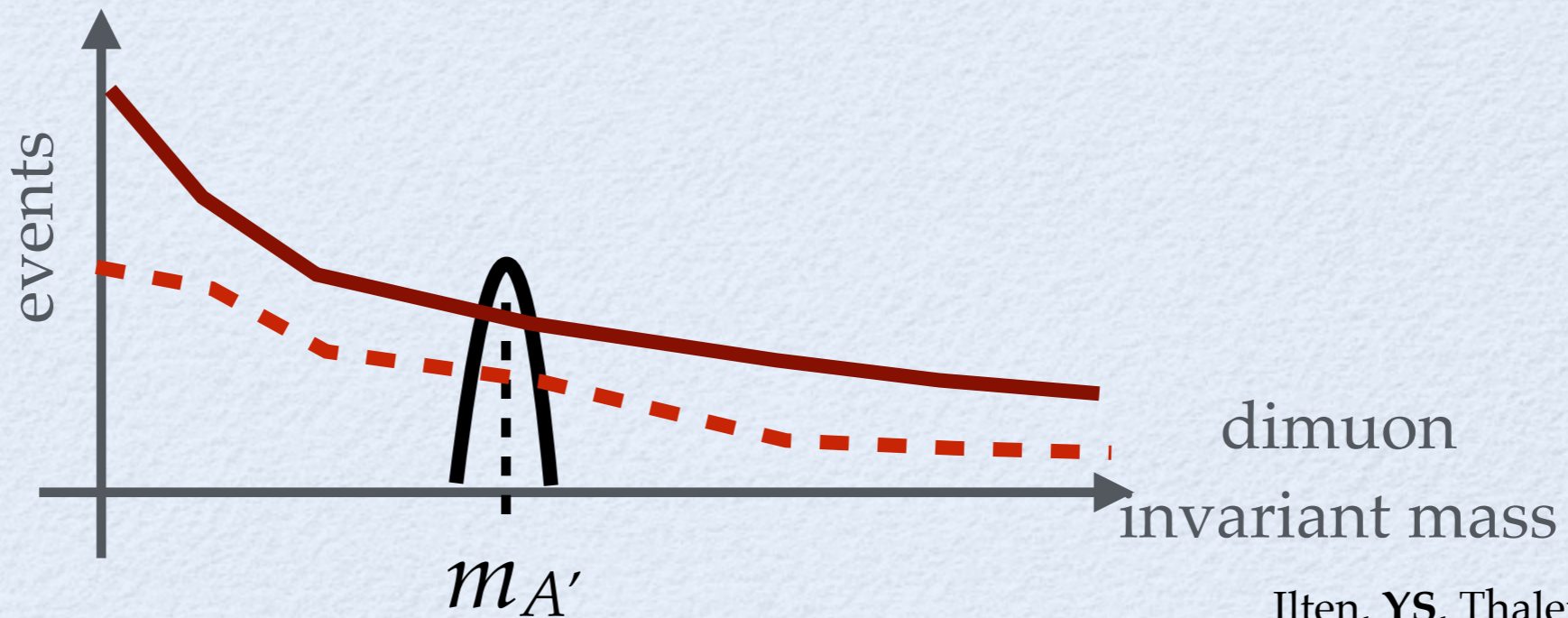
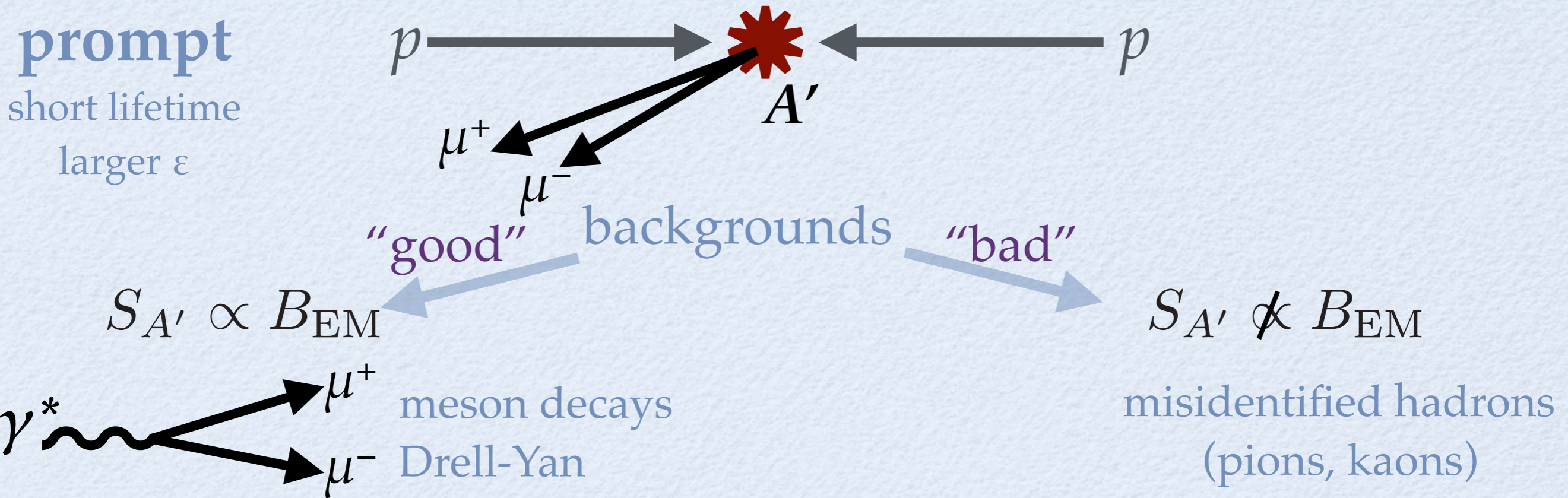
meson decays
Drell-Yan



SIGNAL AND BACKGROUNDS

prompt

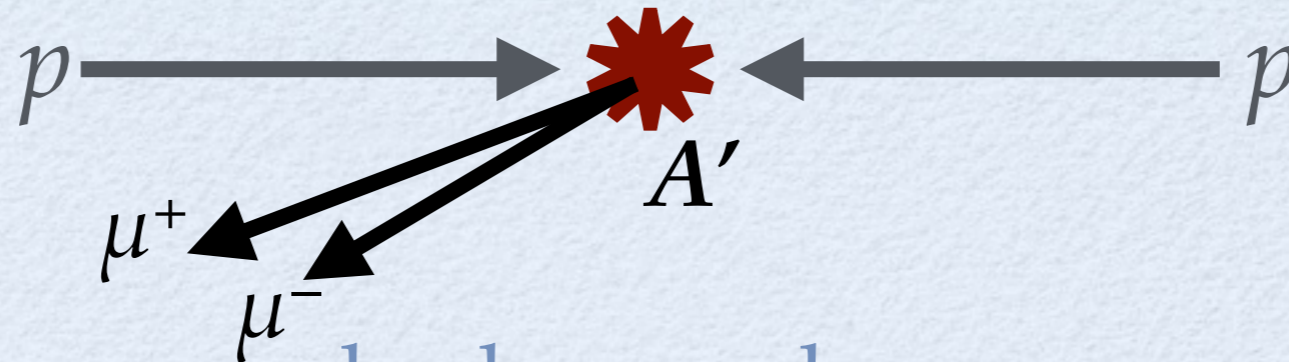
short lifetime
larger ϵ



SIGNAL AND BACKGROUNDS

prompt

short lifetime
larger ϵ

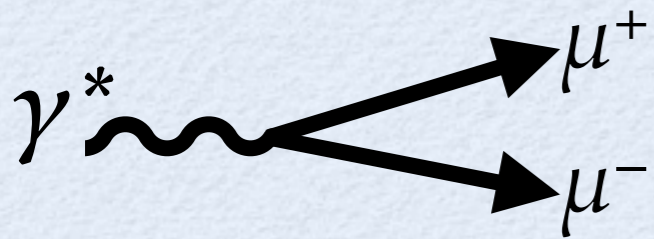


“good” backgrounds

“bad”

$$S_{A'} \propto B_{EM}$$

$$S_{A'} \not\propto B_{EM}$$

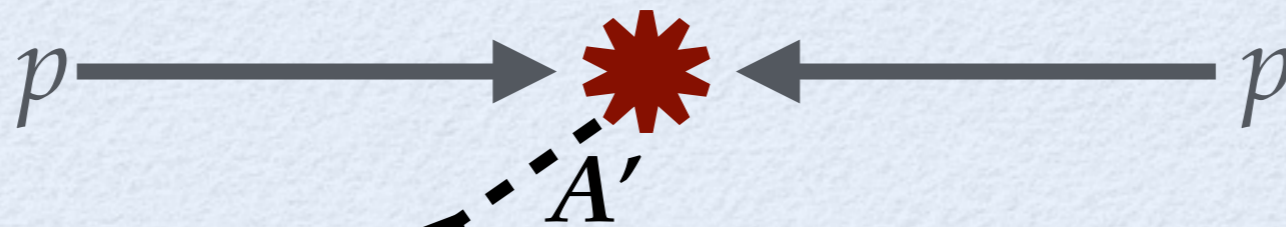


meson decays
Drell-Yan

misidentified hadrons
(pions, kaons)

