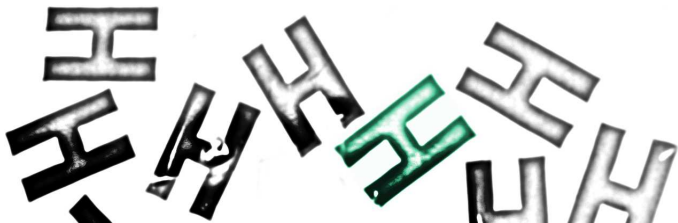


The package TopoID and its applications to Higgs physics

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Theory Seminar, University of Zurich
3rd of November, 2015



- 1** TopoID
 - Polynomial ordering
 - Topology classification
 - Relations among Feynman integrals
 - Topology “merging”
 - Partial fractioning

- 2** $N^3\text{LO}$ Higgs production: qq' -channel
 - Optical theorem and Cutkosky's rules
 - Calculation
 - Results

- 3** NLO and NNLO Higgs pair production
 - Differential factorization
 - Soft-virtual approximation
 - Asymptotic expansion
 - Calculation
 - Results

TopoID

Topology IDentification



Idea: Generic, process independent Mathematica package

Feynman diagrams \rightarrow reduced result (Laporta not included)

- Topology construction
(identification, minimal sets; partial fractioning; factorization, ...)
- Handle properties
(completeness, linear dependence; subtopologies, scalelessness, symmetries; graphs, unitarity cuts, ...)
- FORM code generator
(diagram mapping, topology processing, integral reduction, ...)
- Master integral identification
(base changes, non-trivial relations, ...)

Bring the polynomial P with m terms into unique form \hat{P} by renaming the n variables $\{x_j\}$: [Pak; '12]

- 1 Convert P into $m \times (n + 1)$ matrix $M^{(0)}$
(row: term, 1st column: coefficient, remaining columns: powers of $\{x_j\}$)
- 2 Start with considering the above $M^{(0)}$ and the 2nd column ($k = 1$)
- 3 Compute for all considered matrices $M^{(k),\sigma}$ all transpositions of columns k and $k + 1, \dots$ (and collect permutations σ)
- 4 Sort rows in each matrix lexicographically by the first k columns
- 5 Extract in columns k the lexicographically largest vector
- 6 Keep only matrices with this maximal vector;
If $k < n - 1$: $k \rightarrow k + 1$ and goto Step 3
- 7 Each remaining \hat{P}_σ matrix encodes the same unique \hat{P}_σ and a permutation of variables σ

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(row: term, 1st column: coefficient, remaining columns: powers of $\{x_j\}$)
- 2 Start with considering the above $M^{(0)}$ and the 2nd column ($k = 1$)

$$P = x_1^2 + 2x_1x_2 + x_2^2 + x_3^2 \quad \rightarrow \quad M^{(0)} = \begin{pmatrix} 1 & 2 & 0 & 0 \\ 2 & 1 & 1 & 0 \\ 1 & 0 & 2 & 0 \\ 1 & 0 & 0 & 2 \end{pmatrix} \quad (\text{Step 1})$$

$$S^{(1)} = \{M^{(0)(123)} = M^{(0)}\}, \quad k = 1 \quad (\text{Step 2})$$

- 3 Compute for all considered matrices $M^{(k),\sigma}$ all transpositions of columns k and $k + 1, \dots$ (and collect permutations σ)

$$S^{(1)} : M^{(1)(123)} = \left(\begin{array}{cc|cc} 1 & 2 & 0 & 0 \\ 2 & 1 & 1 & 0 \\ 1 & 0 & 2 & 0 \\ 1 & 0 & 0 & 2 \end{array} \right), \quad M^{(1)(213)} = \left(\begin{array}{cc|cc} 1 & 0 & 2 & 0 \\ 2 & 1 & 1 & 0 \\ 1 & 2 & 0 & 0 \\ 1 & 0 & 0 & 2 \end{array} \right),$$

$$M^{(1)(321)} = \left(\begin{array}{cc|cc} 1 & 0 & 0 & 2 \\ 2 & 0 & 1 & 1 \\ 1 & 0 & 2 & 0 \\ 1 & 2 & 0 & 0 \end{array} \right)$$

(Step 3-1)

- 4 Sort rows in each matrix lexicographically by the first k columns

$$S''^{(1)} : M''^{(1)(123)} = \begin{pmatrix} 1 & \mathbf{0} & 0 & 2 \\ 1 & \mathbf{0} & 2 & 0 \\ 1 & \mathbf{2} & 0 & 0 \\ 2 & \mathbf{1} & 1 & 0 \end{pmatrix}, \quad M''^{(1)(213)} = \begin{pmatrix} 1 & \mathbf{0} & 2 & 0 \\ 1 & \mathbf{0} & 0 & 2 \\ 1 & \mathbf{2} & 0 & 0 \\ 2 & \mathbf{1} & 1 & 0 \end{pmatrix},$$

$$M''^{(1)(321)} = \begin{pmatrix} 1 & 0 & 0 & 2 \\ 1 & 0 & 2 & 0 \\ 1 & 2 & 0 & 0 \\ 2 & 0 & 1 & 1 \end{pmatrix}$$

(Step 4-1)

- 5** Extract in columns k the lexicographically largest vector
- 6** Keep only matrices with this maximal vector;
If $k < n - 1$: $k \rightarrow k + 1$ and goto Step 3

$$\hat{M}''^{(1)} = \begin{pmatrix} 0 \\ 0 \\ 2 \\ 1 \end{pmatrix} \quad (\text{Step 5-1})$$

$$S^{(2)} = \{M''^{(1)(123)}, M''^{(1)(213)}\}, \quad k = 2 \quad (\text{Step 6-1})$$

- 3 Compute for all considered matrices $M^{(k),\sigma}$ all transpositions of columns k and $k + 1, \dots$ (and collect permutations σ)

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$$M^{(2)(213)} = \left(\begin{array}{ccc|c} 1 & 0 & 2 & 0 \\ 1 & 0 & 0 & 2 \\ 1 & 2 & 0 & 0 \\ 2 & 1 & 1 & 0 \end{array} \right), \quad M^{(2)(231)} = \left(\begin{array}{ccc|c} 1 & 0 & 0 & 2 \\ 1 & 0 & 2 & 0 \\ 1 & 2 & 0 & 0 \\ 2 & 1 & 0 & 1 \end{array} \right)$$

(Step 3-2)