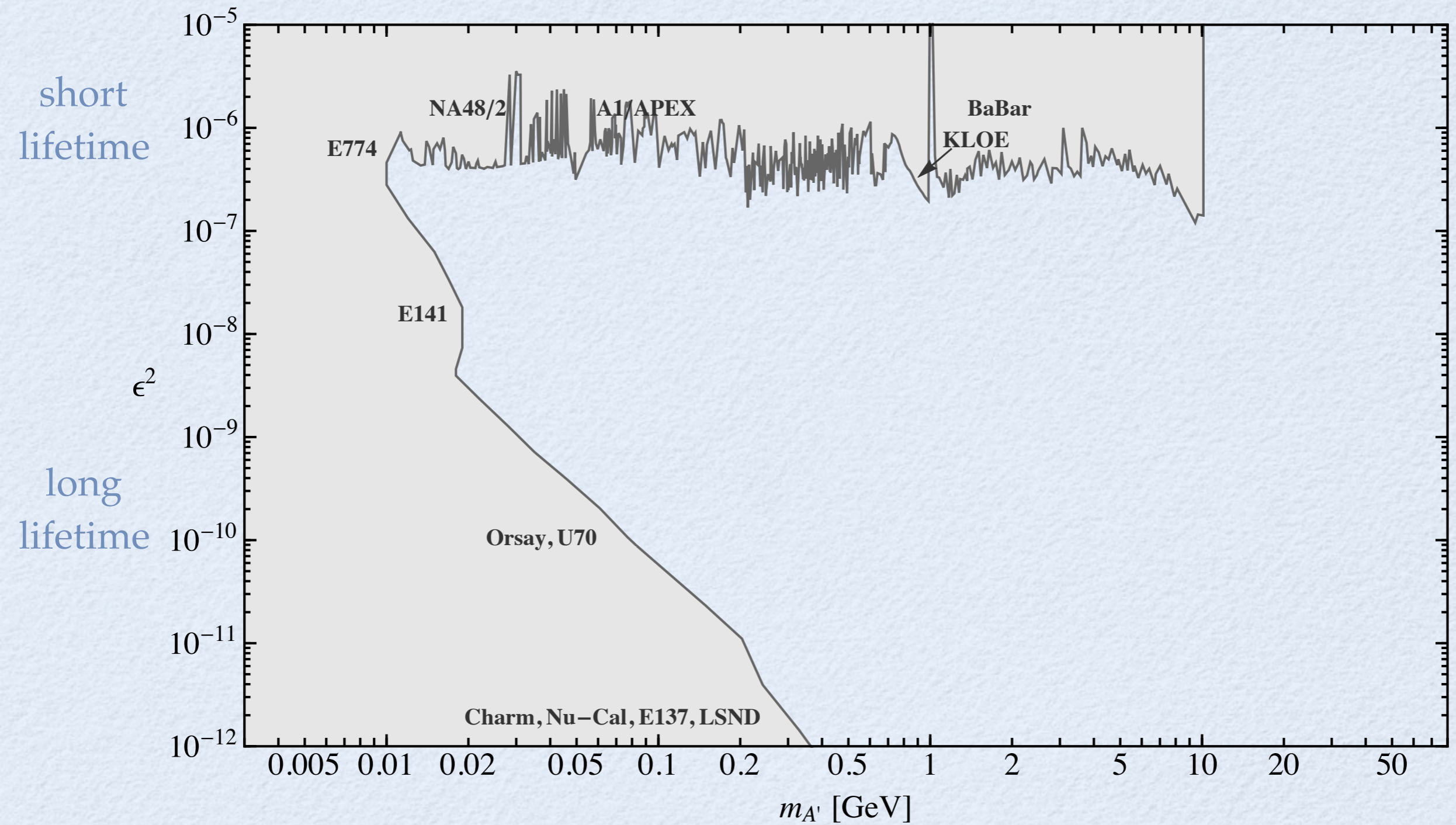
displaced

long lifetime
smaller ϵ

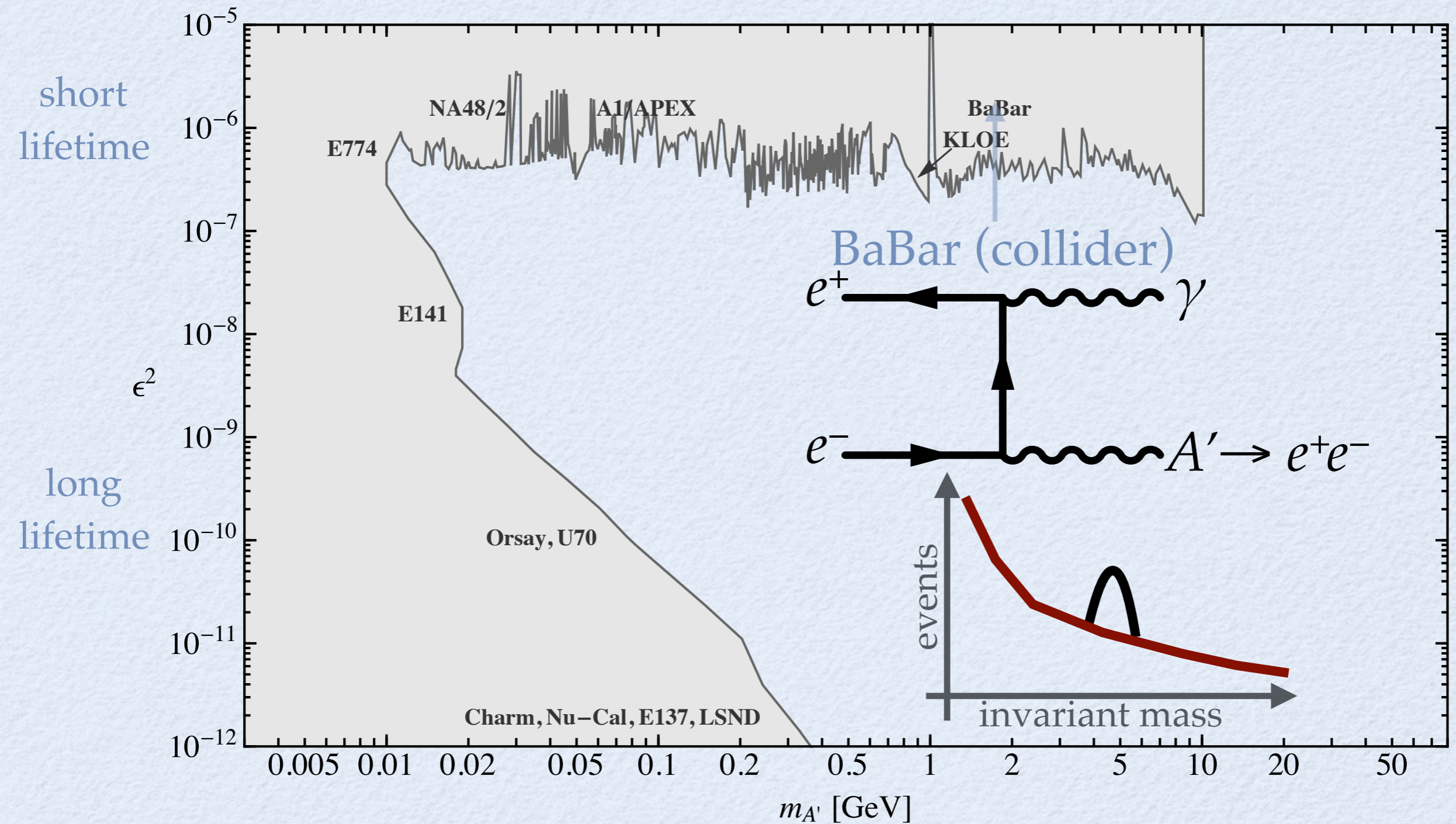


backgrounds
heavy flavor (b, c) decays

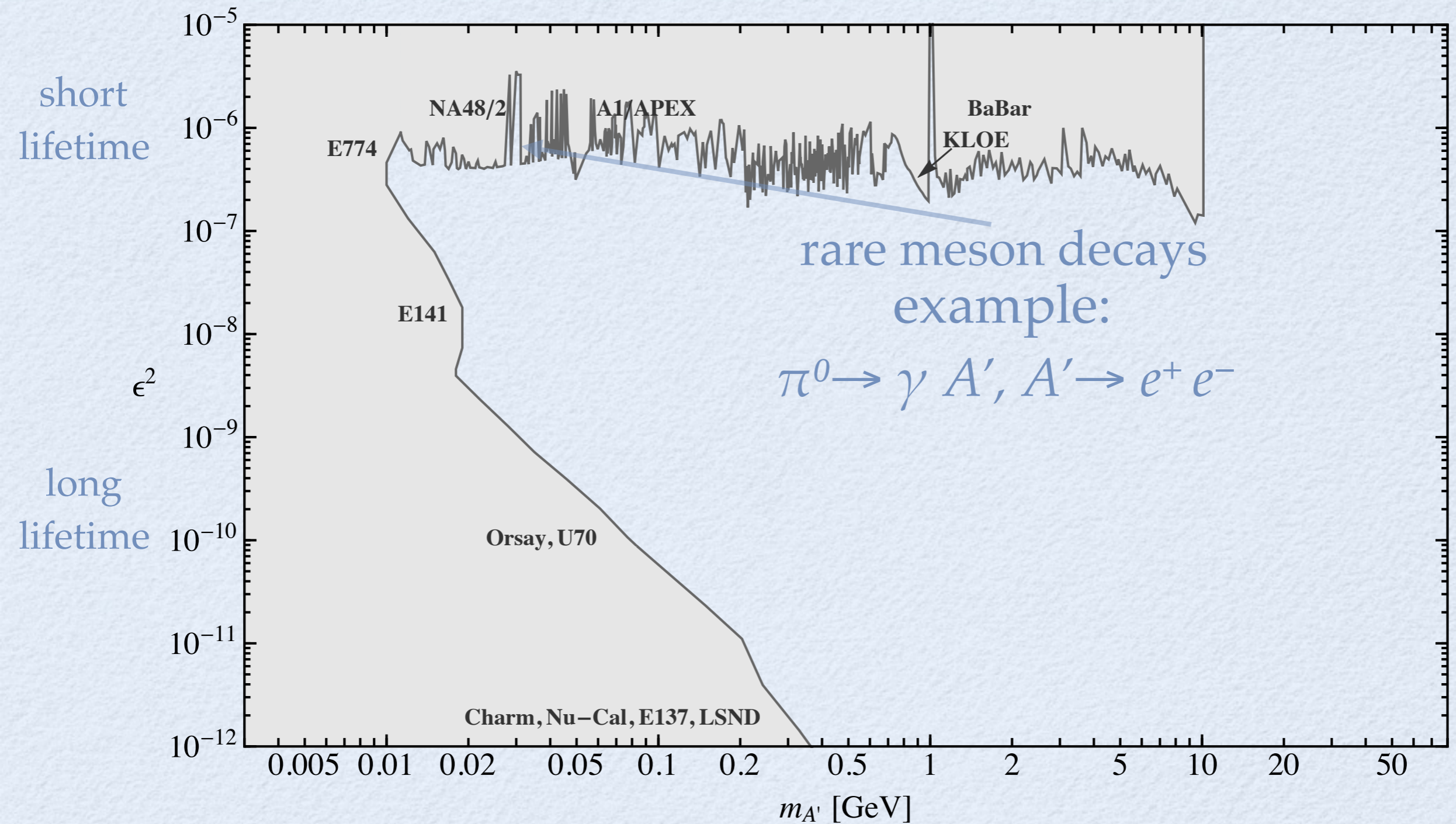
BOUNDS ON A'



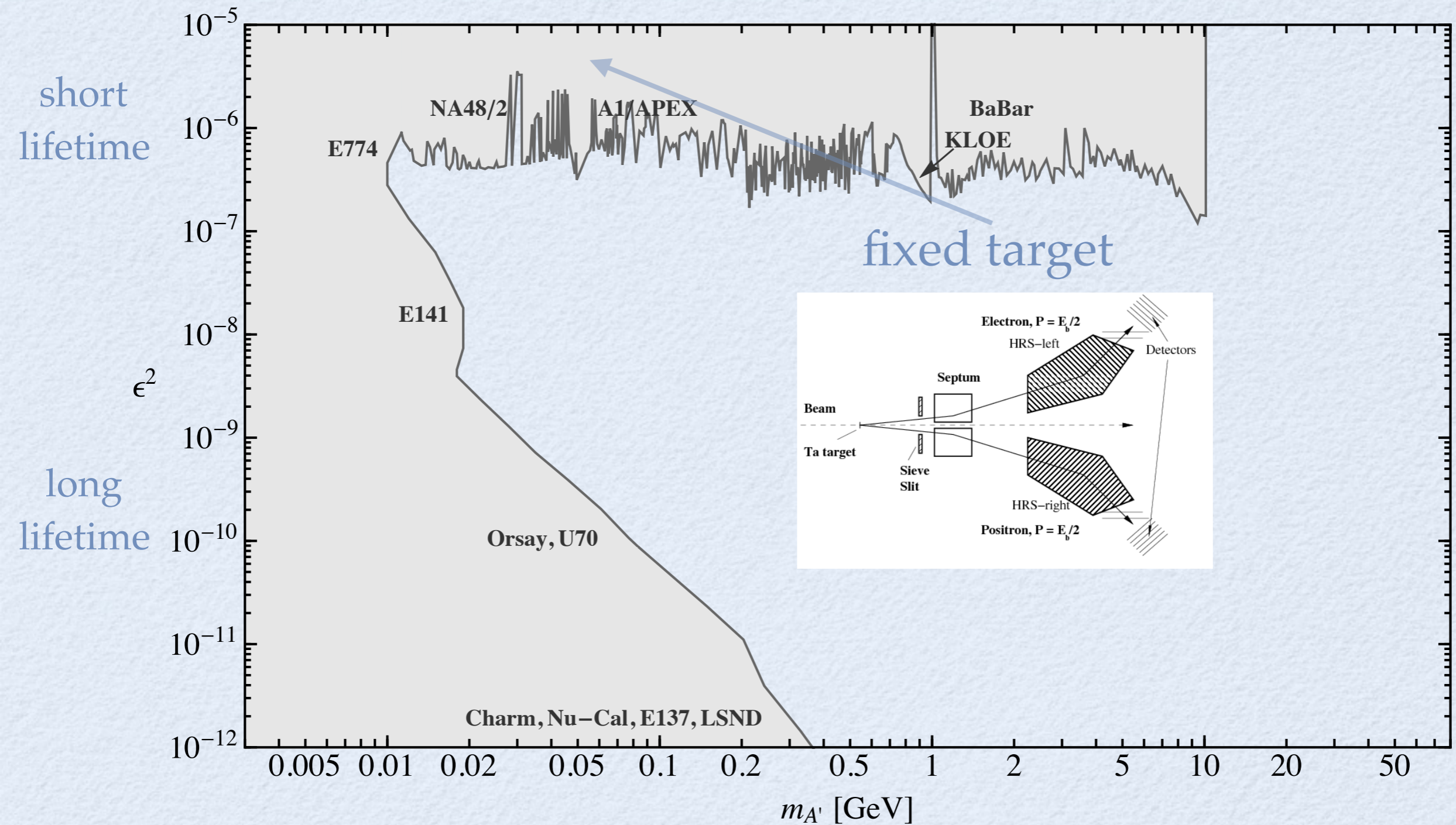
BOUNDS ON A'



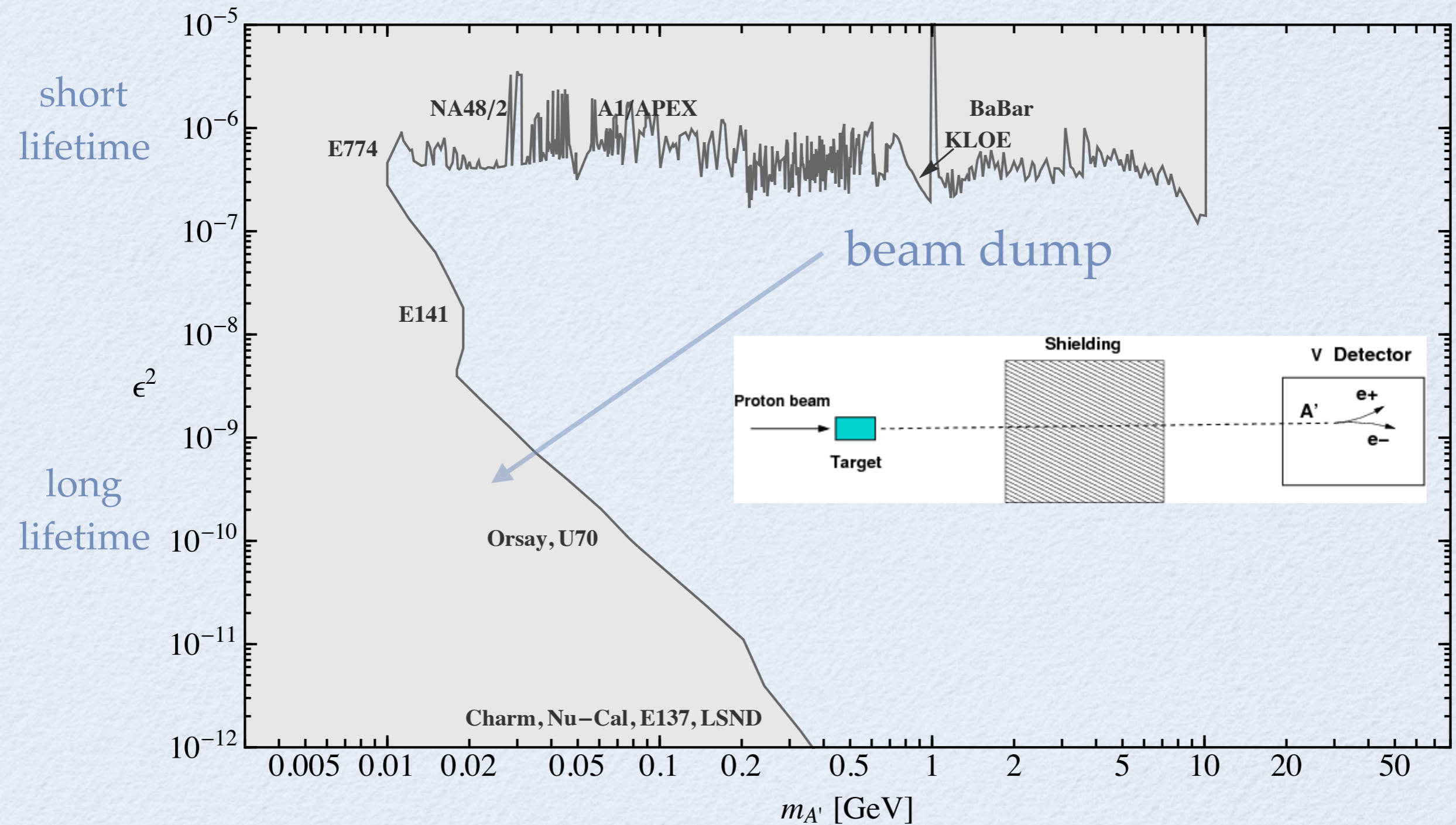
BOUNDS ON A'



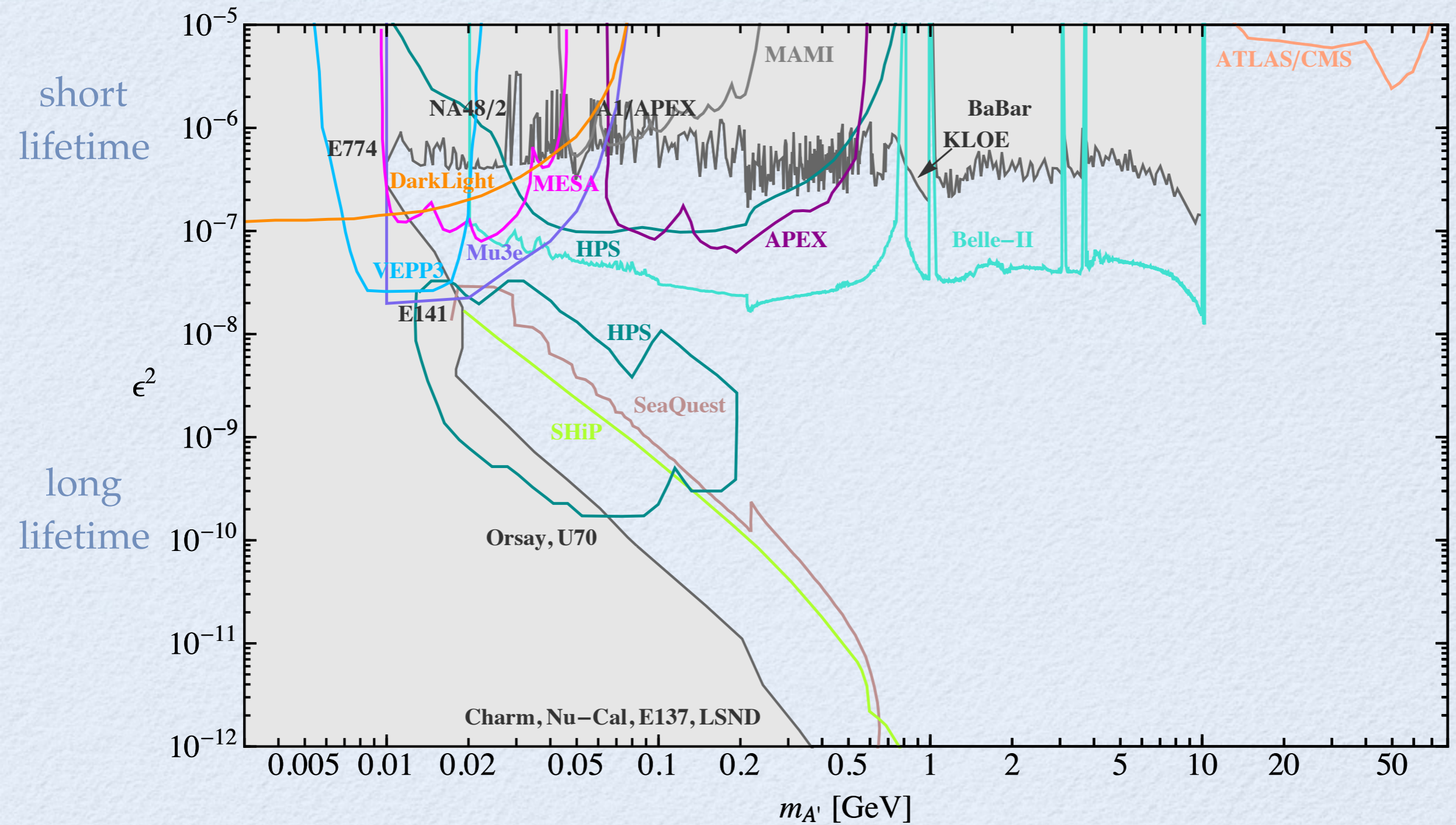
BOUNDS ON A'



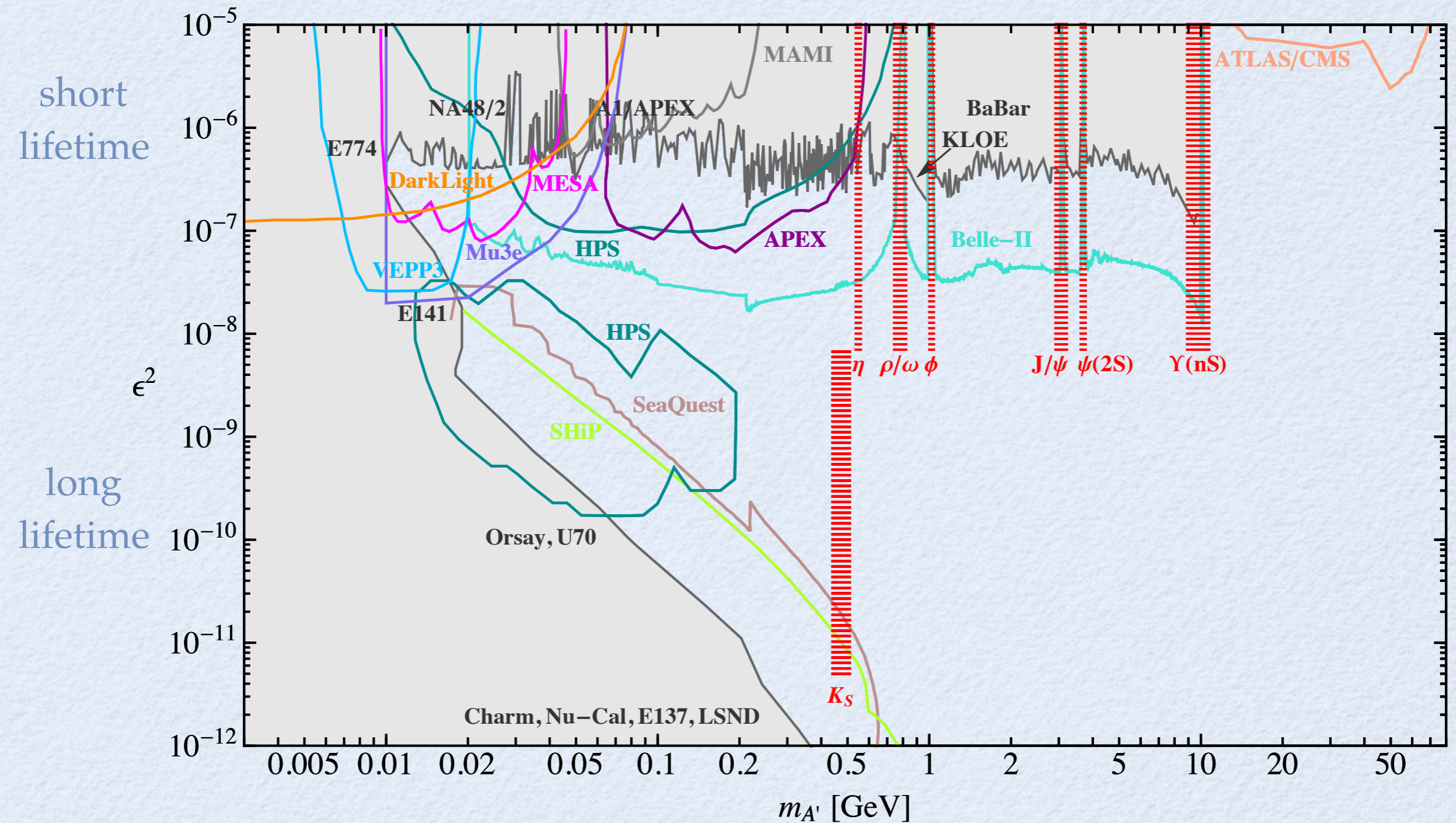
BOUNDS ON A'



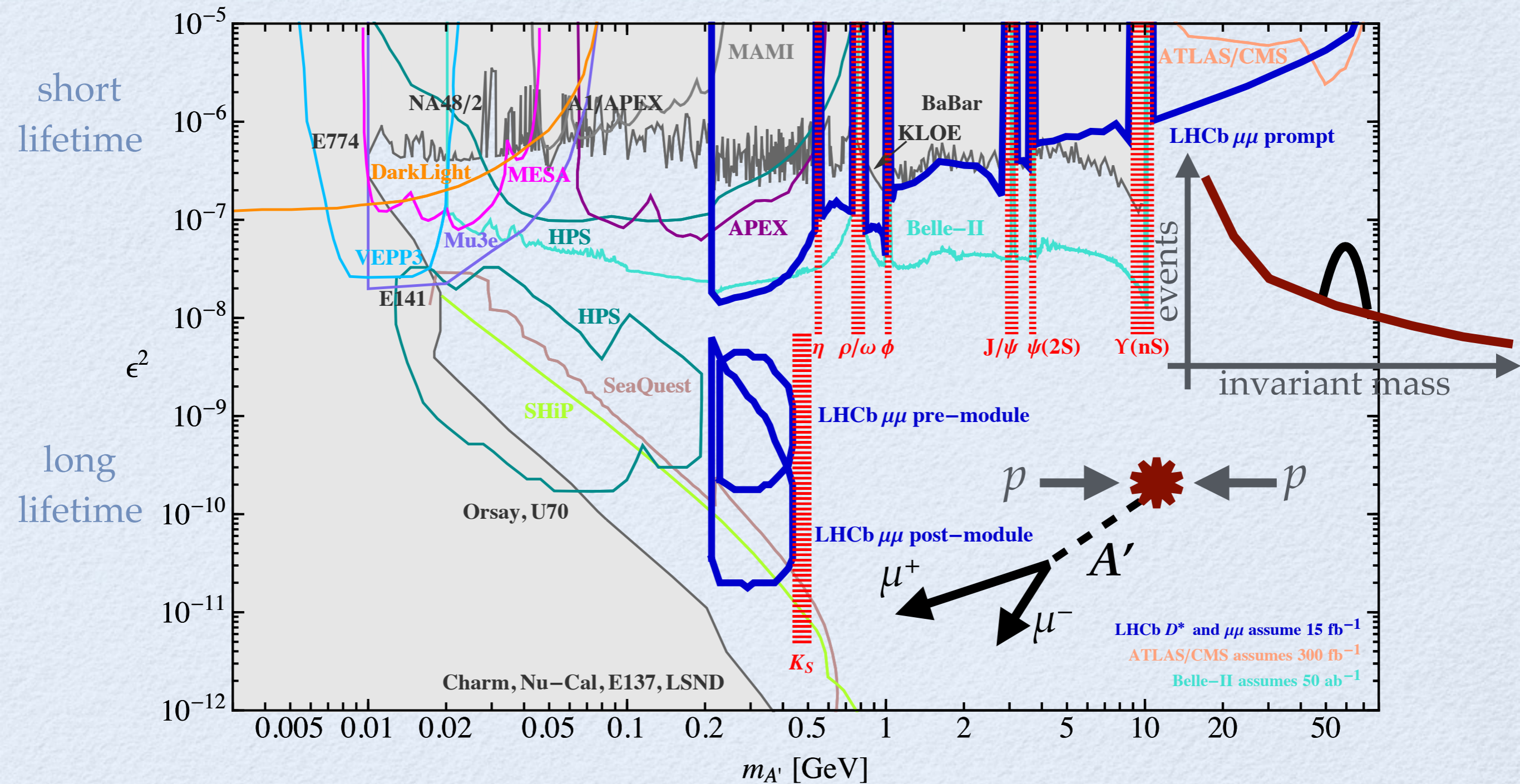
BOUNDS ON A'