- 4 Sort rows in each matrix lexicographically by the first k columns

$$S''^{(2)} : M''^{(2)(123)} = \begin{pmatrix} 1 & 0 & \mathbf{0} & 2 \\ 1 & 0 & \mathbf{2} & 0 \\ 1 & 2 & \mathbf{0} & 0 \\ 2 & 1 & \mathbf{1} & 0 \end{pmatrix}, \quad M''^{(2)(132)} = \begin{pmatrix} 1 & 0 & 0 & 2 \\ 1 & 0 & 2 & 0 \\ 1 & 2 & 0 & 0 \\ 2 & 1 & 0 & 1 \end{pmatrix},$$

$$M''^{(2)(213)} = \begin{pmatrix} 1 & 0 & \mathbf{0} & 2 \\ 1 & 0 & \mathbf{2} & 0 \\ 1 & 2 & \mathbf{0} & 0 \\ 2 & 1 & \mathbf{1} & 0 \end{pmatrix}, \quad M''^{(2)(231)} = \begin{pmatrix} 1 & 0 & 0 & 2 \\ 1 & 0 & 2 & 0 \\ 1 & 2 & 0 & 0 \\ 2 & 1 & 0 & 1 \end{pmatrix}$$

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- 7 Each remaining matrix encodes the same unique \hat{P}_σ and a permutation of variables σ

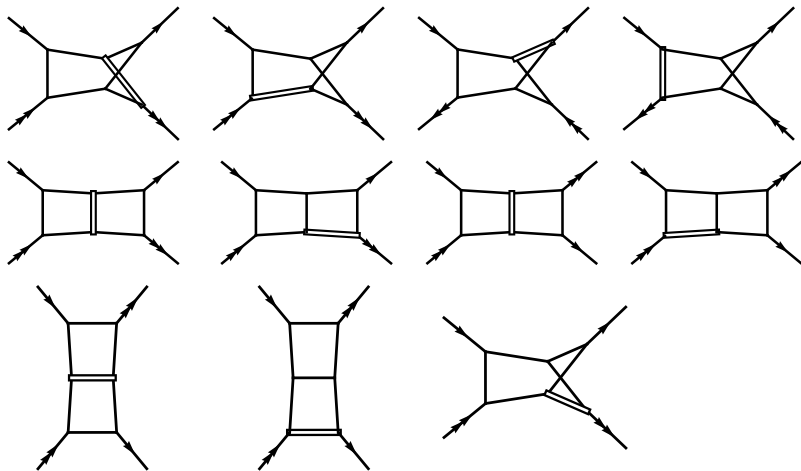
$$\hat{P} = P = x_3^2 + x_2^2 + x_1^2 + 2x_1x_2, \quad \hat{\sigma} = \{(123), (213)\} \quad (\text{Step 7})$$

$$\hat{P} = P = x_3^2 + x_2^2 + x_1^2 + 2x_1x_2, \quad \hat{\sigma} = \{(123), (213)\} \quad (\text{Step 7})$$

- P already in canonical form; two permutations $\{(123), (213)\}$
- (213) denotes $(x_1, x_2, x_3) \rightarrow (x_2, x_1, x_3)$; symmetry of P under $x_1 \leftrightarrow x_2$

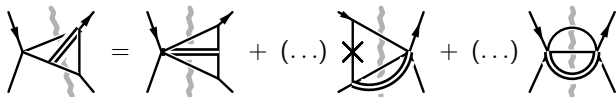
Application to Feynman integrals

- Use $\mathcal{U} + \mathcal{F}$ from the Feynman representation
 - Unique identifier $\hat{\mathcal{U}} + \hat{\mathcal{F}}$; independent of momentum space representation
 - Returned permutations: symmetries of Feynman integrals
- \Rightarrow Many useful applications

Minimal set for NNLO Higgs production:

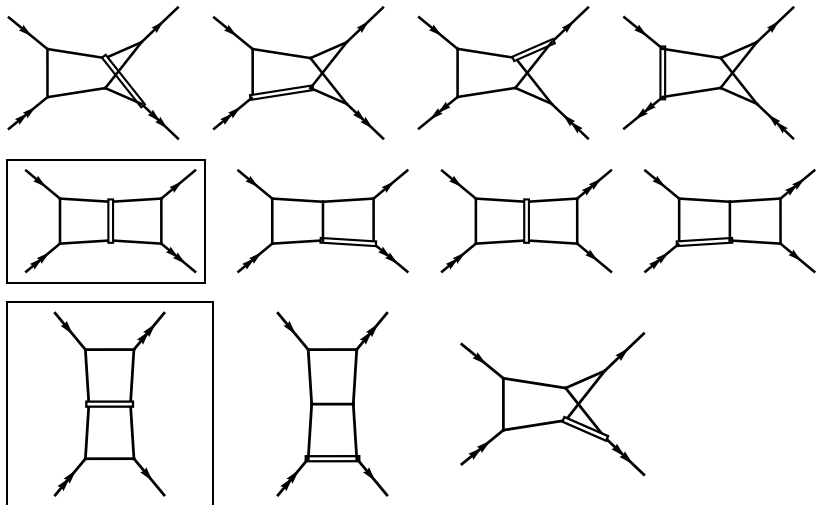
Note: Sufficient for all 2946 diagrams

Non-trivial relation for NNLO Higgs production:



- Cross-topology relations; not from Laporta reduction
- Simplify calculation
- Usefull cross-checks

Minimal set for NNLO Higgs production:



3-loop massless propagators:

$q_1 = k_1,$

$q_4 = p - k_1 - k_2,$

$q_7 = k_1 + k_2 + k_3,$

$q_2 = p - k_1,$

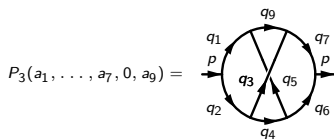
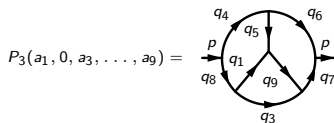
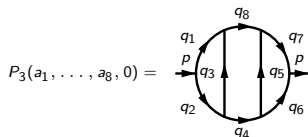
$q_5 = k_3,$

$q_8 = k_1 + k_2,$

$q_3 = k_2,$

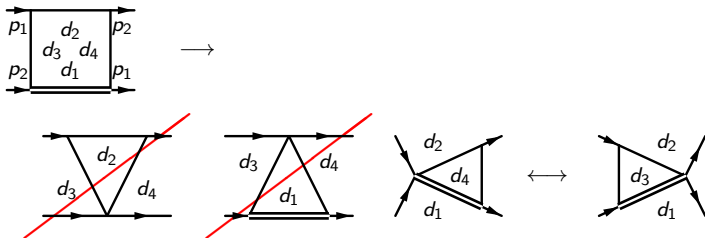
$q_6 = p - k_1 - k_2 - k_3,$

$q_9 = k_1 + k_3.$



- 1 external, 3 internal momenta
- ⇒ 9 scalar products
- 3 incomplete topologies with 8 lines
- Identify "greatest common subtopology"
- Find "supertopology" with these 3 different graphs as subtopologies