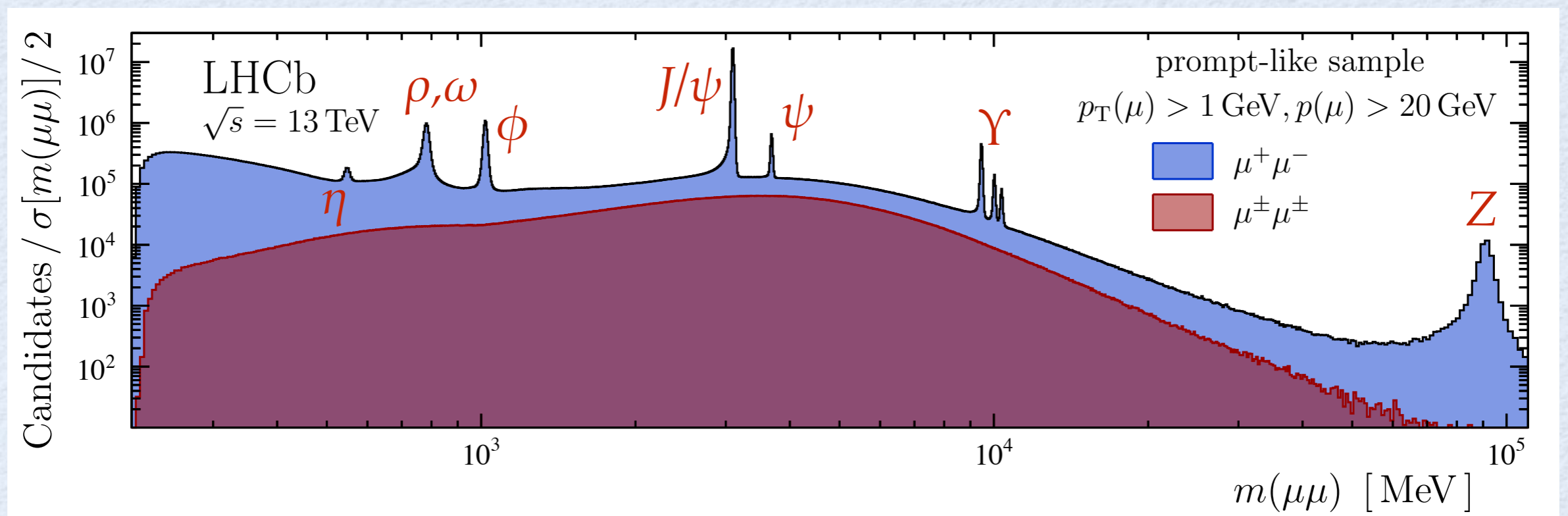
LHCb DARK PHOTON REACH



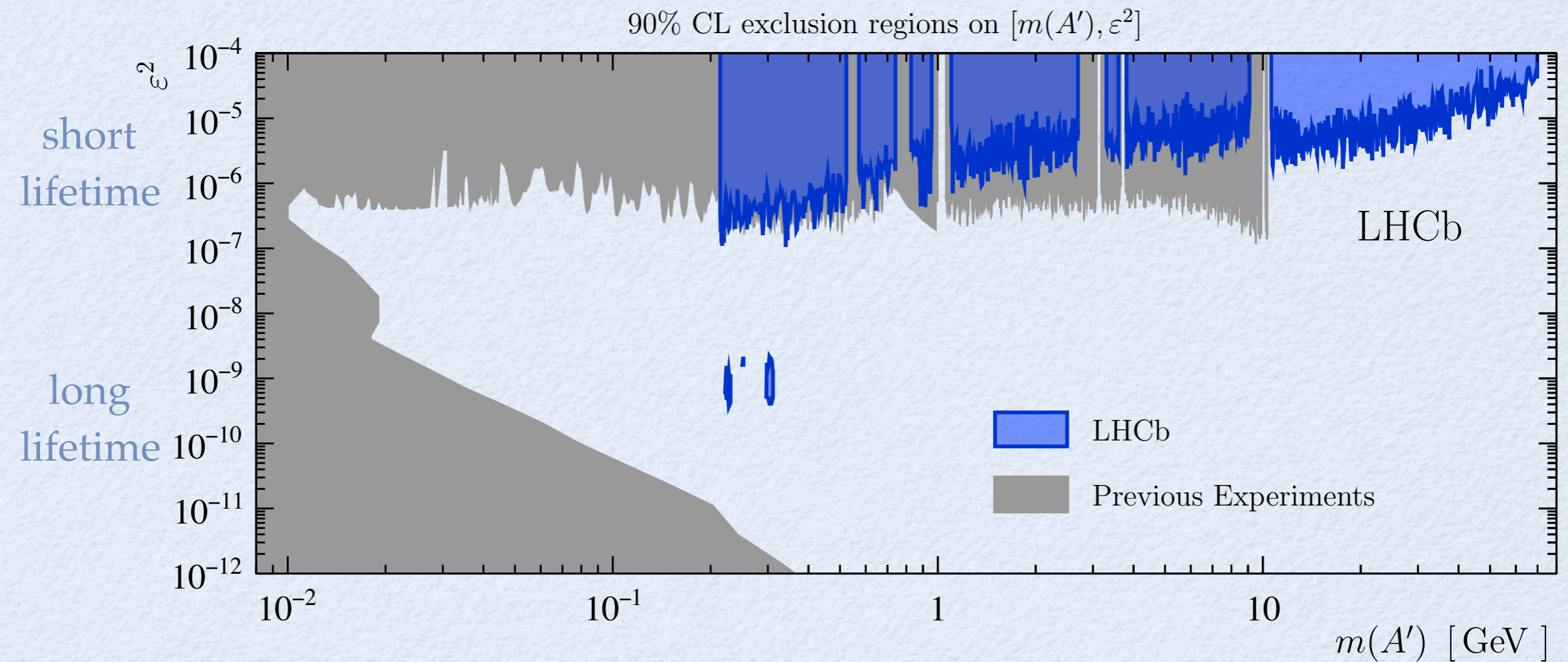
LHCb DARK PHOTON REACH



LHCb RESULT WITH 2016 DATA

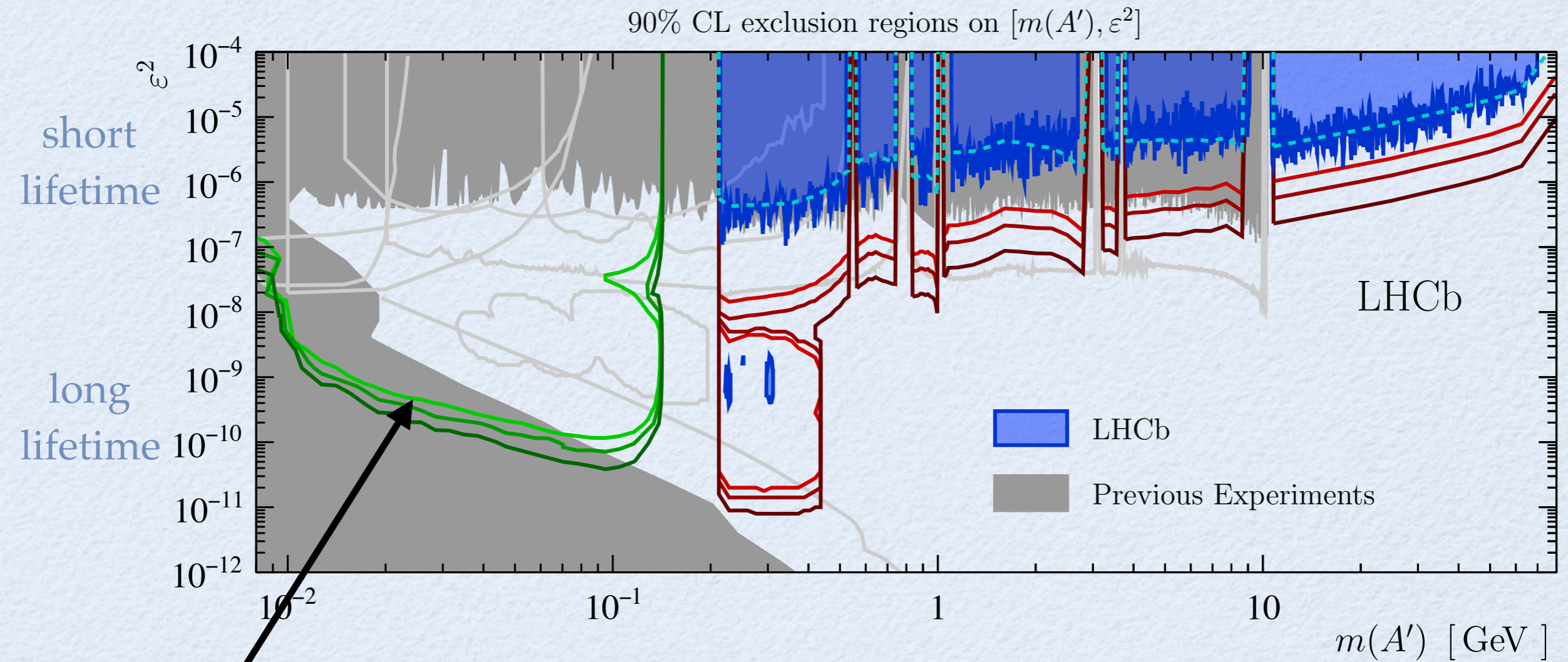


LHCb RESULT WITH 2016 DATA



LHCb, 1710.02867

LHCb RESULT WITH 2016 DATA

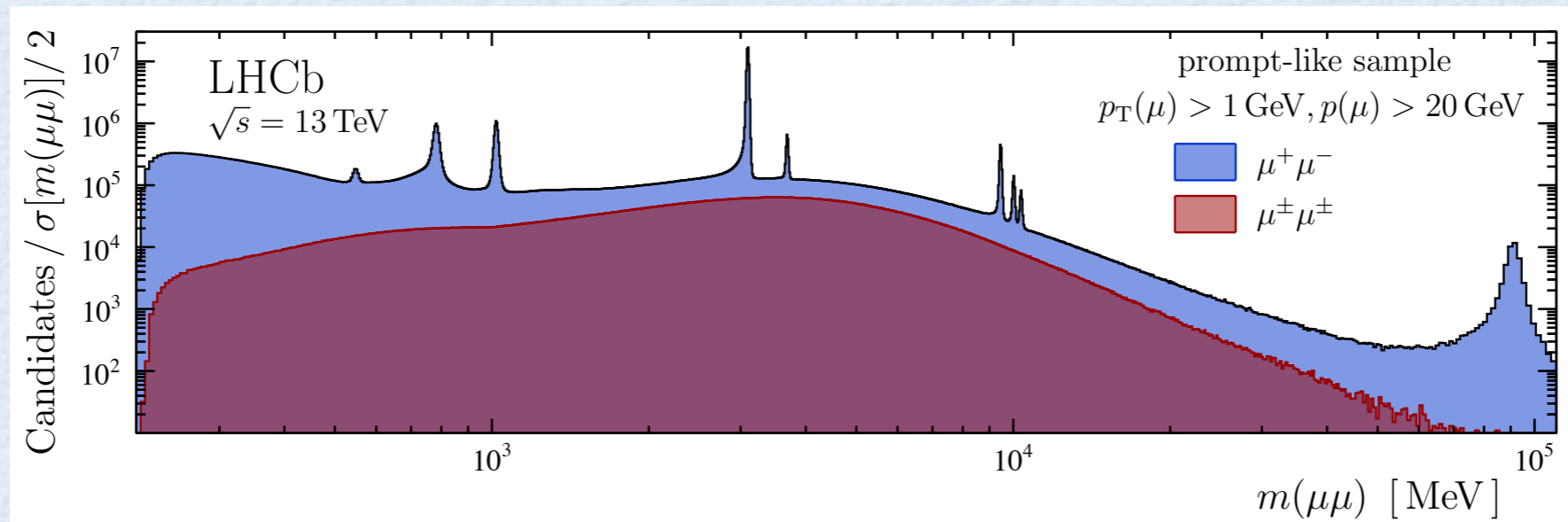


$D^* \rightarrow DA', A' \rightarrow e^+e^-$

Ilten, Thaler, Williams, Xue,
1509.06765

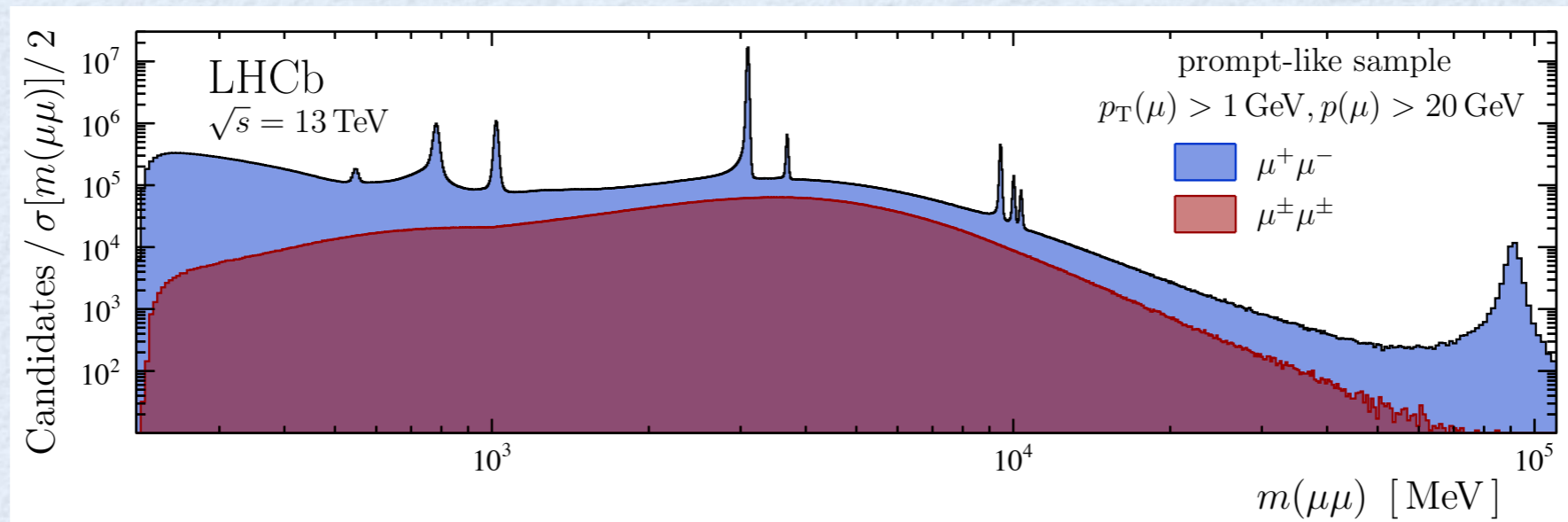
LHCb, 1710.02867

BEYOND THE DARK PHOTON MODEL



LHCb, 1710.02867

BEYOND THE DARK PHOTON MODEL

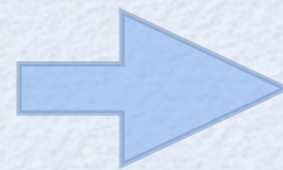


LHCb, 1710.02867

generic vector
resonances

$$\mathcal{L} \subset g_X \sum_f x_f \bar{f} \gamma^\mu f X_\mu + \sum_\chi \mathcal{L}_{X\chi\bar{\chi}}$$

rescaling of the production
and branching ratio



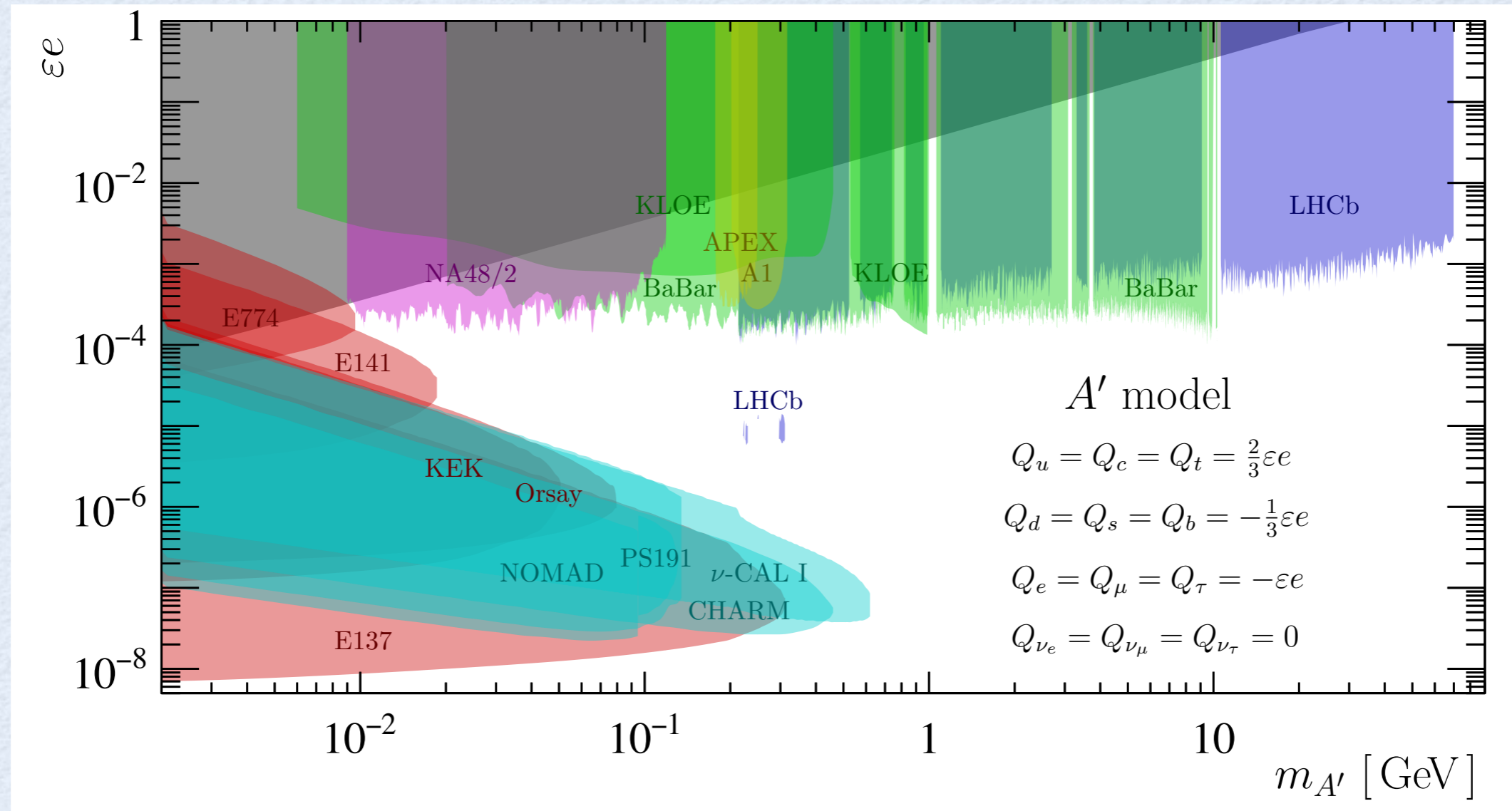
constrain generic vector
resonances - X

$$\sigma_X \mathcal{B}_{X \rightarrow \mathcal{F}} \epsilon(\tau_X) = \sigma_{A'} \mathcal{B}_{A' \rightarrow \mathcal{F}} \epsilon(\tau_{A'})$$

$$g_X = \epsilon^2 \frac{\bar{\sigma}_{A'} \mathcal{B}_{A' \rightarrow \mathcal{F}} \epsilon(\tau_{A'})}{\bar{\sigma}_X \mathcal{B}_{X \rightarrow \mathcal{F}} \epsilon(\tau_X)}$$

see also <https://gitlab.com/philten/darkcast>

BEYOND THE DARK PHOTON MODEL



electron beam dump

proton beam dump

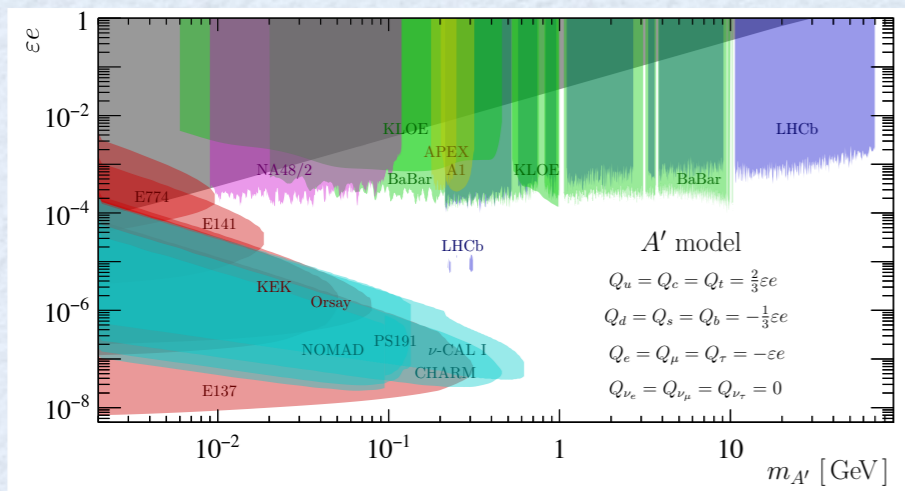
e^+e^- collider

pp collider

meson decays

e on fixed target

BEYOND THE DARK PHOTON MODEL



electron beam dump

proton beam dump

e^+e^- collider

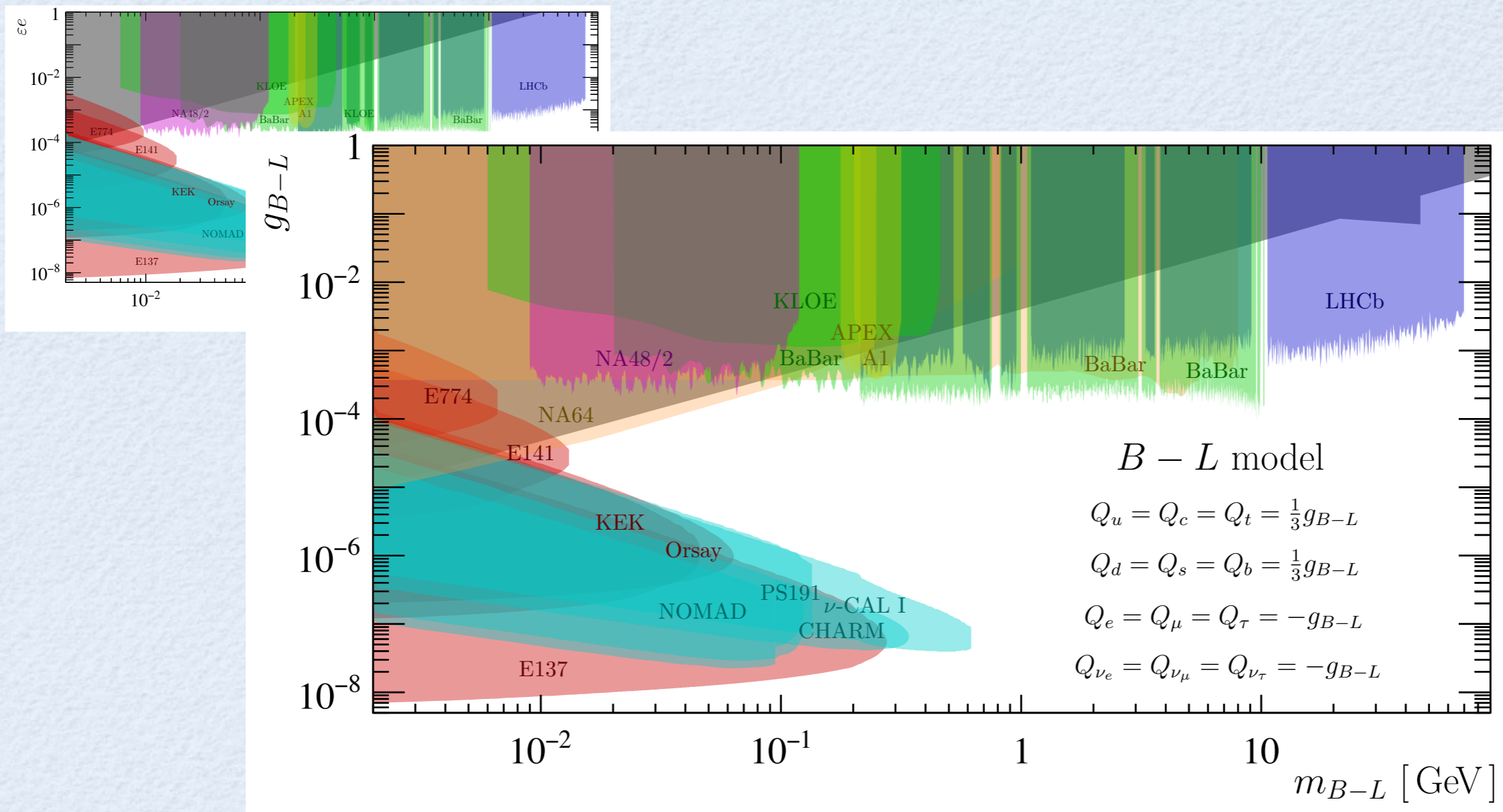
pp collider

meson decays

e on fixed target

invisible decays

BEYOND THE DARK PHOTON MODEL



electron beam dump

proton beam dump

e^+e^- collider

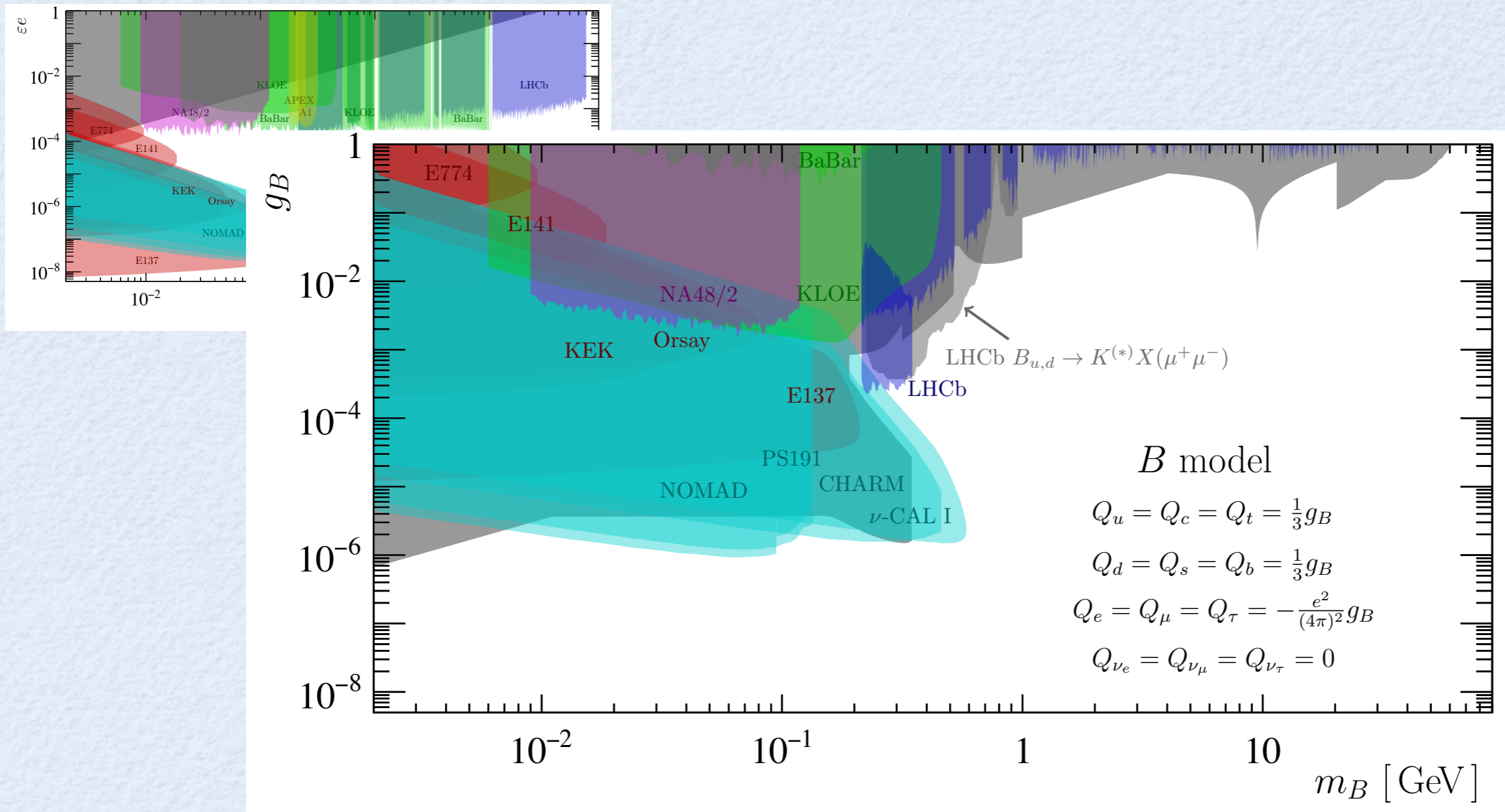
pp collider

meson decays

e on fixed target

invisible decays

BEYOND THE DARK PHOTON MODEL



electron beam dump

proton beam dump

e^+e^- collider

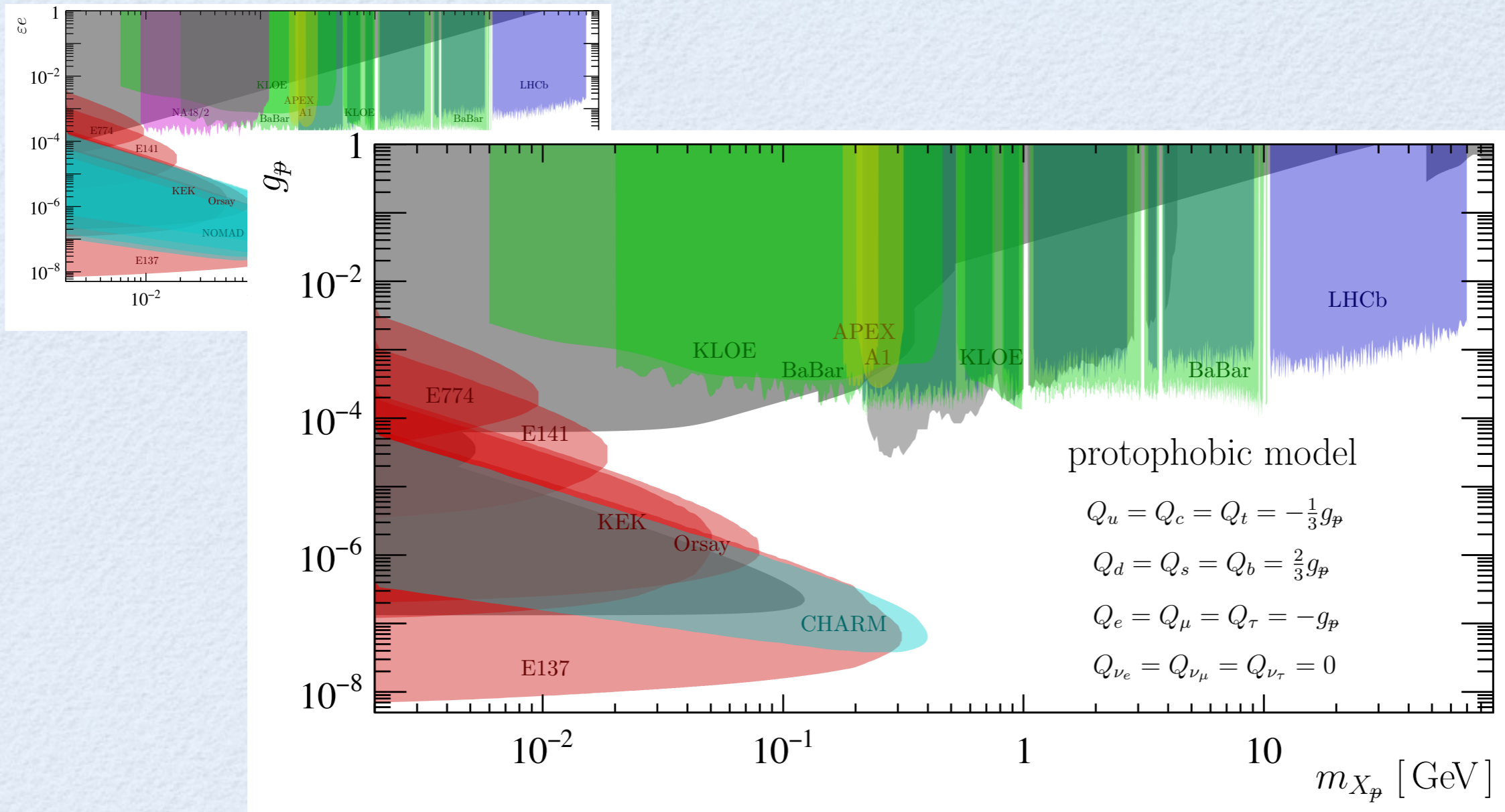
pp collider

meson decays

e on fixed target

invisible decays

BEYOND THE DARK PHOTON MODEL



electron beam dump

proton beam dump

e^+e^- collider

pp collider

meson decays

e on fixed target

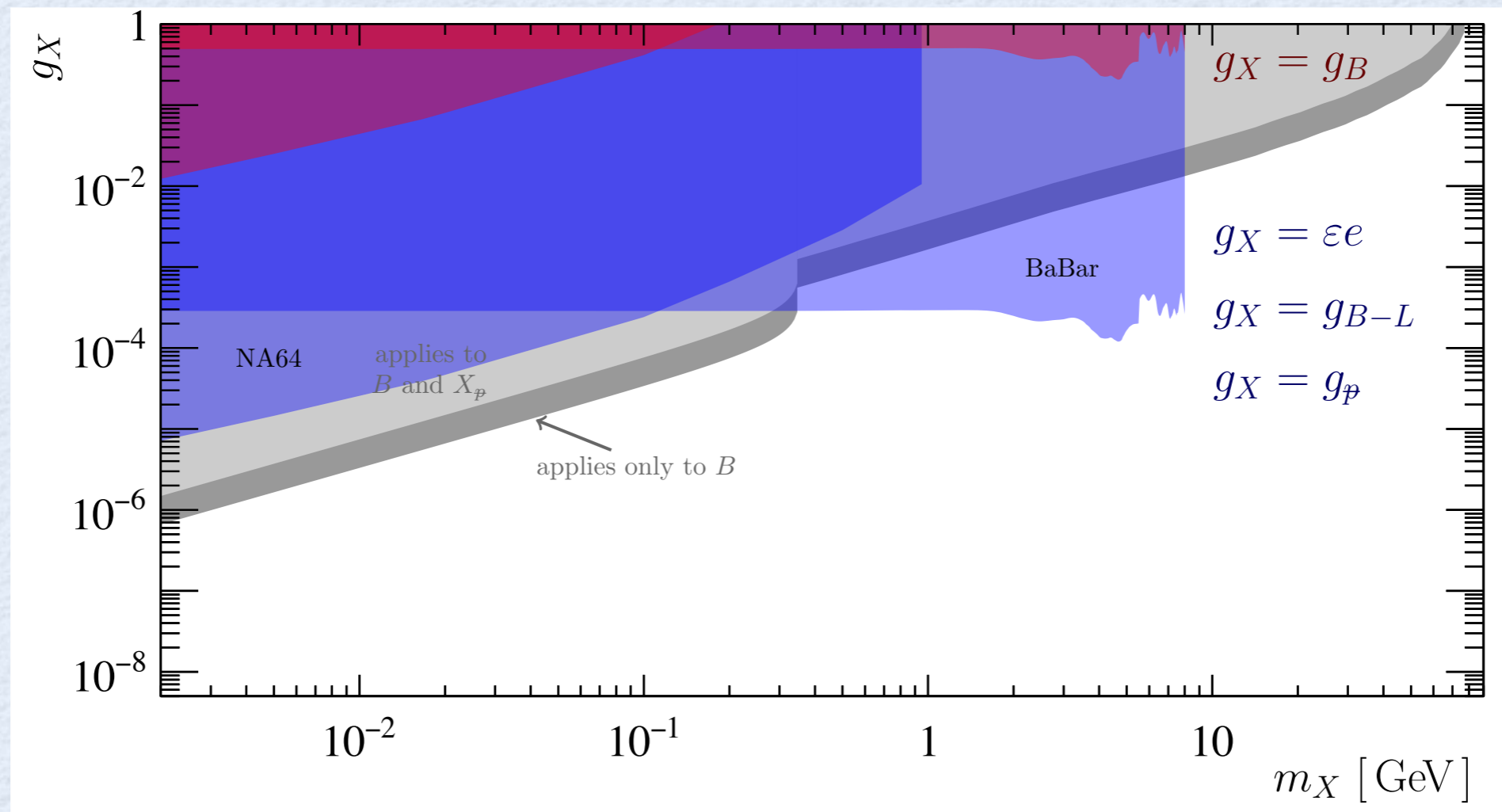
invisible decays

BEYOND THE DARK PHOTON MODEL

assuming dominant invisible decay

BEYOND THE DARK PHOTON MODEL

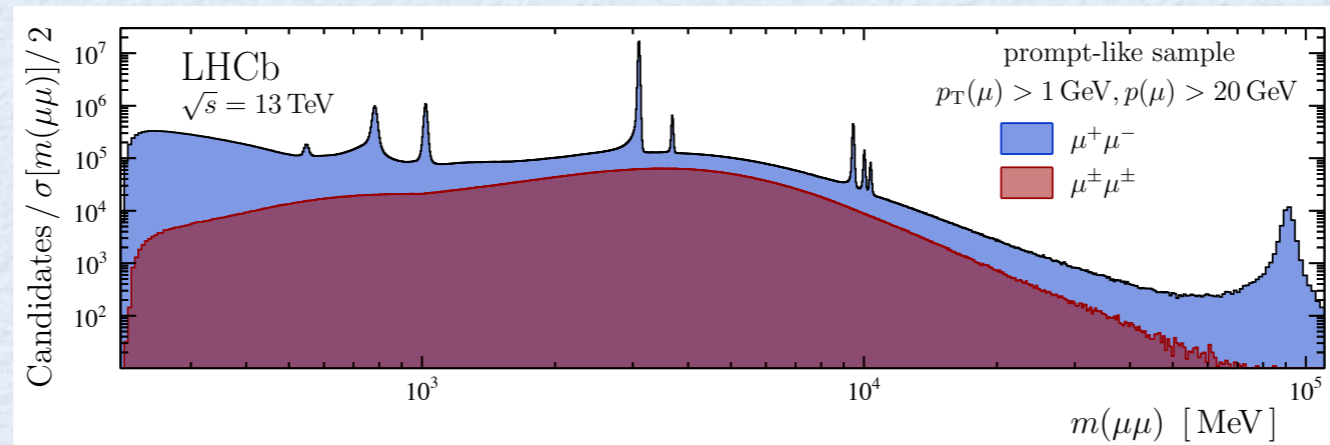
assuming dominant invisible decay



DARK PHOTON AT LHCb

a portal, dark photon

standard model \longleftrightarrow dark matter