Partial fractioning for NLO Higgs production:

Via Gröbner basis:

$$\begin{aligned}
 d_4 &\rightarrow -m_H^2 + s + d_1 + d_2 - d_3 \\
 \frac{d_3}{d_4} &\rightarrow \frac{1}{d_4} (-m_H^2 + s + d_1 + d_2 - d_4) \\
 \frac{d_2}{d_3 d_4} &\rightarrow \frac{1}{d_3 d_4} (m_H^2 - s - d_1 + d_3 + d_4) \\
 \frac{d_1}{d_2 d_3 d_4} &\rightarrow \frac{1}{d_2 d_3 d_4} (m_H^2 - s - d_2 + d_3 + d_4) \\
 \frac{1}{d_1 d_2 d_3 d_4} &\rightarrow \frac{1}{(m_H^2 - s) d_1 d_2 d_3 d_4} (d_1 + d_2 - d_3 - d_4)
 \end{aligned}$$

N³LO Higgs production: qq' -channel

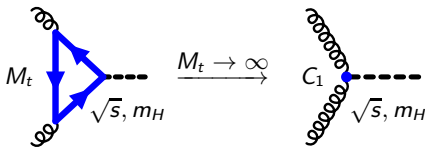
Motivation and introduction

- Higgs production at the LHC dominated by gluon fusion
- After [Anastasiou, Duhr, Dulat, Furlan, Gehrmann, Herzog, Mistlberger; '14, '15]:
2.2% corrections, 3% uncertainty at N³LO

⇒ Cross-check

- Loop-induced process, dominated by top quark mass
- Effective field theory with top quark integrated out:

$$\mathcal{L}_{Y,\text{eff}} = -\frac{H}{v} C_1 \mathcal{O}_1 \quad \text{and} \quad \mathcal{O}_1 = \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu}$$



- Reduced numbers of scales and loops:
single dimensionless variable $x = m_H^2/s$ (soft: $x \rightarrow 1$)
- Finite matching coefficient C_1 needed to 4-loop

[Chetyrkin, Kniehl, Steinhauser; '98] [Schröder, Steinhauser; '06] [Chetyrkin, Kühn, Sturm; '06]

N³LO Higgs production: qq' -channel

Status

- LO calculation (exact)

[Ellis et al.; '76] [Wilczek et al.; '77] [Georgi et al.; '78] [Rizzo; '80]

- NLO (exact)

[Dawson; '91] [Djouadi, Spira, Zerwas; '91]

- NNLO (EFT) \Rightarrow soft expansion to 3rd order valid to $\mathcal{O}(1\%)$

[Harlander, Kilgore; '02] [Anastasiou, Melnikov; '02] [Ravindran, Smith, van Neerven; '03]

- NNLO $\mathcal{O}(1/M_t^2)$ corrections \approx NNLO +1%

[Pak, Rogal, Steinhauser; '09-'11] [Harlander, Mantler, Marzani, Ozeren; '09-'10]

- N³LO IR counterterms

- 3-loop splitting functions

[Moch, Vermaseren, Vogt; '02]

- NNLO master integrals to higher orders in ϵ

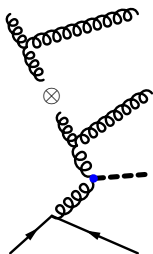
[Pak, Rogal, Steinhauser; '11] [Anastasiou, Bühler, Duhr, Herzog; '12]

- Cross sections and convolution integrals

[Höschele, JH, Pak, Steinhauser, Ueda; '12, '13] [Bühler, Lazopoulos; '13]

- N³LO scale variation $\Rightarrow \mathcal{O}(2\% - 8\%)$

[Bühler, Lazopoulos; '13]



N^3 LO Higgs production: qq' -channel

Status

■ N^3 LO corrections

- VV^2 and V^3 – 3-loop gluon form factor

[Baikov, Chetyrkin, Smirnov², Steinhauser; '09] [Gehrmann, Glover, Huber, Izkizlerli, Studerus; '09]

- VRV exact in x

[Anastasiou, Duhr, Dulat, Herzog, Mistlberger; '13] [Kilgore; '14]

- V^2R exact in x

[Dulat, Mistlberger; '14] [Duhr, Gehrmann, Jaquier; '14]

- VR^2 expansion in $x \rightarrow 1$

[Anastasiou, Duhr, Dulat, Furlan, Gehrmann; '14] [Li, von Manteuffel, Schabinger, Zhu; '14],

[Anastasiou, Duhr, Dulat, Furlan, Herzog, Mistlberger; '15]

- R^3 expansion in $x \rightarrow 1$

[Anastasiou, Duhr, Dulat, Mistlberger; '13]

- 37 terms in the $x \rightarrow 1$ expansion

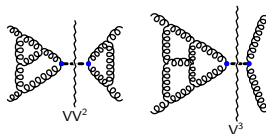
[Anastasiou, Duhr, Dulat, Furlan, Gehrmann, Herzog, Mistlberger; '14, '15]

⇒ Sufficient for phenomenology

- qq' -channel exact in x (VR^2 , R^3)

[Höchele, JH, Ueda; '14] [Anzai, Hasselhuhn, Höchele, JH, Kilgore, Steinhauser, Ueda; '15]

⇒ Independent cross-check



■ Many different resummations

[Catani et al.] [Moch et al.] [Ahrens et al.] [de Florian et al.] [Ball et al.] [...]

$N^3\text{LO}$ Higgs production: qq' -channel

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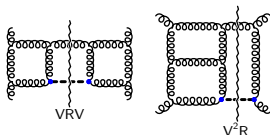
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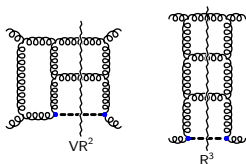
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$N^3\text{LO}$ Higgs production: qq' -channel

Generalities

- 1 Reduction of integrals with full x -dependence
→ Only contributing cuts
- 2 Construct differential equations for master integrals
→ Canonical basis (in general: coupled system)
- 3 Soft limit $x \rightarrow 1$ as boundary condition
→ Leading term using Mellin-Barnes representation

Canonical differential equations:

$$\frac{d}{dx} m_i(x, \epsilon) = \epsilon A_{ij}(x) m_j(x, \epsilon) \quad \text{with} \quad d = 4 - 2\epsilon$$

[Henn et al.; '13-...]