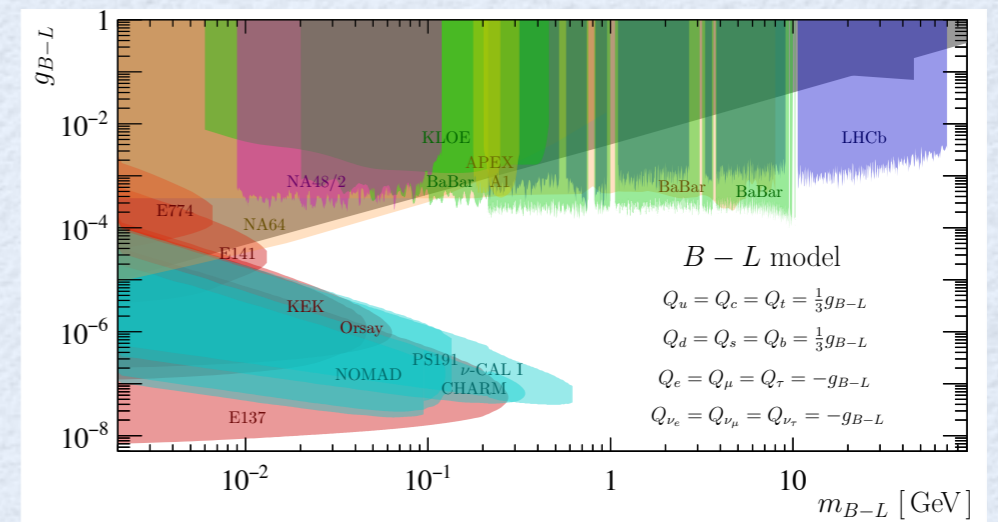
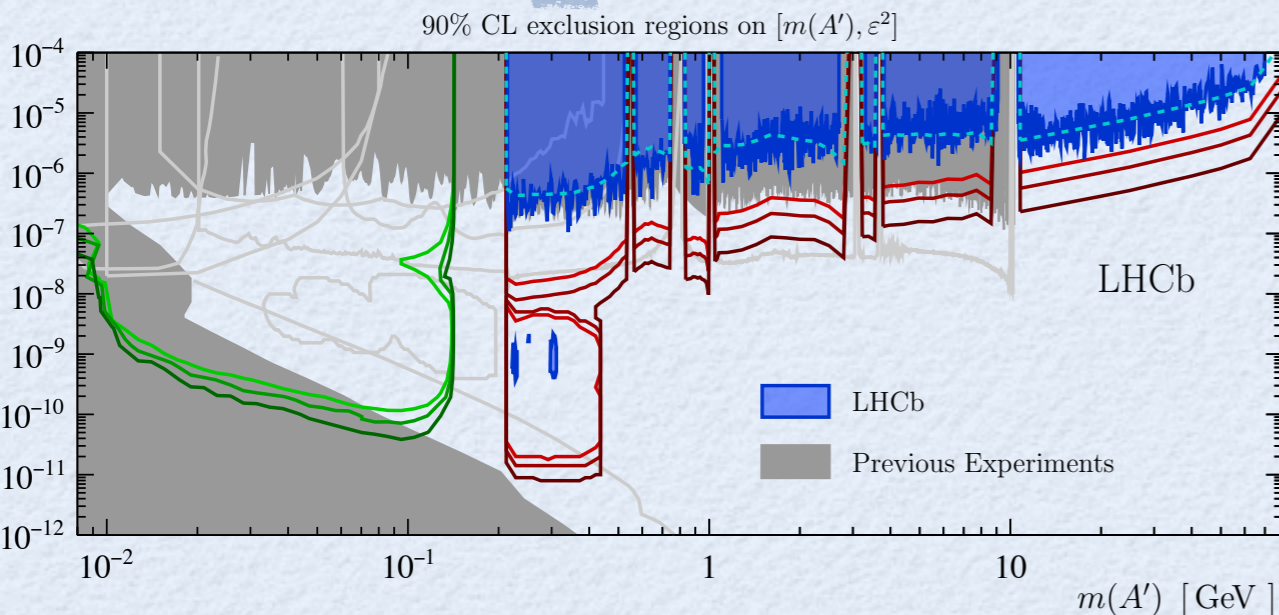


Ilten, YS, Thaler, Williams, Xue,
PRL 1603.08926

LHCb, 1710.02867

dark photon

generic models



Ilten, YS, Williams, Xue,
1801.04847

AXION LIKE PARTICLES

pseudo scalars

(appears in different BSM models)

$$\mathcal{L}_{\text{eff}} = -\frac{4\pi\alpha_s c_g}{\Lambda} a G^{\mu\nu} \tilde{G}_{\mu\nu} + \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

$$c_g \neq 0 \text{ or } c_\gamma \neq 0$$

Aloni, YS, Williams, 1811.03474

Aloni, Fanelli, YS, Williams, work in progress

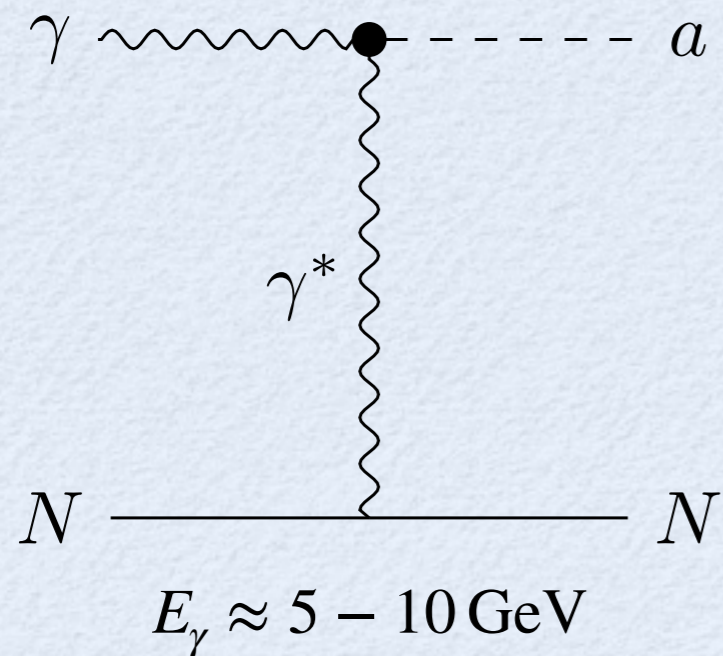
PRIMAKOFF ALP PRODUCTION

$$\mathcal{L}_{\text{eff}} = \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

PRIMAKOFF ALP PRODUCTION

$$\mathcal{L}_{\text{eff}} = \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

photon on fixed target

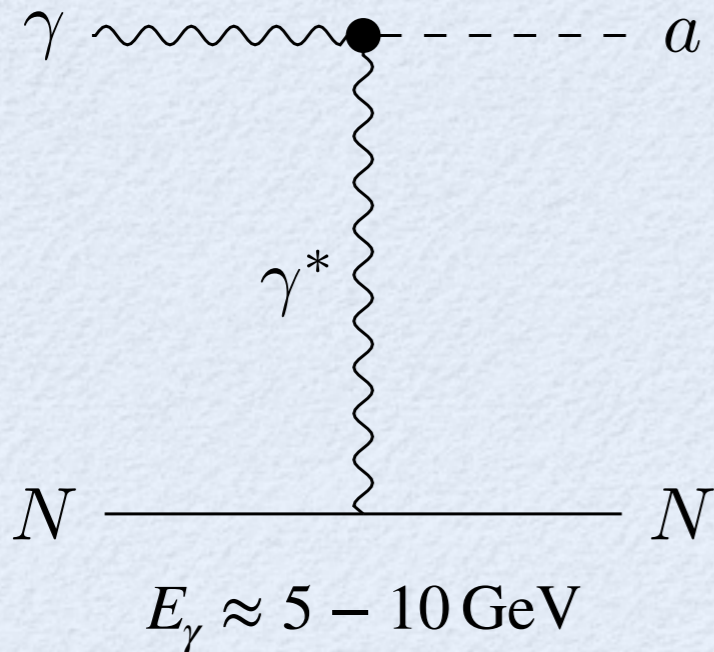


data driven signal
estimation

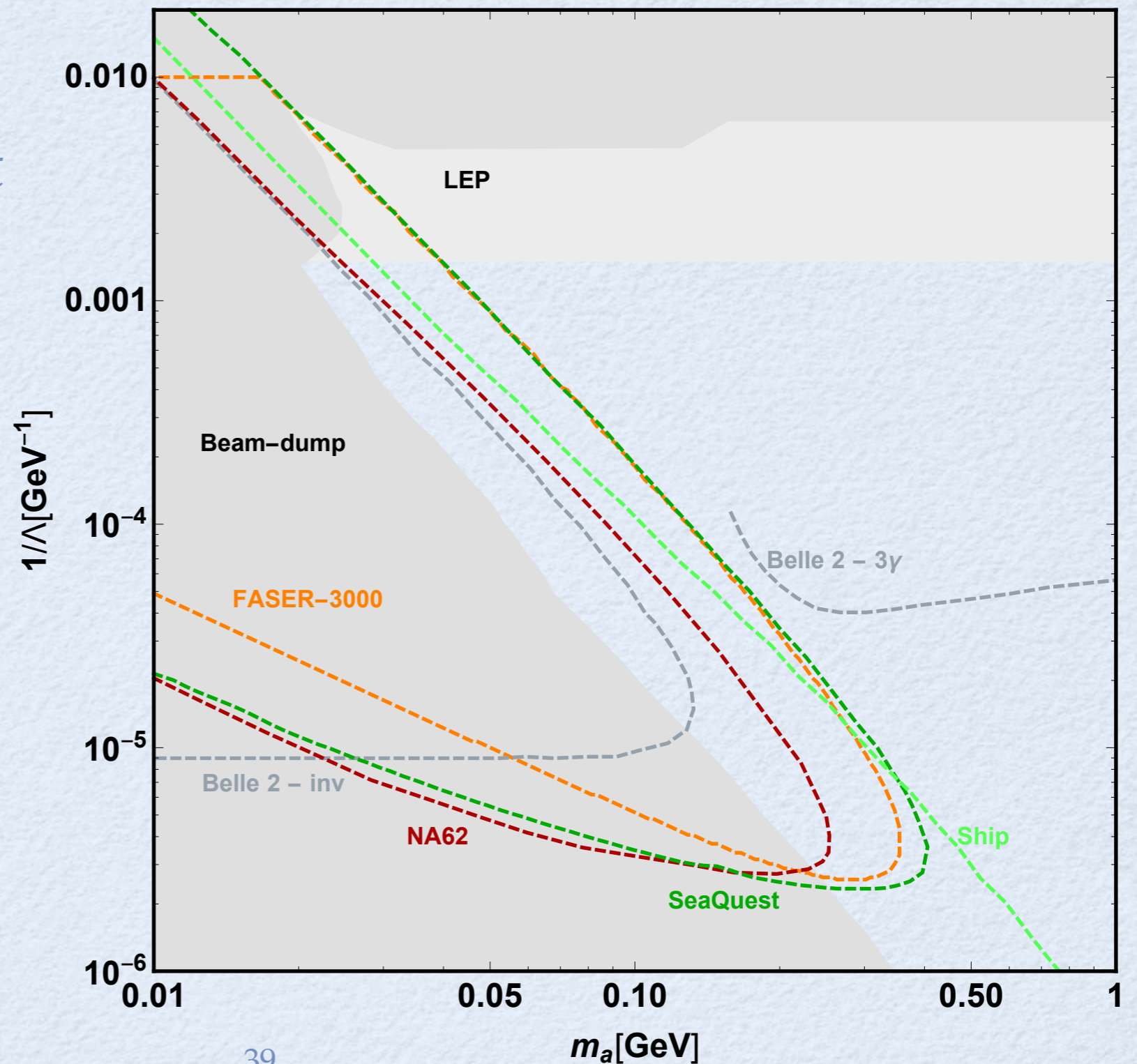
PRIMAKOFF ALP PRODUCTION

$$\mathcal{L}_{\text{eff}} = \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

photon on fixed target



data driven signal estimation

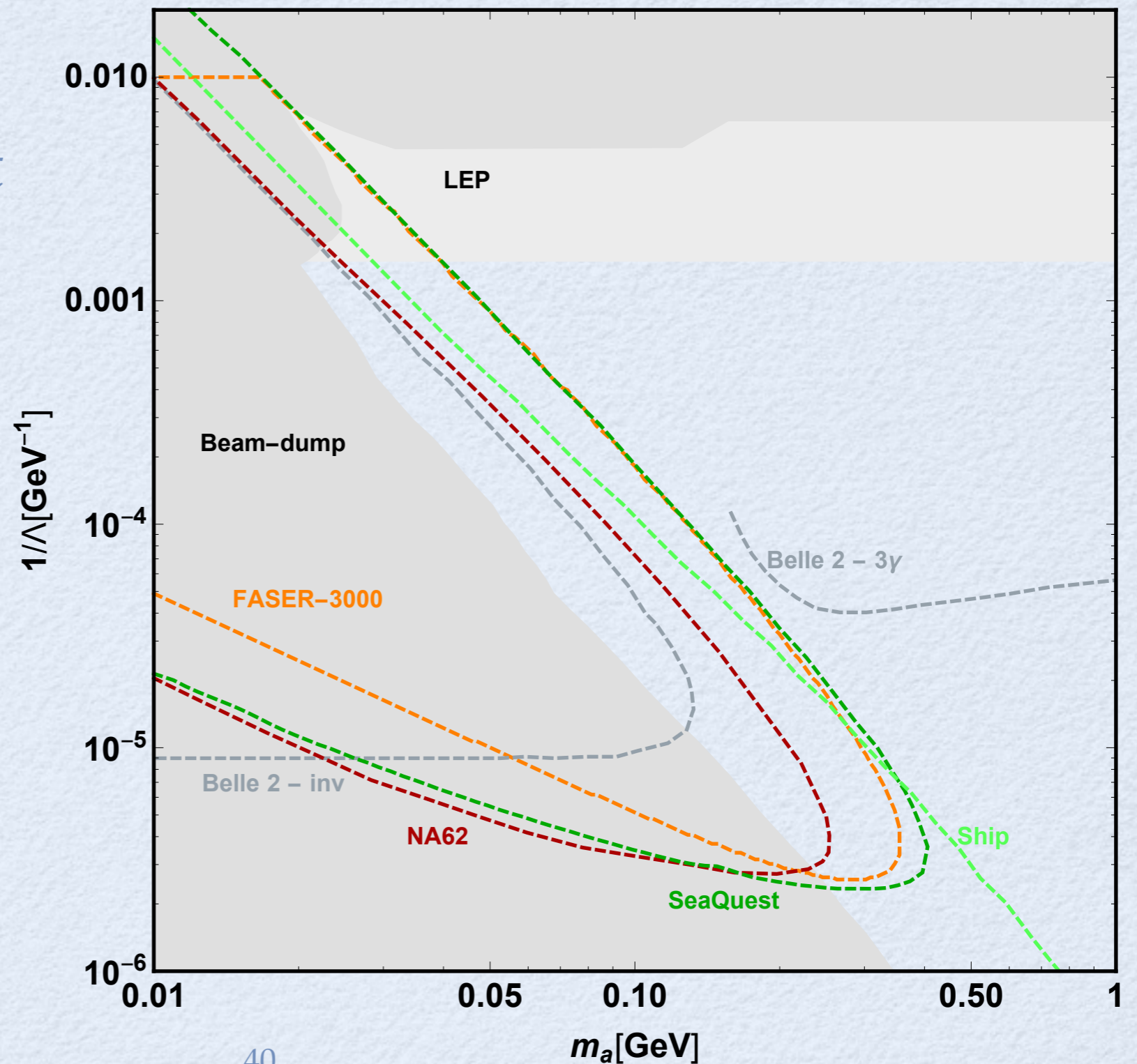
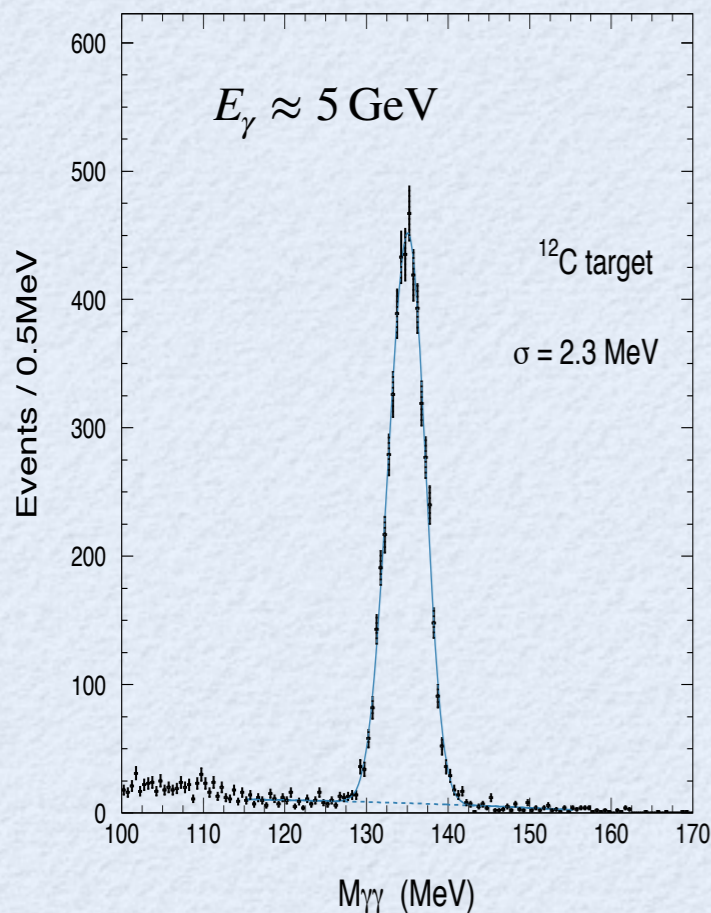


PRIMAKOFF ALP PRODUCTION

$$\mathcal{L}_{\text{eff}} = \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

photon on fixed target

PrimeEx, 1009.1681

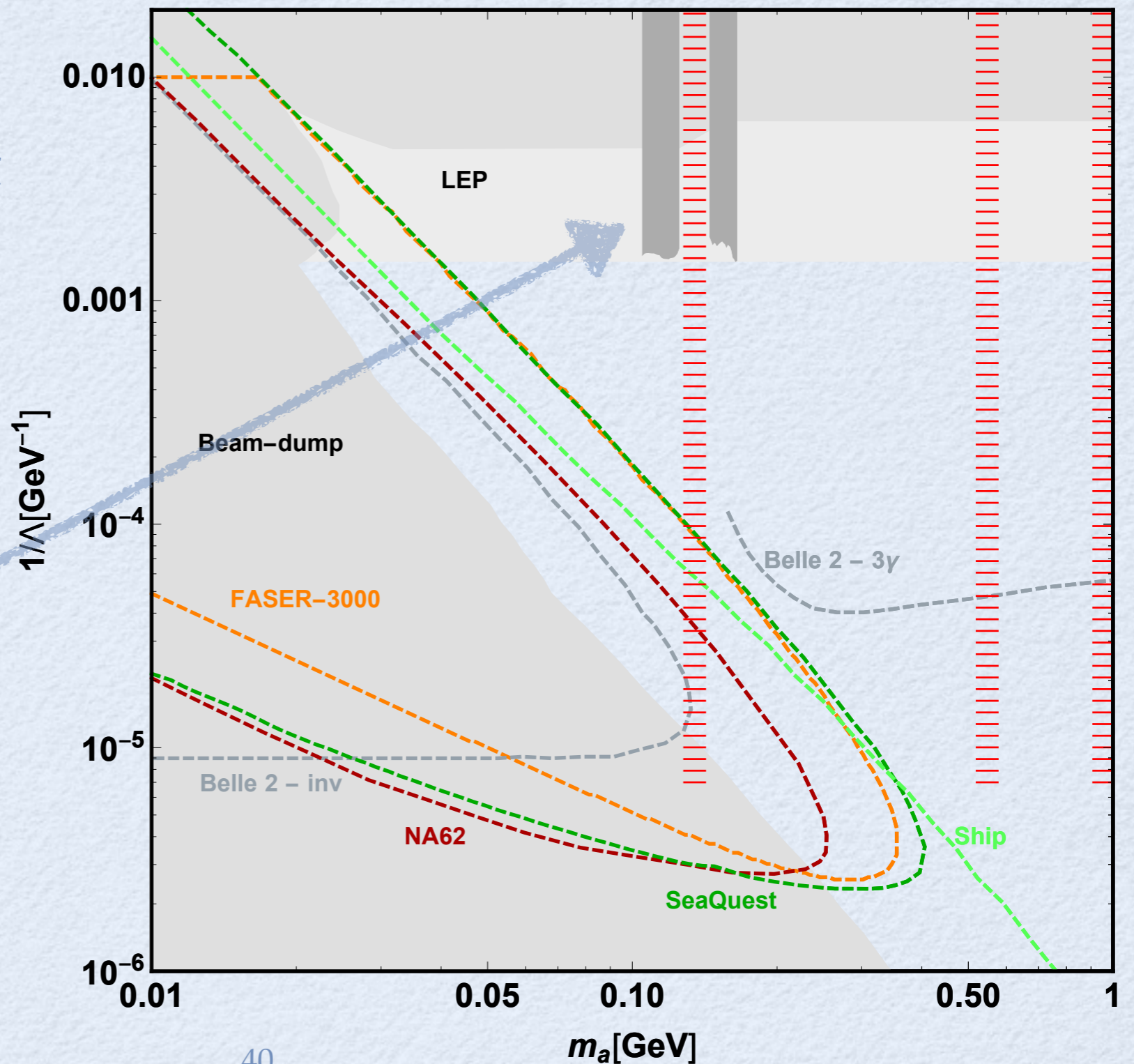
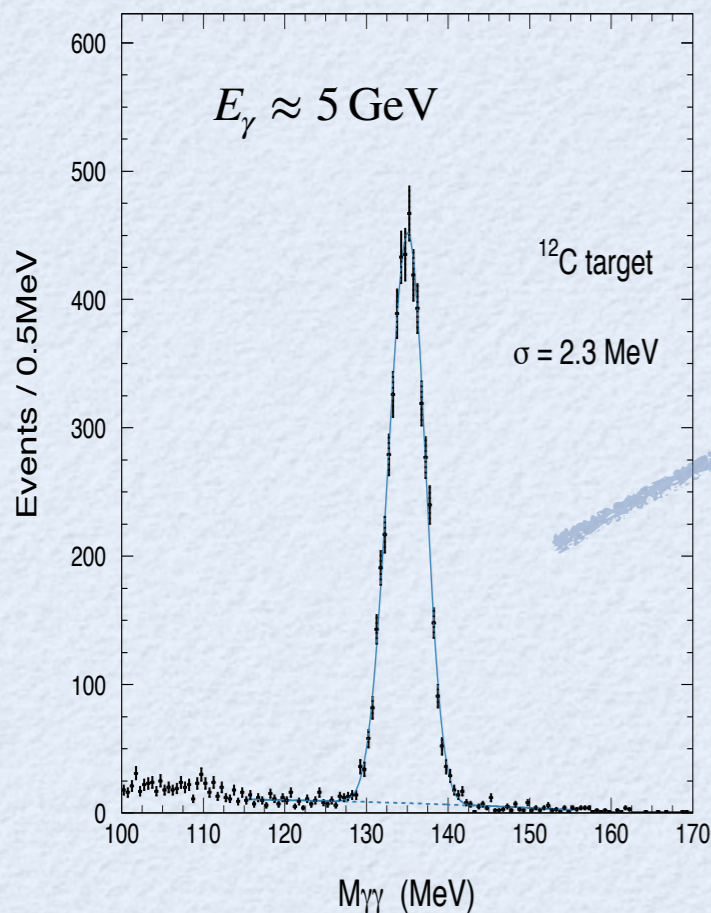


PRIMAKOFF ALP PRODUCTION

$$\mathcal{L}_{\text{eff}} = \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

photon on fixed target

PrimeEx, 1009.1681



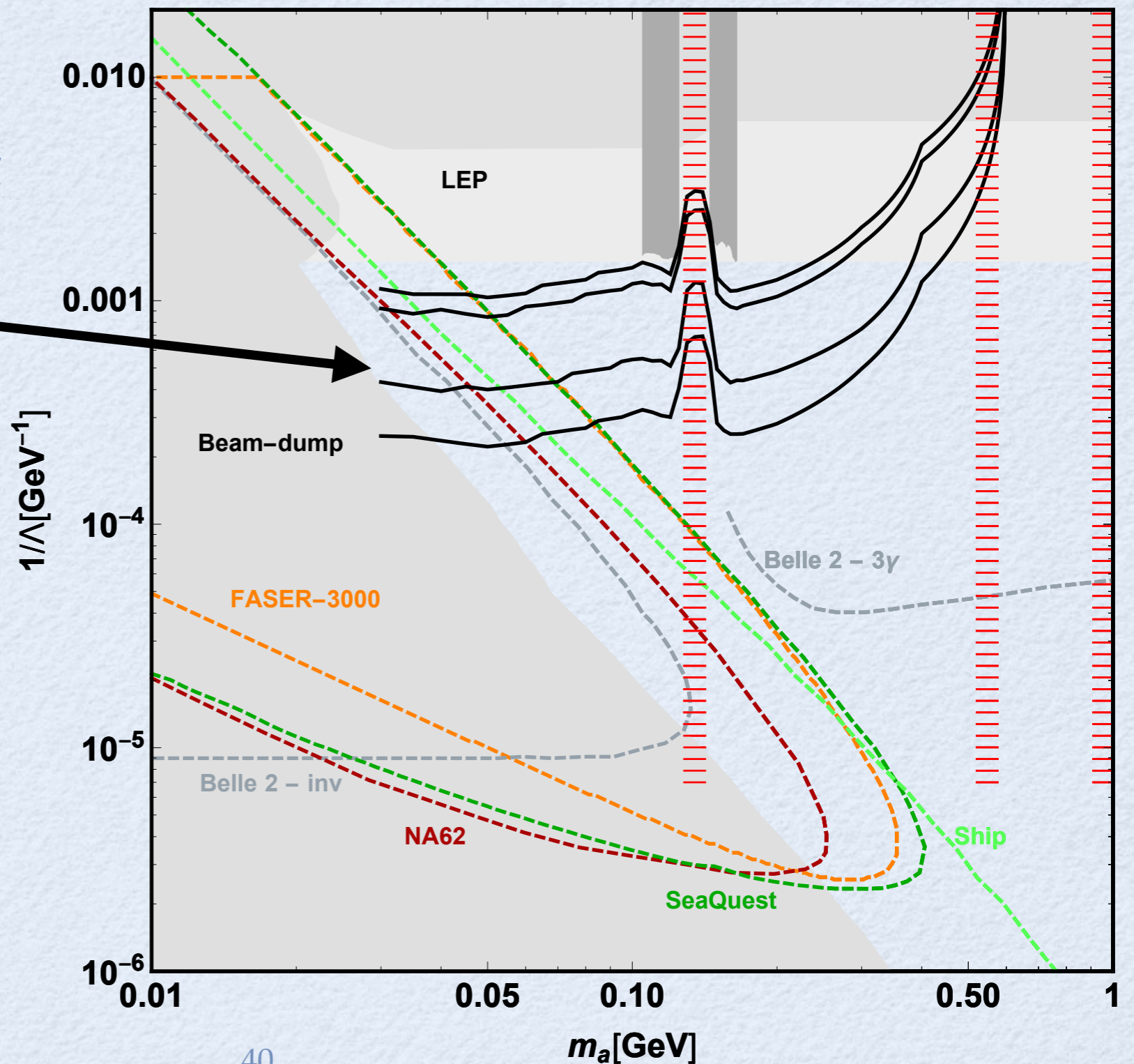
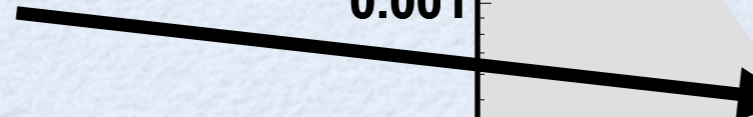
PRIMAKOFF ALP PRODUCTION

$$\mathcal{L}_{\text{eff}} = \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

photon on fixed target

PrimeEx, 5GeV

(C, Si, Pb
data on tape)



PRIMAKOFF ALP PRODUCTION

$$\mathcal{L}_{\text{eff}} = \frac{c_\gamma}{4\Lambda} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

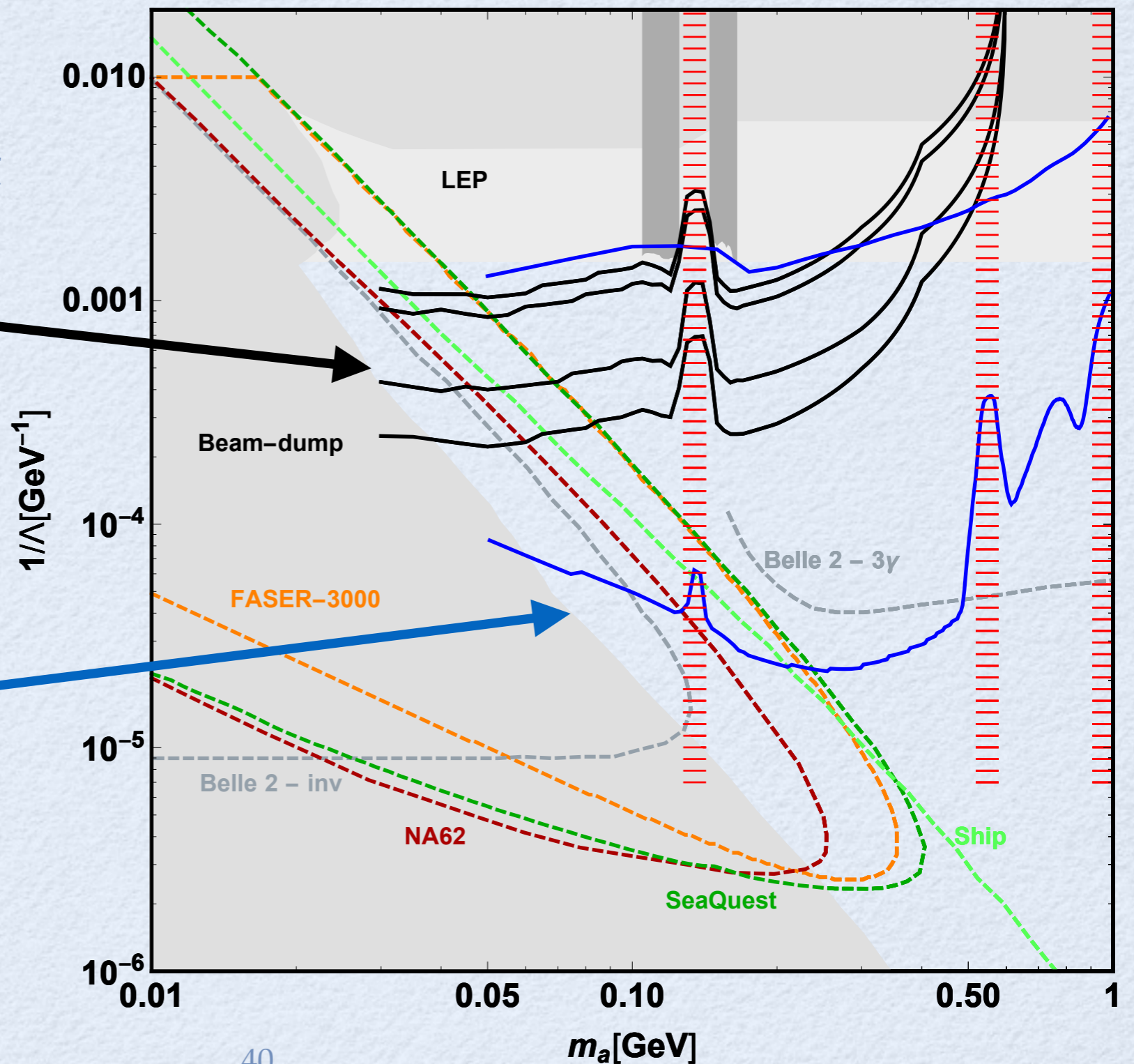
photon on fixed target

PrimeEx, 5GeV

(C, Si, Pb
data on tape)

GlueX, 8.5GeV

(p and Pb future run
1/fb)



ALPS AT THE GeV SCALE

$$\mathcal{L}_{\text{eff}} = -\frac{4\pi\alpha_s c_g}{\Lambda} a G^{\mu\nu} \tilde{G}_{\mu\nu}$$

how to estimate hadronic rates for
ALPs with GeV scale mass?

ALPS AT THE GEV SCALE

$$\mathcal{L}_{\text{eff}} = -\frac{4\pi\alpha_s c_g}{\Lambda} a G^{\mu\nu} \tilde{G}_{\mu\nu}$$

how to estimate hadronic rates for
ALPs with GeV scale mass?

$m_a \lesssim \text{GeV}$
chiral PT

???