- ϵ - and x -dependence factorize
- Solve order-by-order in ϵ
- A_{ij} : alphabet of appearing functions

E.g. Harmonic Polylogarithms (HPLs):

$$H_{\vec{w}}(x) = \int_0^x dx' f_{w_1}(x') H_{\vec{w}_{n-1}}(x') \quad \text{and} \quad f_0(x) = \frac{1}{x}, \quad f_{\pm 1}(x) = \frac{1}{1 \mp x}$$

Optical theorem and Cutkosky's rules

Higher-order corrections:

Virtual More loop integrations

Real Additional final state particles (different phase spaces)

Optical theorem:

$$\sigma(i \rightarrow f) \sim \sum_f \int d\Pi_f |\mathcal{M}(i \rightarrow f)|^2 \sim \text{Disc } \mathcal{M}(i \rightarrow i)$$

Cutkosky's rules: Consider only valid diagrammatic cuts for Disc

- Two connectivity components
- Separate in- and outgoing momenta (s -channel)
- Contribute to the process (1 or 2 Higgs, 0 to 3 parton lines)

Optical theorem and Cutkosky's rules

E.g. to NNLO:

$$\begin{aligned} & \int d\Pi_1 \left| \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \\ \dots \end{array} \right|^2 + \int d\Pi_2 \left| \begin{array}{c} \text{Diagram 3} \\ \text{Diagram 4} \\ \dots \end{array} \right|^2 \\ & + \int d\Pi_3 \left| \begin{array}{c} \text{Diagram 5} \\ \text{Diagram 6} \\ \dots \end{array} \right|^2 + \dots \\ & = \begin{array}{c} \text{Diagram 7} + \text{Diagram 8} + \dots + \text{Diagram 9} \\ \text{Diagram 10} + \dots + \text{Diagram 11} + \dots \end{array} \end{aligned}$$

The diagrams are Feynman diagrams for a scattering process. Diagrams 1-6 are tree-level diagrams with a dashed line representing a cut. Diagrams 7-11 are diagrams with a dashed line representing a cut, where the cut is connected to a loop structure. The diagrams are arranged in a grid-like structure, with the first row showing the squared magnitudes of the tree-level diagrams, the second row showing the squared magnitudes of the loop diagrams, and the third row showing the sum of the diagrams with a cut, which is equal to the sum of the squared magnitudes of the diagrams with a cut.

Optical theorem and Cutkosky's rules

Optical theorem:

$$\sigma(i \rightarrow f) \sim \sum_f \int d\Pi_f |\mathcal{M}(i \rightarrow f)|^2 \sim \text{Disc } \mathcal{M}(i \rightarrow i)$$

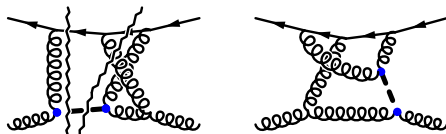
- Pros**
 - Forward scattering \Rightarrow simplified kinematics
 - Common treatment of loop and phase space integrals
 - Calculation of Disc only for master integrals
- Cons**
 - More loops and diagrams
 - Only total cross section (naively)

Approach first used in [\[Anastasiou, Melnikov; '02\]](#)

Optical theorem and Cutkosky's rules

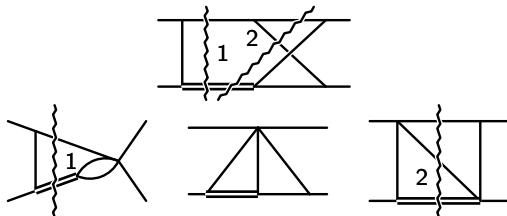
Handling cut-diagrams:

Filter diagrams



- Fast graph-based algorithm build into Per1 script to process QGRAF output
- N^3 LO Higgs: 860 118 \rightarrow 174 938
- NNLO Higgs pair (SV): 17 667 600 \rightarrow 42 252

Assist reduction



- Build also into TopoID \Rightarrow pass to Laporta reduction
- Typically: only $\mathcal{O}(10\%)$ of subtopologies (or relations)

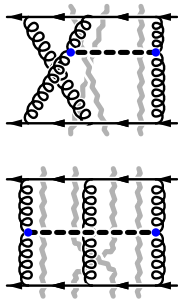
$N^3\text{LO}$ Higgs production: qq' -channel

Calculation: toolchain

Reduction

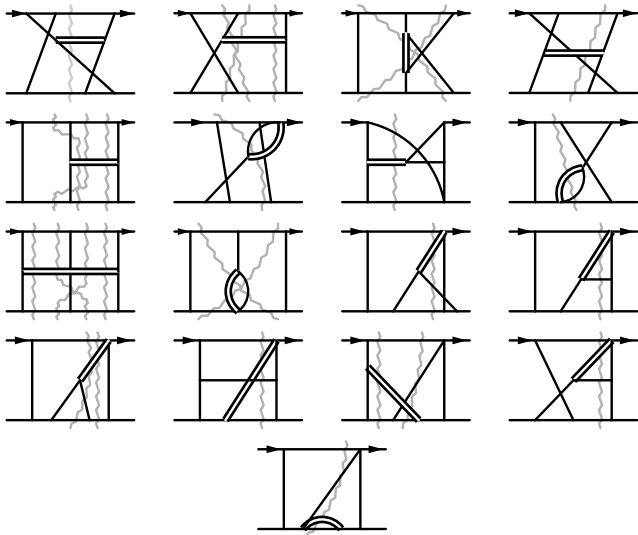
- 1 Generate Feynman diagrams
QGRAF [Nogueira; '93]
- 2 Select diagrams with specific cuts
filter [JH, Pak; (unpublished)]
- 3 Map diagrams to topologies (← graph information)
exp [Harlander, Seidensticker, Steinhauser; '98]
reg [Pak; (unpublished)]
- 4 Reduction to scalar integrals (← generic topologies)
FORM [Kuipers, Ueda, Vermaseren, Vollinga; '13]
- 5 Reduction to master integrals (← basic topologies)
rows [JH, Pak; (unpublished)]
FIRE [Smirnov]
- 6 Minimal basis of master integrals
TopoID [JH, Pak; (unpublished)]

220 diagrams, e.g.



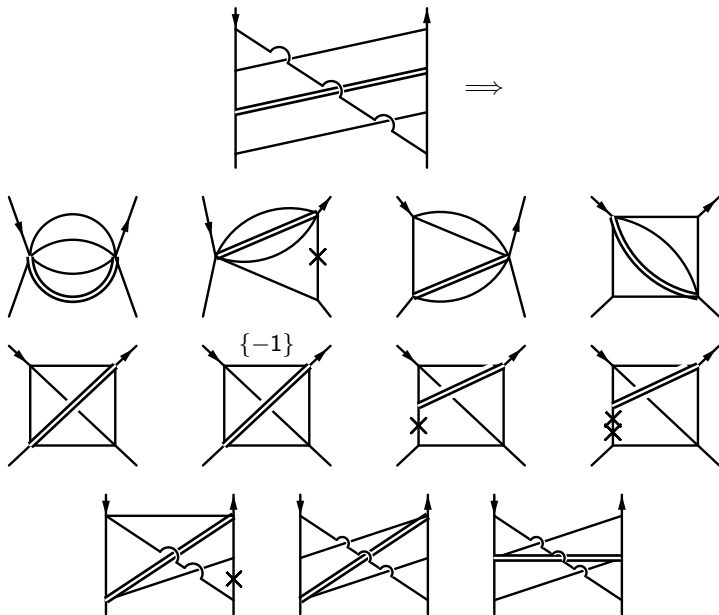
$N^3\text{LO}$ Higgs production: qq' -channel

Calculation: 17 topologies with 3- and 4-particle cuts



$N^3\text{LO}$ Higgs production: qq' -channel

Calculation: e.g. "sea snake" topology



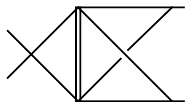
N^3LO Higgs production: qq' -channel

Results: functions beyond HPLs

- Laporta: 332 master integrals; TopoID: 108 and cancellation of ξ
- Some Feynman integrals generate functions with alphabet beyond HPLs:

$$\left\{ \frac{1}{x}, \frac{1}{1-x}, \frac{1}{1+x}, \frac{1}{1+4x}, \frac{1}{x\sqrt{1+4x}} \right\}$$

- Traced back to common subtopology:



- Numerical implementation in Mathematica:

- Change alphabet to $\left\{ \frac{1}{x}, \frac{1}{1-x}, \frac{1}{1+x}, \frac{1}{1+4x}, \frac{1}{x} \left(\frac{1}{\sqrt{1+4x}} - 1 \right) \right\}$

- Use series expansions $x \rightarrow 0$ and $x \rightarrow 1 \Rightarrow 10$ digits in 1 second

- Note: $x \rightarrow (1-x)/x^2 \Rightarrow$ “Cyclotomic Polylogarithms”

- Representation as “Goncharov Polylogarithms” (linear denominators)

- Letters of 6th roots of unity; here only: $\left\{ \frac{1}{x}, \frac{1}{1-x}, \frac{1}{(-1)^{\frac{1}{3}} - x} \right\}$

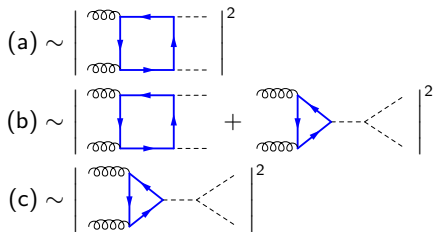
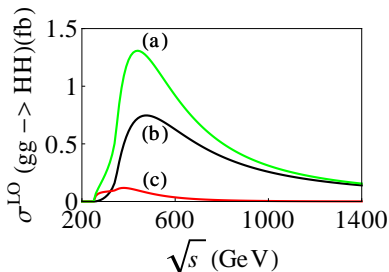
NLO and NNLO Higgs pair production

Motivation

Higgs Potential in the Standard Model:

$$V(H) = \frac{1}{2} m_H^2 H^2 + \lambda v H^3 + \frac{1}{4} \lambda H^4, \quad \lambda^{\text{SM}} = \frac{m_H^2}{2v^2} \approx 0.13, \quad v: \text{Higgs vev.}$$

- Verify mechanism of spontaneous symmetry breaking in the SM
- Measure the Higgs self-coupling \Rightarrow sensitive process



NLO and NNLO Higgs pair production

Theory status

Prospects for the LHC @ 14 TeV:

- $b\bar{b}\gamma\gamma$ -channel, 600 fb^{-1} : $\lambda \neq 0$ [Baur, Plehn, Rainwater; '04]
- $b\bar{b}\gamma\gamma$ -, $b\bar{b}\tau^+\tau^-$ -channels: “promising”;
 $b\bar{b}W^+W^-$ -channel: “not promising” [Baglio, Djouadi, Gröber, Mühlleitner, Quevillon, Spira; '13]
- 600 fb^{-1} : $\lambda > 0$; 3000 fb^{-1} : $\lambda_{-20\%}^{+30\%}$ (ratio with Higgs cross section) [Goertz, Papaefstathiou, Yang, Zurita; '13]
- And many others, e.g.:
[Dolan, Englert, Spannowsky; '12] [Papaefstathiou, Yang, Zurita; '13]
[Barr, Dolan, Englert, Spannowsky; '13] [Barger, Everett, Jackson, Shaughnessy; '14]
[Englert, Krauss, Spannowsky, Thompson; '15] [...]
- Until now: Higgs pair production not observed in $b\bar{b}b\bar{b}$ - and $b\bar{b}\gamma\gamma$ -channels (as expected in the SM) [ATLAS; '15] [CMS; '15]

⇒ **Wait for HL-LHC**

NLO and NNLO Higgs pair production

Theory status

Known since long:

- LO result with exact M_t dependence [Glover, van der Bij; '88] [Plehn, Spira, Zerwas; '98]
- NLO result in $M_t \rightarrow \infty$ limit [Dawson, Dittmaier, Spira; '98]

$$\sigma_H \approx 20^{\text{LO}} \text{ fb} + 20^{\text{NLO}, M_t \rightarrow \infty} \text{ fb} \quad \text{for } \sqrt{s_H} = 14 \text{ TeV}, \mu = 2m_H$$

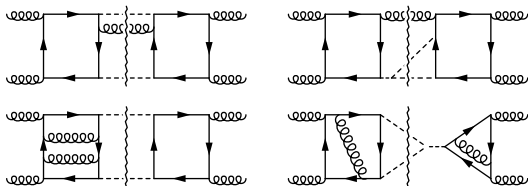
More recently:

- NLO + NNLL ($M_t \rightarrow \infty$) \approx NLO +20% [Shao, Li, Li, Wang; '13]
- NNLO w/ or w/o soft-virtual approx. ($M_t \rightarrow \infty$) \approx NLO +20% [de Florian, Mazzitelli; '13]
- $\mathcal{O}(1/M_t^8)$ corrections at NLO \approx NLO +10% [Grigo, JH, Melnikov, Steinhauser; '13]
- NLO real exact in M_t , NLO virt. for $M_t \rightarrow \infty \approx$ NLO **-10%** [Maltoni, Vryonidou, Zaro; '14]
- Cross-check of virtual NNLO corr.; NNLO matching coefficient for $ggHH$ -coupling \approx NNLO +1% [Grigo, Melnikov, Steinhauser; '14]
- Improved $\mathcal{O}(1/M_t^{12})$ NLO, $\mathcal{O}(1/M_t^4)$ NNLO soft-virt. corrections [Grigo, JH, Steinhauser; 15']