$m_a \gtrsim 2\text{GeV}$
pQCD

ALPS AT THE GEV SCALE

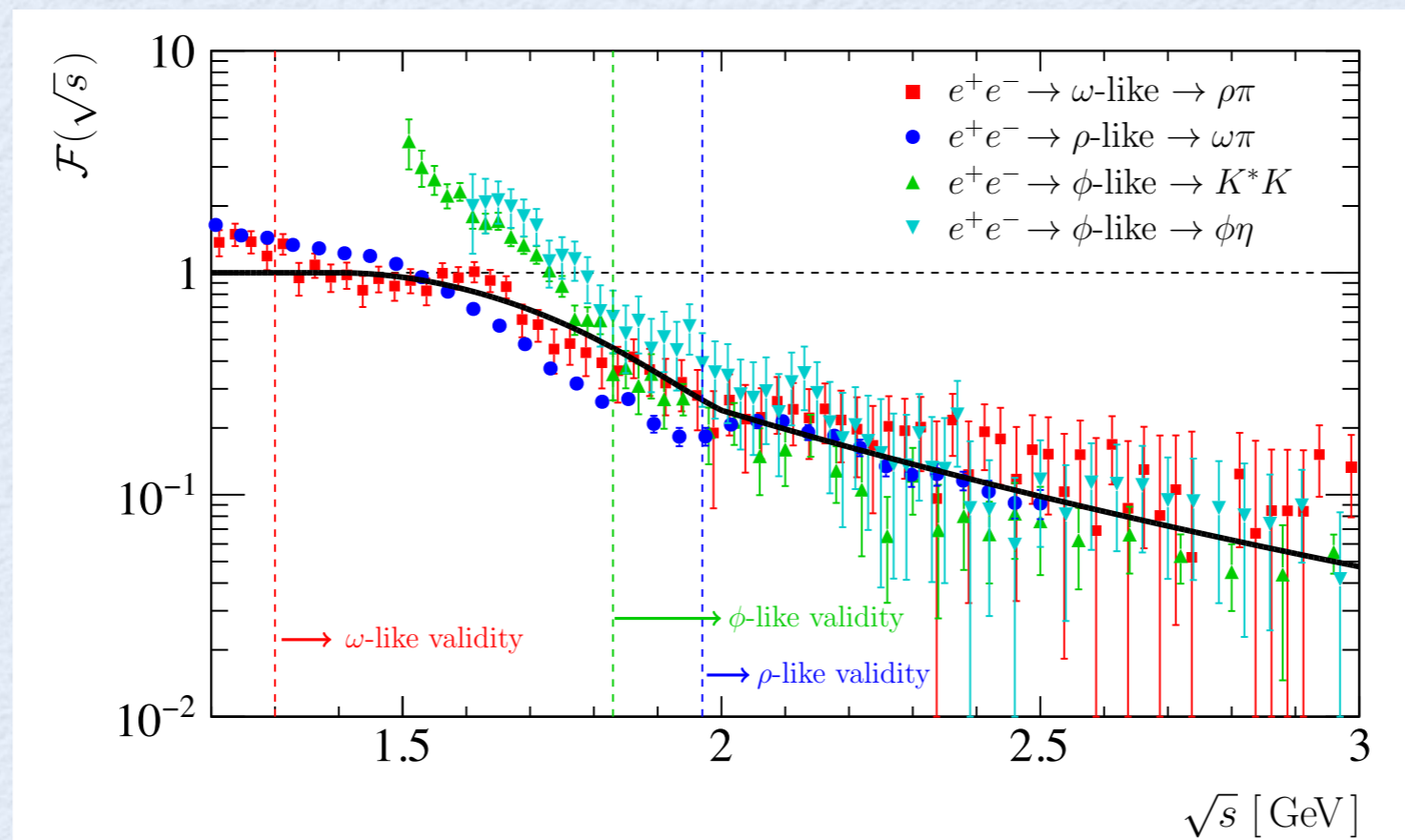
$$\mathcal{L}_{\text{eff}} = -\frac{4\pi\alpha_s c_g}{\Lambda} a G^{\mu\nu} \tilde{G}_{\mu\nu}$$

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use data!

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ALPS AT THE GEV SCALE

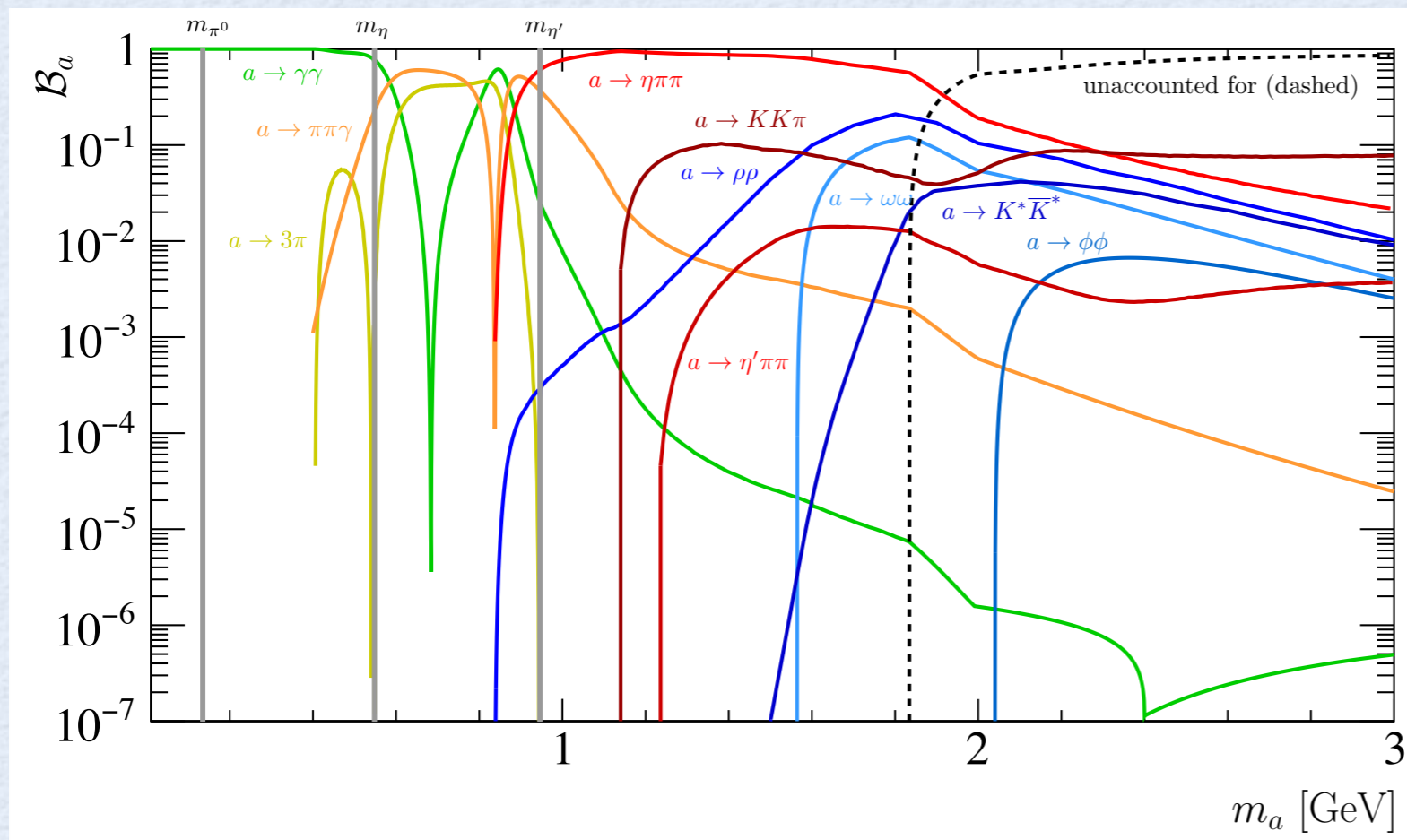
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ALPS AT THE GEV SCALE

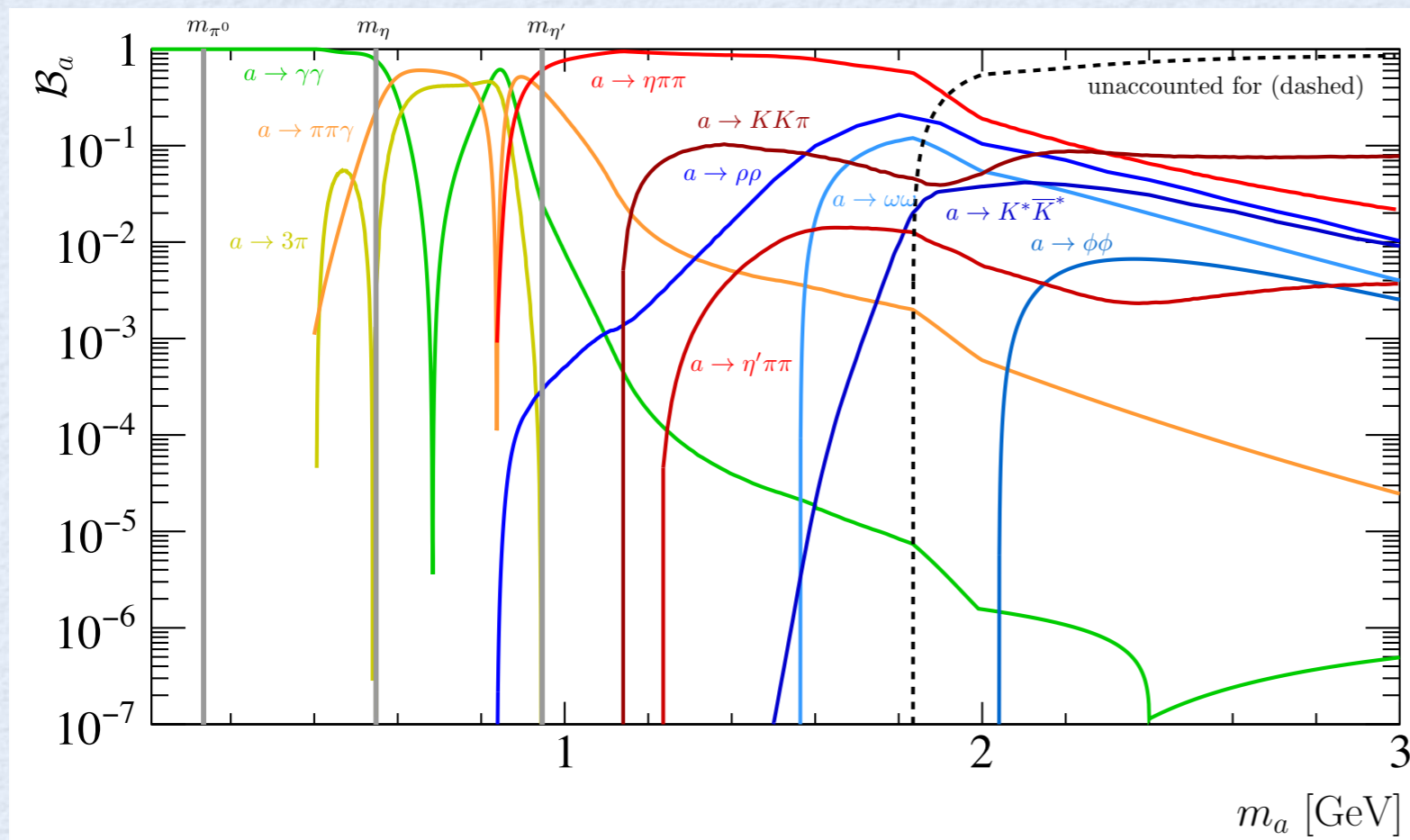
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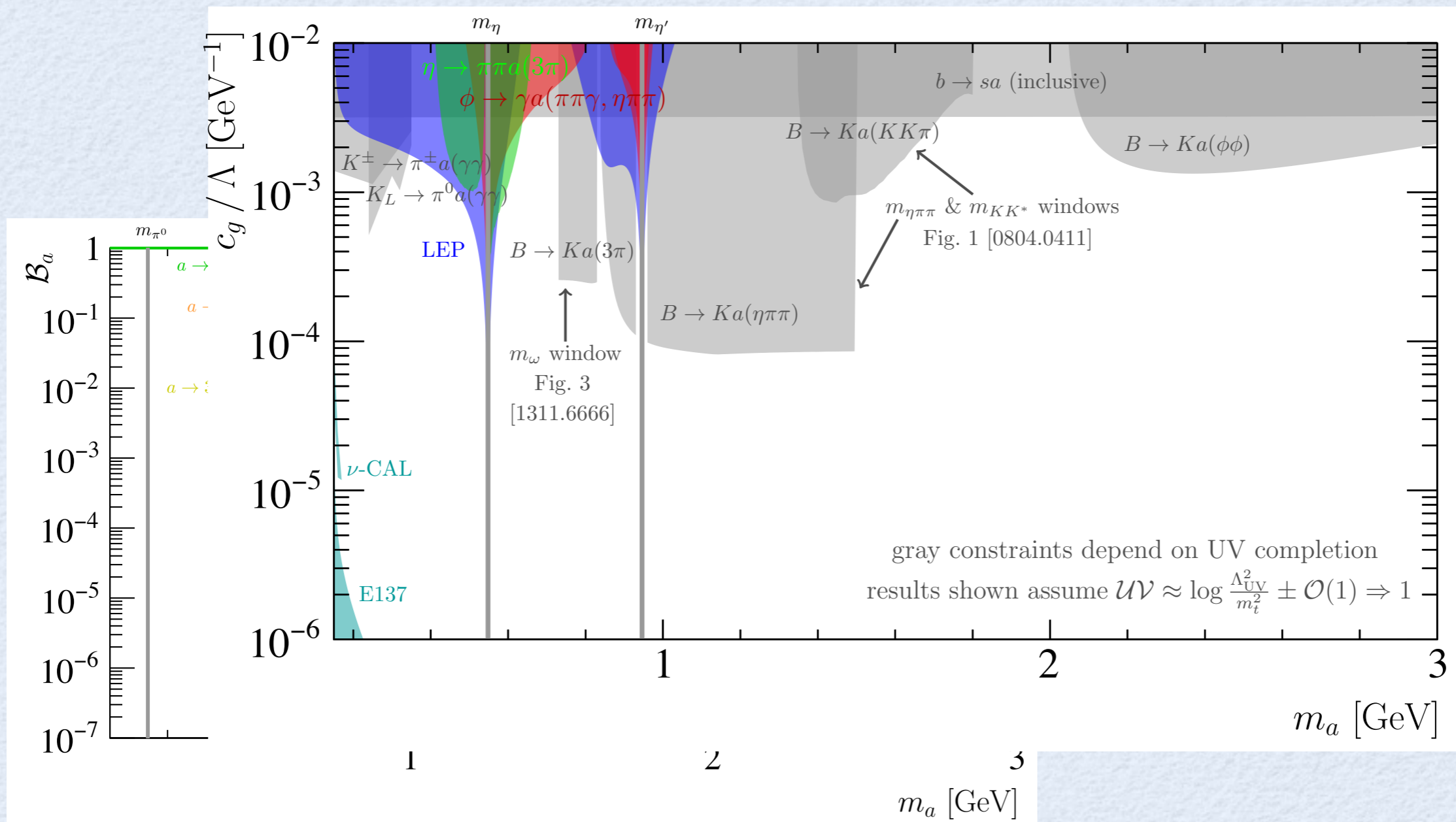
η_c cross check

	This Work VMD $\times \mathcal{F}(m) ^2$	Experiment Average	Experiment $SU(3)$
$\mathcal{B}(\eta_c \rightarrow \rho\rho)$	1.0%	$1.8 \pm 0.5\%$	$1.10 \pm 0.14\%$
$\mathcal{B}(\eta_c \rightarrow \omega\omega)$	0.40%	$0.20 \pm 0.10\%$	$0.44 \pm 0.06\%$
$\mathcal{B}(\eta_c \rightarrow \phi\phi)$	0.25%	$0.28 \pm 0.04\%$	$0.28 \pm 0.04\%$
$\mathcal{B}(\eta_c \rightarrow K^* \bar{K}^*)$	0.91%	$0.91 \pm 0.26\%$	$1.00 \pm 0.13\%$

ALPS AT THE GEV SCALE

$$\mathcal{L}_{\text{eff}} = -\frac{4\pi\alpha_s c_g}{\Lambda} a G^{\mu\nu} \tilde{G}_{\mu\nu}$$

how to estimate hadronic rates for ALPs with GeV scale mass?



beck

Experiment	$SU(3)$
average	$1.10 \pm 0.14\%$
	$0.10\% \pm 0.06\%$
	$0.04\% \pm 0.04\%$
	$0.26\% \pm 0.13\%$



OUTLOOK

- new physics beyond the standard model is well motivated, but with unknown scale
- we saw examples how to probe new forces in intensity and precision frontiers
- each of these examples probes unexplored territories and improve our understanding of Nature