NLO and NNLO Higgs pair production

Generalities

- Operate on full-theory diagrams at NLO and NNLO
- Virt. corr. in two independent calculations:
 - amplitude (differential; 2-/3-loop)
 - forward scattering (total; 4-/5-loop)
- Real corr.: only via forward scattering at NLO
- Perform expansion for $M_t \rightarrow \infty$; improve upon effective theory results for NLO [Dawson, Dittmaier, Spira; '98], NNLO [de Florian, Mazzitelli; '13]
- Laporta reduction to master integrals for the “soft” subdiagrams
- Remaining “hard” massive tadpoles via MATAD
- Master integrals as series around $\sqrt{s} = 2m_H$ (not in this talk)



Differential factorization

Factorization of the LO result exact in M_t for:

Total cross section

$$\sigma^{(i)} = \Delta^{(i)} \sigma_{\text{exact}}^{(0)} = \frac{\sigma_{\text{exact}}^{(0)}}{\sigma_{\text{exp}}^{(0)}} \int_{4m_H^2}^s dQ^2 \frac{d\sigma_{\text{exp}}^{(i)}}{dQ^2}$$

$$\text{with } \Delta^{(i)} = \frac{\sigma_{\text{exp}}^{(i)}}{\sigma_{\text{exp}}^{(0)}}, \quad \sigma_{\text{exp}}^{(i)} = \sum_{n=0}^N c_n^{(i)} \rho^n, \quad \rho = \frac{m_H^2}{M_t^2}$$

Differential cross section

$$\sigma^{(i)} = \int_{4m_H^2}^s dQ^2 \frac{\left(\frac{d\sigma_{\text{exact}}^{(0)}}{dQ^2} \right)}{\left(\frac{d\sigma_{\text{exp}}^{(0)}}{dQ^2} \right)} \frac{d\sigma_{\text{exp}}^{(i)}}{dQ^2}$$

“Cure” the invalidity of the $M_t \rightarrow \infty$ expansion for the large- Q^2 region

- Virt. corr. via amplitude: **access to Q^2 -dependence $\sim \delta(s - Q^2)$**
- Real corr. via optical theorem (naively): **only total cross section**
⇒ Use the soft-virtual approximation [de Florian, Mazzitelli; '12]

Soft-virtual approximation

- Split σ into its contributions (works also for $d\sigma/dQ^2$):

$$\begin{aligned}\sigma &= \sigma^{\text{virt+ren}} + \sigma^{\text{real+split}} = \text{finite} \\ &= \underbrace{\Sigma_{\text{div}} + \Sigma_{\text{fin}}}_{=\Sigma_{\text{SV}}=\text{finite}} + \underbrace{\Sigma_{\text{soft}} + \Sigma_{\text{hard}}}_{=\Sigma_{\text{H}}=\text{finite}}\end{aligned}$$

- Σ_{div} universal for color-less final state [de Florian, Mazzitelli; '12]

- Compute $\sigma^{\text{virt+ren}}$ as ρ -expansion

- Solve $\sigma^{\text{virt+ren}} = \Sigma_{\text{div}} + \Sigma_{\text{fin}}$ for Σ_{fin}

- Σ_{div} and Σ_{soft} (soft coll. counterterms + soft real corr.)
 \sim exact σ^{LO} (include M_t effects)

$$Q^2 \frac{d\sigma}{dQ^2} = \sigma^{\text{LO}} z G(z) \quad \text{with} \quad z = \frac{Q^2}{s}, \quad G(z) = G_{\text{SV}}(z) + G_{\text{H}}(z)$$

$$\sigma_{(\text{SV})} = \int_{1-\delta}^1 dz \sigma^{\text{LO}}(zs) G_{(\text{SV})}(z) \quad \text{with} \quad \delta = 1 - \frac{4m_H^2}{s}$$

- $G_{\text{SV}}^{(i)}(z)$ constructed from $\sigma_{\text{fin}}^{(i)}$ and σ^{LO} only

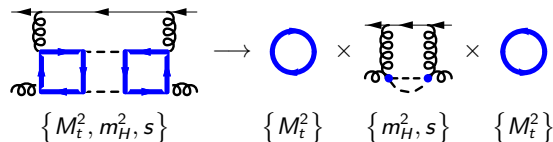
[de Florian, Mazzitelli; '12] [Grigo, JH, Steinhauser; '15]

Asymptotic expansion

- Expand at integrand level for all contributing regions
≡ series expansion in analytic result
- Hierarchy: $M_t^2 \gg s, m_H^2 \Rightarrow$ series in $\rho = m_H^2/M_t^2$
- Effectively reduce number of loops and scales

- Here: regions correspond to subgraphs (in general more than one)
- Hard mass expansion: subgraphs must contain all heavy lines

Example: NLO real with one region

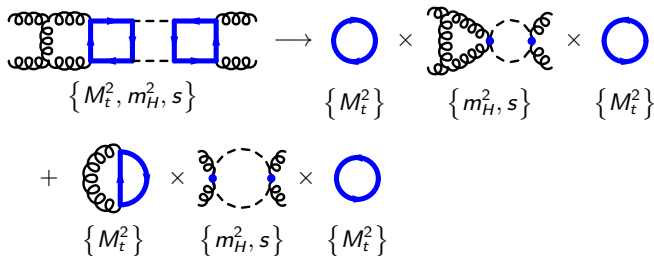


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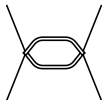
Example: NLO virt. with two regions



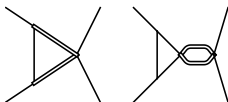
NLO and NNLO Higgs pair production

Calculation (via optical theorem)

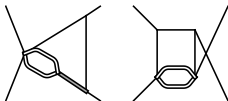
LO topology:



Virt. NLO topologies:



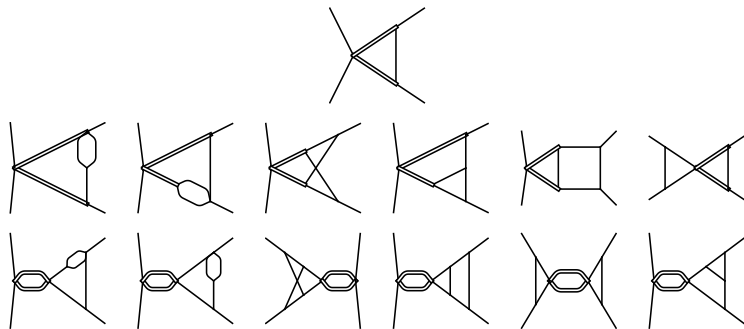
Real NLO topologies:



NLO and NNLO Higgs pair production

Calculation (via optical theorem)

Virt. NNLO topologies:



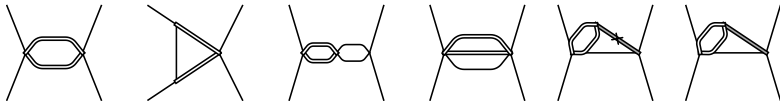
Note:

- Different regions in asymptotic expansion \Rightarrow different loop-orders
- Here: multiplied with 1- to 3-loop massive tadpoles

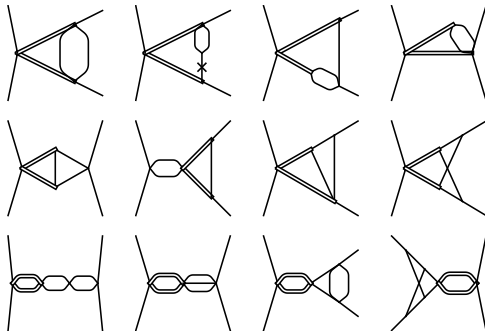
NLO and NNLO Higgs pair production

Calculation (via optical theorem)

Virt. and real LO-NLO master integrals:



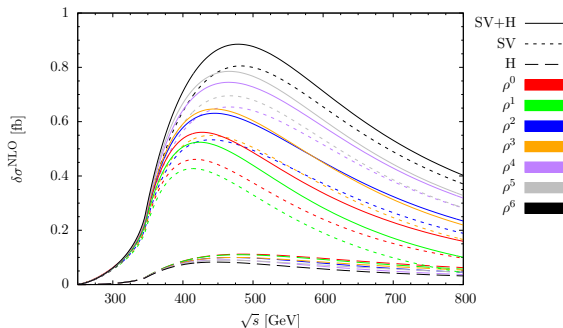
Virt. NNLO master integrals (in addition):



NLO and NNLO Higgs pair production

Splitting in soft-virtual and real contributions at NLO

■ Partonic NLO correction; total factorization



Using $\mu = 2m_H$ (also in the following)

■ Different behavior for higher orders in ρ expansion:

SV increasing

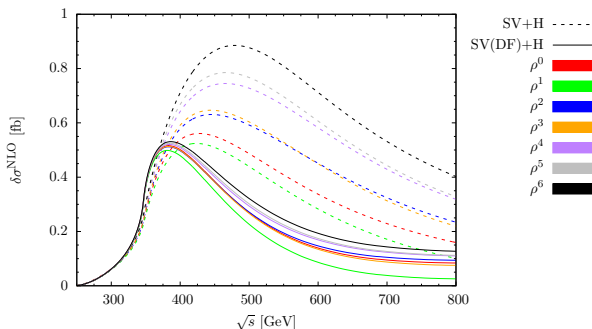
H decreasing (flat for $\sqrt{s} \gtrsim 400$ GeV)

⇒ SV numerically dominant

NLO and NNLO Higgs pair production

Total vs. differential factorization (DF) at NLO

- DF applied only to SV part;
H treated via total factorization (i.e. identical)



- Maxima of DF curves at lower \sqrt{s} ; smaller cross sections
- ⇒ Improvement of convergence:
difference of ρ^0 and corrections to ρ^6 (for $\sqrt{s} = 400$ GeV):
0.25 fb vs. 0.05 fb
- (Partonic K-factor: behavior at top quark pair threshold not washed out)

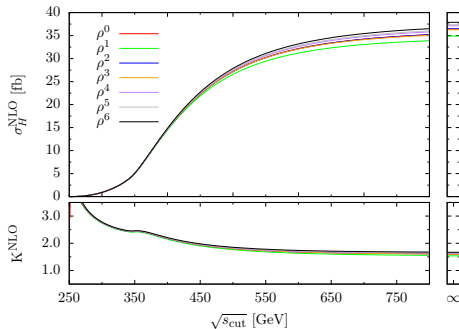
NLO and NNLO Higgs pair production

Hadronic cross section and K-factor

- Technical upper cut on \sqrt{s} (good proxy to Q^2):

$$\sigma_H(s_H, s_{\text{cut}}) = \int_{4m_H^2/s_H}^1 d\tau \left(\frac{d\mathcal{L}_{gg}}{d\tau} \right) (\tau) \sigma(\tau s_H) \theta(s_{\text{cut}} - \tau s_H)$$

- $\sqrt{s_{\text{cut}}} \rightarrow \infty$: total cross section for 14 TeV



- Spread of ρ -orders \Rightarrow $\pm 10\%$ uncertainty of EFT at NLO due to M_t

NLO and NNLO Higgs pair production

Revisiting NLO

Lessons from NLO for NNLO:

- SV approximation constructed for $z \rightarrow 1$;
 $G_{SV}(z)$ can be replaced by $f(z) G_{SV}(z)$

Splitting into SV and H not unique

- Tune $f(z)$ at NLO such that $\sigma \approx \Sigma_{SV}$

$\Rightarrow f(z) = z$ accurate to 2%

- Replace RGE logarithms ($\sqrt{s} \approx Q^2$ in the soft limit):

$\Rightarrow \log(\mu^2/s) \rightarrow \log(\mu^2/Q^2)$

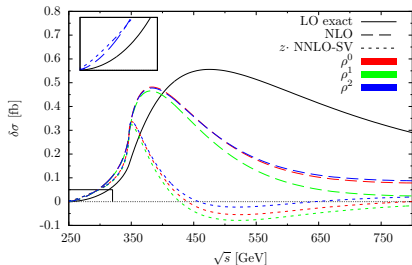
Discrepancy to [\[Maltoni, Vryonidou, Zaro; '14\]](#):

- Real corr.: treated exactly; Virt. corr.: EFT result
- Claim: **-10%** correction at NLO

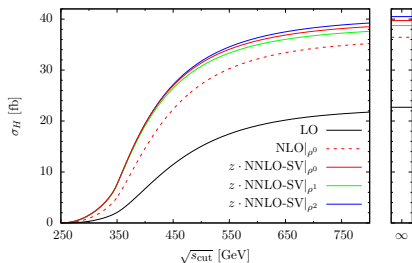
But: Dominant positive shift from virtual $1/M_t$ -corrections (c.f. backup slide)

NLO and NNLO Higgs pair production

NNLO SV corrections



- EFT result plus ρ - and ρ^2 -terms
- Peaks at smaller \sqrt{s}
- Same pattern of ρ -corrections



- Conv. up to $\sqrt{s_{\text{cut}}} \approx 400$ GeV
- ρ - and ρ^2 -corrections: $\pm 2.5\%$

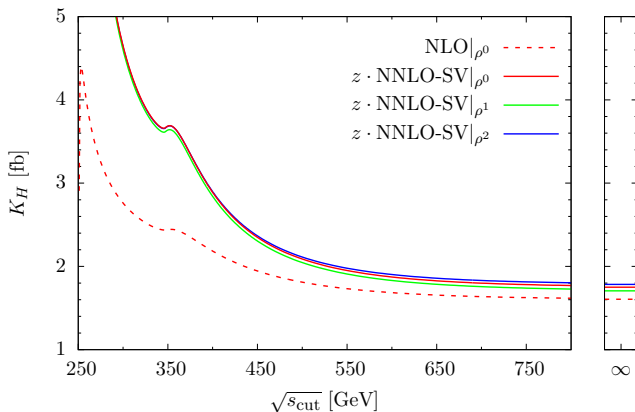
⇒

M_t -uncertainty at NNLO: 5%

(NNLO corr. $\approx 20\%$)

NLO and NNLO Higgs pair production

NNLO SV K-factor



■ Strong raise close to threshold \Leftarrow steeper NNLO correction

■ For total cross section: $K_H^{\text{NNLO}} \approx 1.7 - 1.8$

Conclusion

TopoID:

- Generic, process independent Mathematica package for multiloop calculations; especially for many topologies
- Until now: two successful applications
- Works also for 5-loop propagators
- Publication planned for ACAT 2016

qq' -channel in N^3 LO Higgs production:

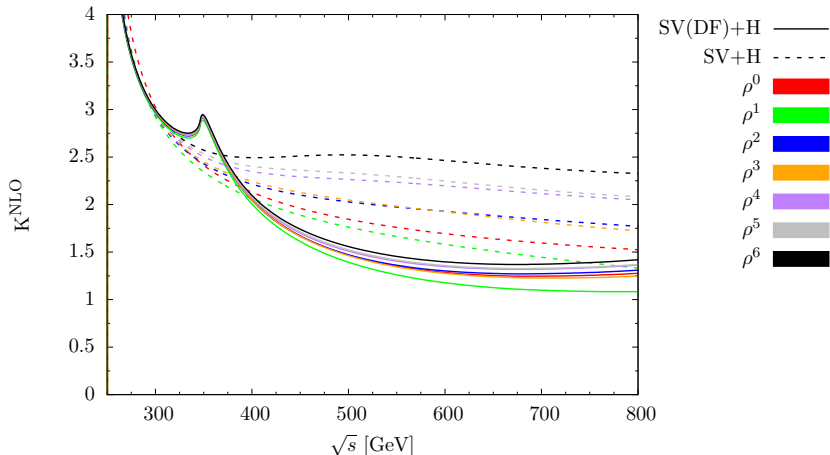
- New iterated integrals beyond HPLs appear
- Agreement of leading logarithms with
[\[Anastasiou, Duhr, Dulat, Furlan, Herzog, Mistlberger; '15\]](#)
- Full calculation underway . . .

Higgs pair production:

- Real corrections at NLO via opt. theorem; employ SV approx. at NNLO
- Computed expansion to $\mathcal{O}(1/M_t^{12})$ at NLO, $\mathcal{O}(1/M_t^4)$ at NNLO
- Uncertainty of EFT results due to M_t : 10% at NLO, 5% at NNLO
- Small- \sqrt{s} behavior as benchmark for exact calculations

Backup: NLO and NNLO Higgs pair production

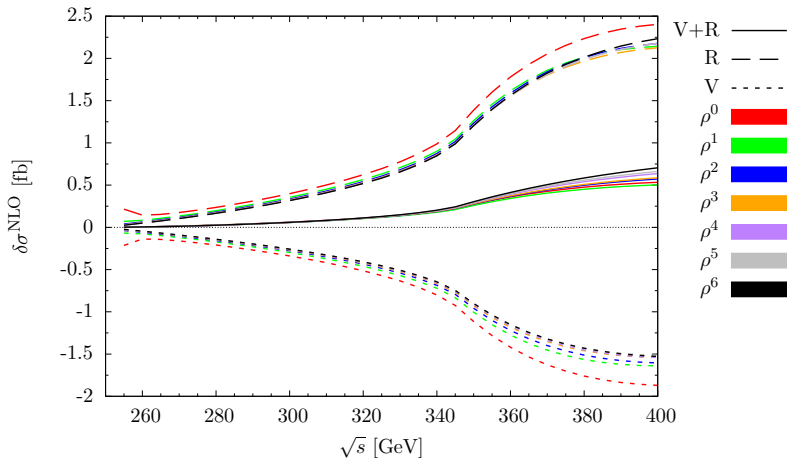
Partonic NLO K-factor



Note: Behavior around top quark pair threshold not washed out

Backup: NLO and NNLO Higgs pair production

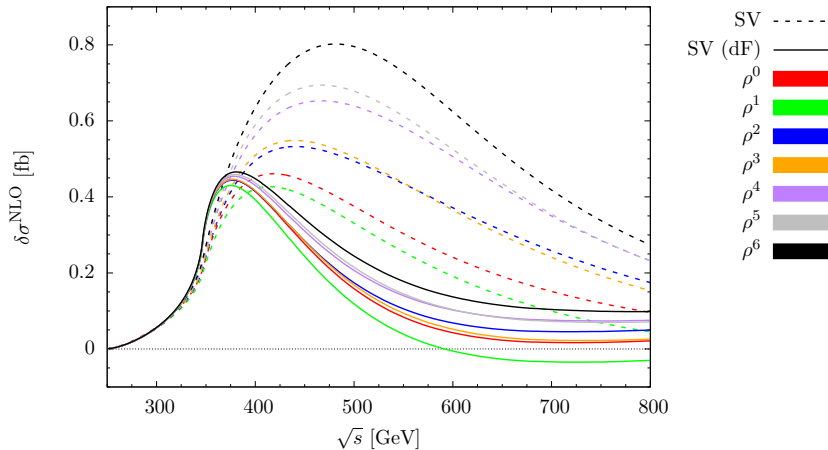
Splitting into real and virtual corrections at NLO



Note: R and V separately divergent; only finite contributions shown

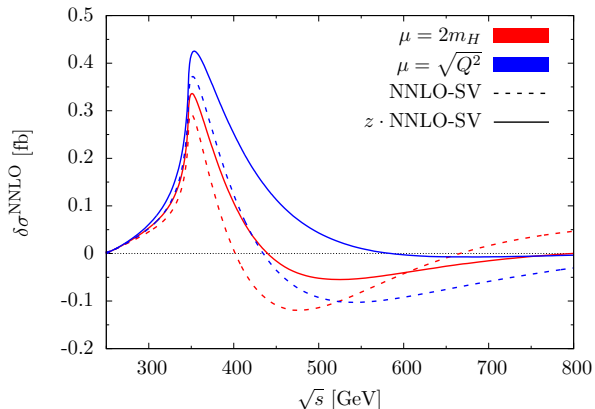
Backup: NLO and NNLO Higgs pair production

Total vs. differential factorization without hard contributions at NLO



Backup: NLO and NNLO Higgs pair production

Partonic NNLO cross section for different scales choices and $f(z)$



Note: $f(z) = z$ (better proxy) leads to higher values